

General Relativistic neutrino-Radiation MagnetoHydroDynamics (GRRMHD) simulations of binary neutron star mergers

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GW170817: dawn of Multi-messenger astrophysics with GW

- **Detection of GW170817**

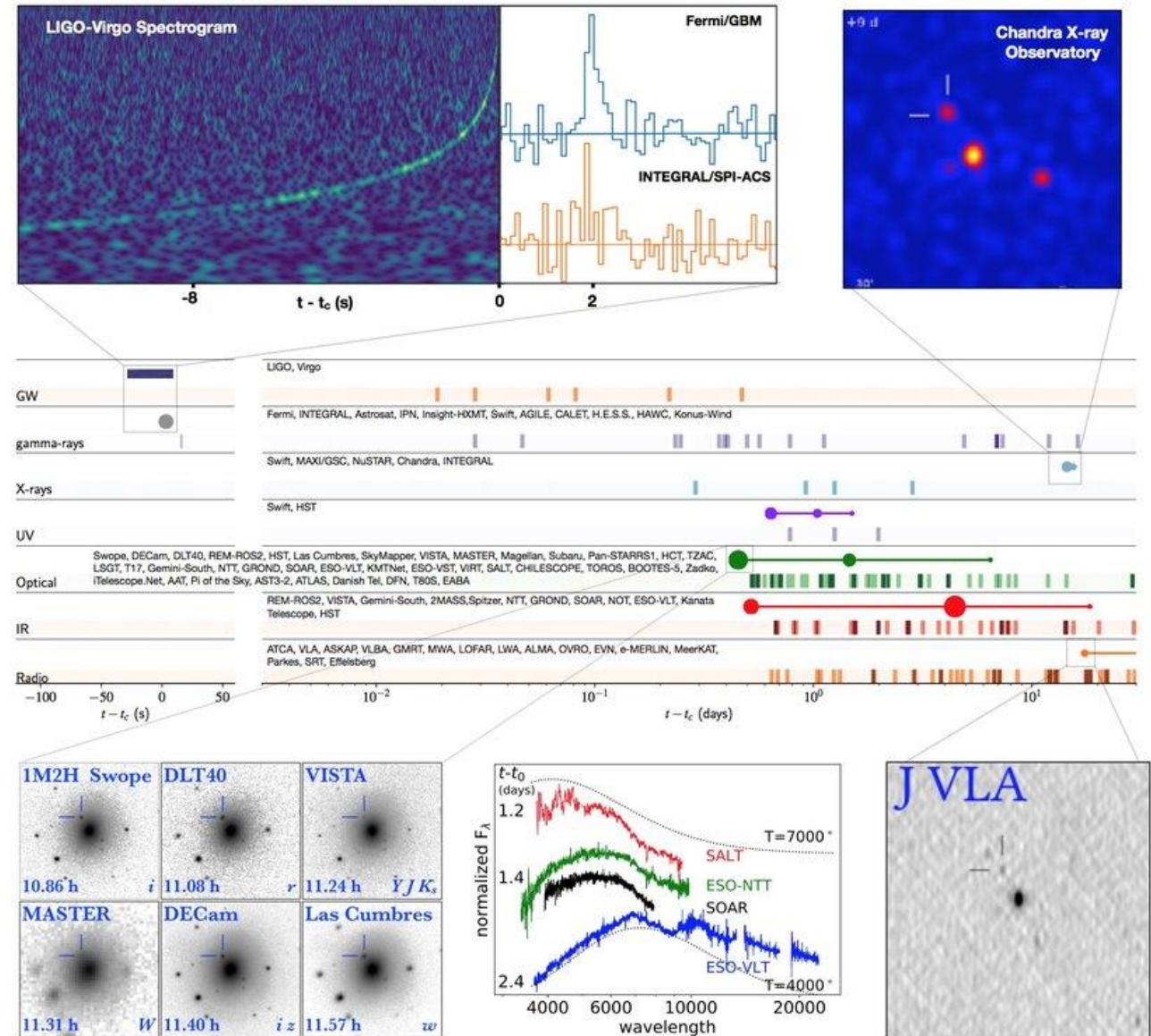
- constraint on neutron star (NS) equation of state by tidal deformability of NS in the late inspiral stage
- measurement of Hubble constant

- **Observation of AT2017gfo**

- the origin of r-process elements like rare earth elements (Lanthanides), Au, Pt, and U is likely to be binary NS

- **Association of GRB170817**

- the central engine of (at least a part of) short hard GRB is binary neutron star merger



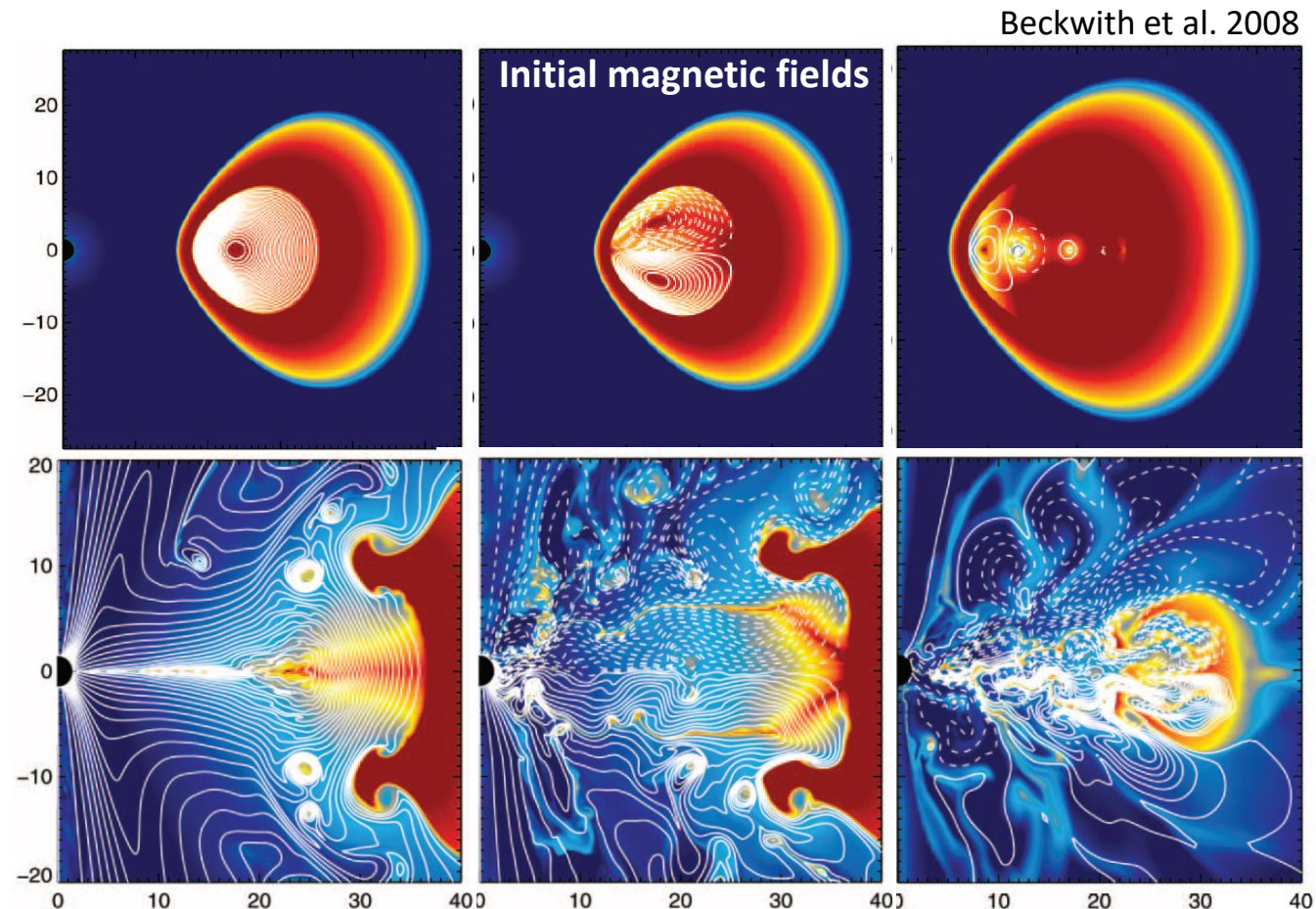
How to drive a short GRB jet ?

- **Blandford-Znajek mechanism** is a promising mechanism to launch the short GRB jet
- **Strong ($\gtrsim 10^{15}$ G) and coherent magnetic fields** which thread the BH horizon are necessary to launch an energetic jet
- Poloidal magnetic fields of binary pulsars estimated by the spin-down period : $B_p \sim 10^{8-12}$ G $\ll 10^{15}$ G

Tauris et al. 2017

Key question

How to make such a strong coherent magnetic field from NS magnetic field



Generation of coherent magnetic fields

e.g., Moffatt (1978) "Magnetic field generation in electrically conducting fluids"

The averaged induction equation

$$\partial_t \bar{\mathbf{B}} = \nabla \times (\bar{\mathbf{U}} \times \bar{\mathbf{B}} + \bar{\boldsymbol{\varepsilon}})$$

Electromotive force

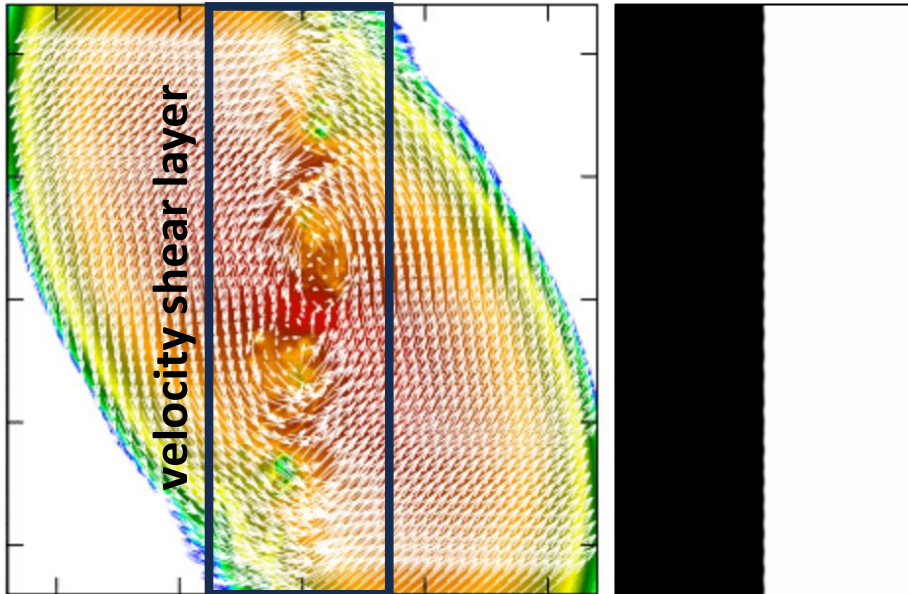
$$\bar{\boldsymbol{\varepsilon}} = \overline{\mathbf{u} \times \mathbf{b}}$$

mean field random field

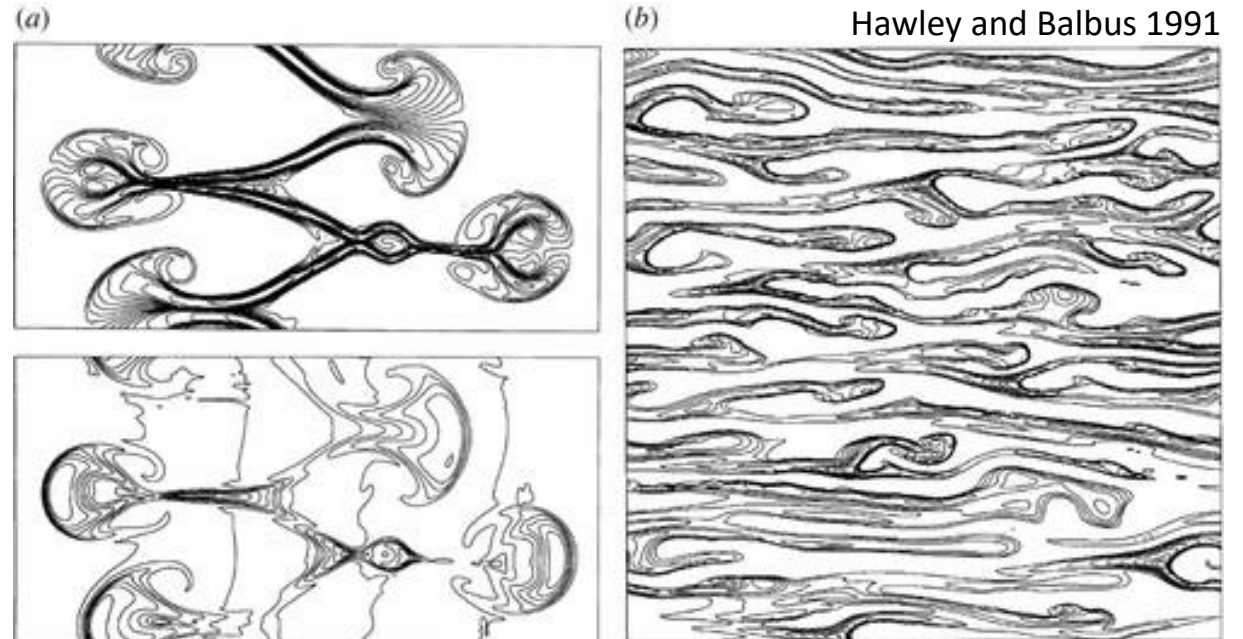
$$\mathbf{Q} = \overline{\mathbf{Q}} + \mathbf{q}$$

⇒ small scale turbulent velocity and magnetic fields can generate coherent fields

Kiuchi et al. 2015, 2018



Kelvin-Helmholtz (KH) instability at the contact shear



Magneto-rotational instability (MRI) in the torus

Generation of coherent magnetic fields

e.g., Moffatt (1978) "Magnetic field generation in electrically conducting fluids"

The averaged induction equation

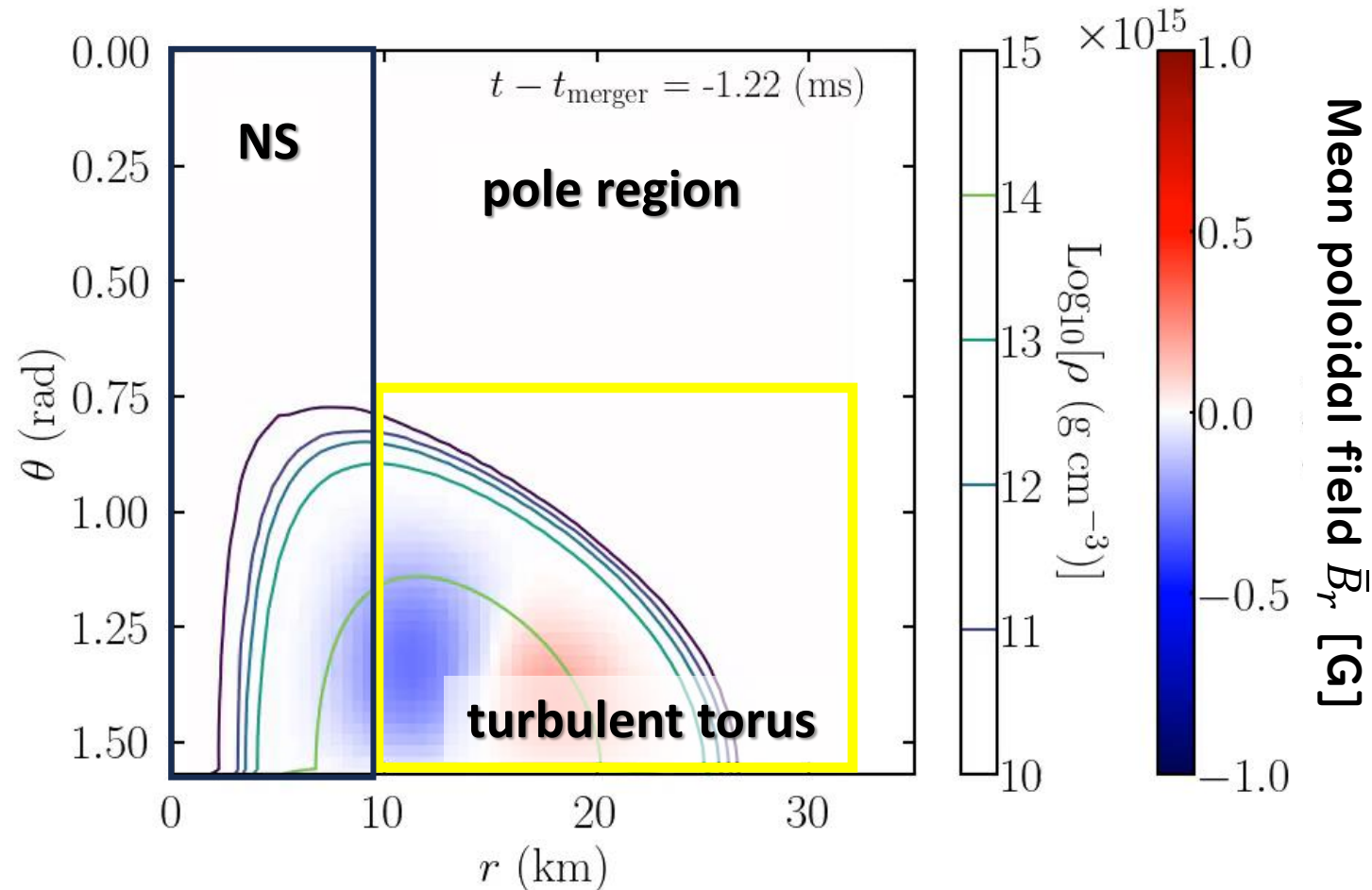
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mean field random field

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Phenomenological Prescription

$$\bar{\boldsymbol{\varepsilon}}_i = \alpha_{ij} \bar{B}_j + \beta_{ij} (\nabla \times \bar{\mathbf{B}})_j$$

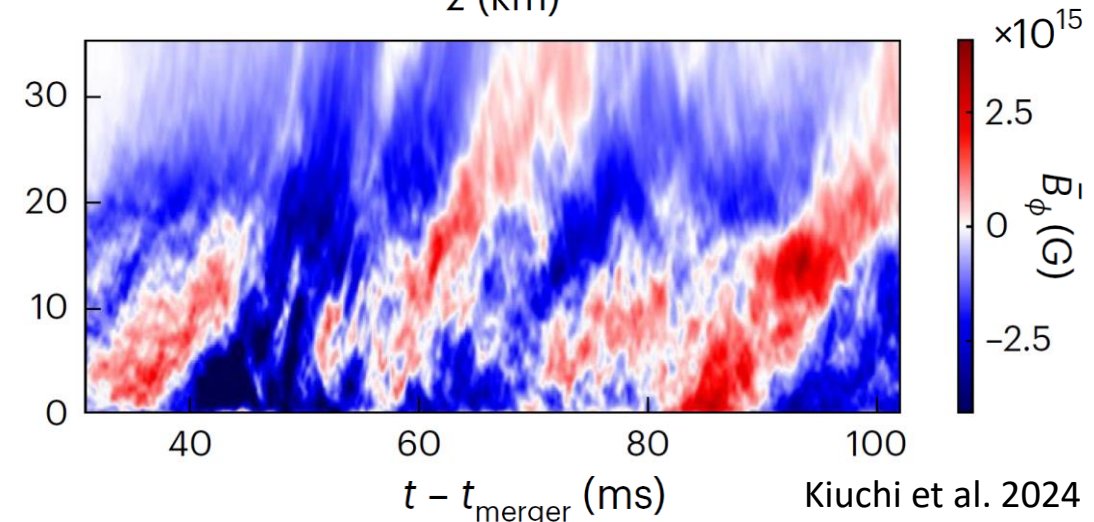
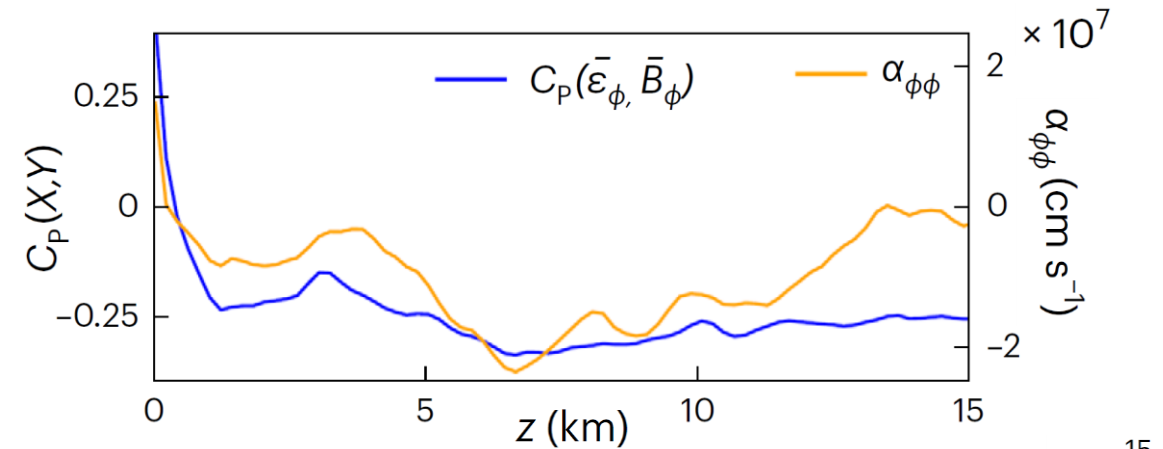
$$\partial_t \bar{B}_R = -\partial_z (\alpha_{\phi\phi} \bar{B}_\phi) \quad R = \sqrt{x^2 + y^2}$$

$$\partial_t \bar{B}_\phi = r \sin \theta (\bar{\mathbf{B}}_{\text{pol}} \cdot \nabla \Omega)$$

$\alpha\Omega$ dynamo

- ✓ (anti-correlation between \bar{B}_ϕ and $\bar{\boldsymbol{\varepsilon}}_\phi$)
- ✓ dynamo-wave propagate to $\alpha_{\phi\phi} \nabla \Omega \times \mathbf{e}_\phi \propto -\mathbf{e}_\theta$
- ✓ dynamo-cycle period :

$$\tau_{\text{dynamo}} \sim \sqrt{2\pi} \left(\alpha_{\phi\phi} k_z \frac{d\Omega}{d \ln R} \right)^{-1/2} \sim \text{a few } 10 \text{ ms}$$



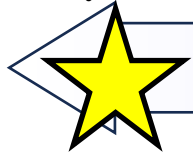
Kiuchi et al. 2024

prompt collapse

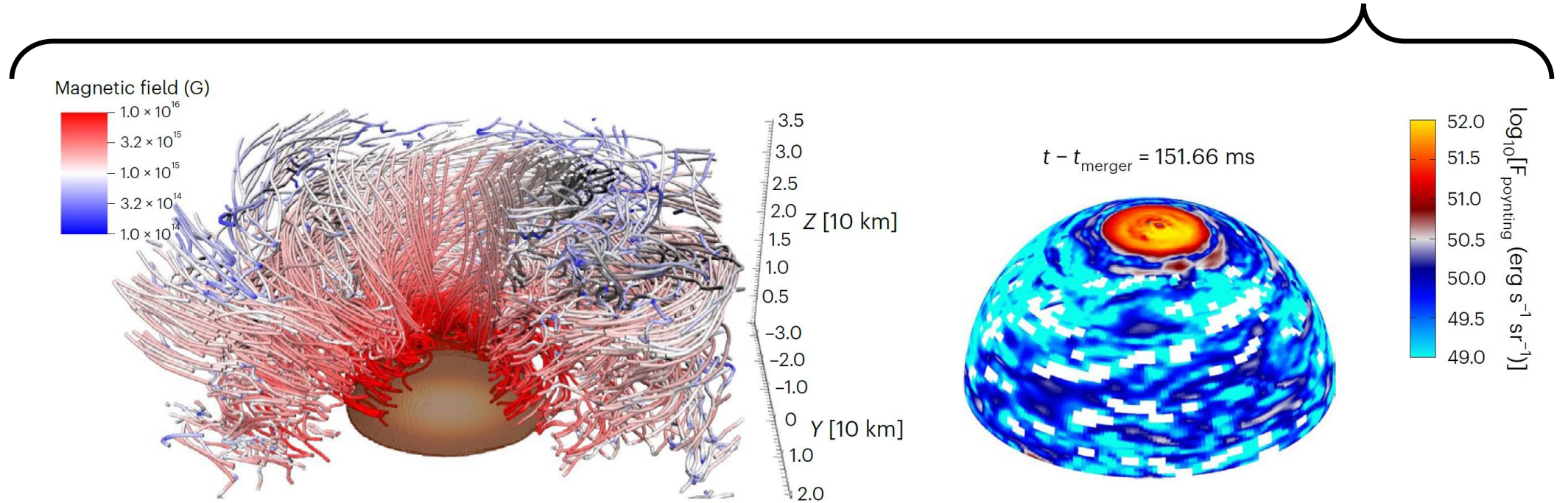
$O(0.01 \text{ sec})$

$O(0.1 \text{ sec})$

$< O(1 \text{ sec})$



Remnant NS lifetime



- ✓ Magnetic field amplification by Kelvin-Helmholtz instability and magneto-rotational instability
- ✓ Subsequent mean field generation by $\alpha\Omega$ dynamo
- ✓ Collimated ($\theta_{\text{jet}} \approx 12^\circ$), Poynting flux dominated jet launched with $L_{\text{poy}} \sim 10^{51} \text{ erg/s}$ (this is NOT the isotropic-equivalent luminosity)
- ✓ Mildly neutron-rich ($X_{n,\text{ave}} \sim 0.7$) ejecta with $M_{\text{ej}} \geq 0.1 M_{\odot}$

GW190425 and Prompt collapse to a BH

Brief summary of GW190425

Abbott et al. 2020

- ✓ Total mass of BNS : $M_{\text{total}} = 3.3 \sim 3.4 M_{\odot}$
 \Rightarrow expected to collapse promptly to a BH
- ✓ Poor sky localization due to a single detector event
- ✓ no electromagnetic counterpart is detected

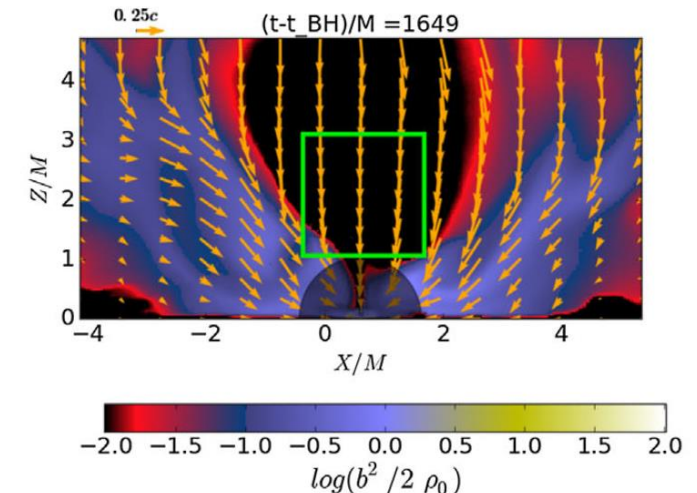
Previous GRMHD simulation for prompt collapse

Ruiz and Shapiro 2017

- ✓ Poynting flux dominated jet are NOT launched
- ✓ No evidence for coherent magnetic field formation
- ✓ (Nearly) **equal mass binary** \Rightarrow small disk mass $\lesssim 10^{-3} M_{\text{tot}}$
- ✓ **Short-term simulations** up to 26 ms after the merger

Source Properties for GW190425

	Low-spin Prior ($\chi < 0.05$)	High-spin Prior ($\chi < 0.89$)
Primary mass m_1	1.60–1.87 M_{\odot}	1.61–2.52 M_{\odot}
Secondary mass m_2	1.46–1.69 M_{\odot}	1.12–1.68 M_{\odot}
Chirp mass \mathcal{M}	$1.44^{+0.02}_{-0.02} M_{\odot}$	$1.44^{+0.02}_{-0.02} M_{\odot}$
Detector-frame chirp mass	$1.4868^{+0.0003}_{-0.0003} M_{\odot}$	$1.4873^{+0.0008}_{-0.0006} M_{\odot}$
Mass ratio m_2/m_1	0.8 – 1.0	0.4 – 1.0
Total mass m_{tot}	$3.3^{+0.1}_{-0.1} M_{\odot}$	$3.4^{+0.3}_{-0.1} M_{\odot}$
Effective inspiral spin parameter χ_{eff}	$0.012^{+0.01}_{-0.01}$	$0.058^{+0.11}_{-0.05}$
Luminosity distance D_L	159^{+69}_{-72} Mpc	159^{+69}_{-71} Mpc
Combined dimensionless tidal deformability $\tilde{\Lambda}$	≤ 600	≤ 1100



Set-up of simulation

- **Einstein's equations :**

- ✓ BSSN formalism (Shibata and Nakamura 1995; Baumgarte and Shapiro 1998)
- ✓ Moving puncture method (Campanelli et al. 2006; Baker et al. 2006)
- ✓ Z4c constraint propagation (Hilditch et al. 2013)

- **Magnetohydrodynamics :** (Kiuchi et al. 2022)

- ✓ HLLD Riemann solver (Mignone et al. 2009)
- ✓ Divergence-B constraint transport (Gardiner and Stone 2008)
- ✓ Magnetic-flux preserving mesh refinement (Balsara 2009)

- **Neutrino transfer :** (Sekiguchi et al. 2012)

- ✓ M1 closure (Shibata et al. 2011)
- ✓ Neutrino heating (Fujibayashi et al. 2017)

- **Prescription of BNS :**

- ✓ SFHo equation of state (Steiner et al. 2013) : $M_{\max} \approx 2.1 M_{\odot}$
- ✓ $1.25 M_{\odot} - 1.65 M_{\odot}$ unequal mass binary ($M_{\text{tot}} = 2.9 M_{\odot}$)
- ✓ prompt collapse to a BH with $M_{\text{BH}} \approx 2.8 M_{\odot}$, $a_{\text{BH}} = 0.76$
- ✓ accretion disk with $M_{\text{disk}} \approx 0.06 M_{\odot}$ is formed

- **Magnetic field :**

- ✓ poloidal magnetic field is superimposed inside the NSs

$$A_j = A[(x - x_{\text{NS}})\delta_j^y - (y - y_{\text{NS}})\delta_j^x] \cdot \max(P/P_{\text{max}} - 2 \cdot 10^{-4}, 0)^{1/2}$$

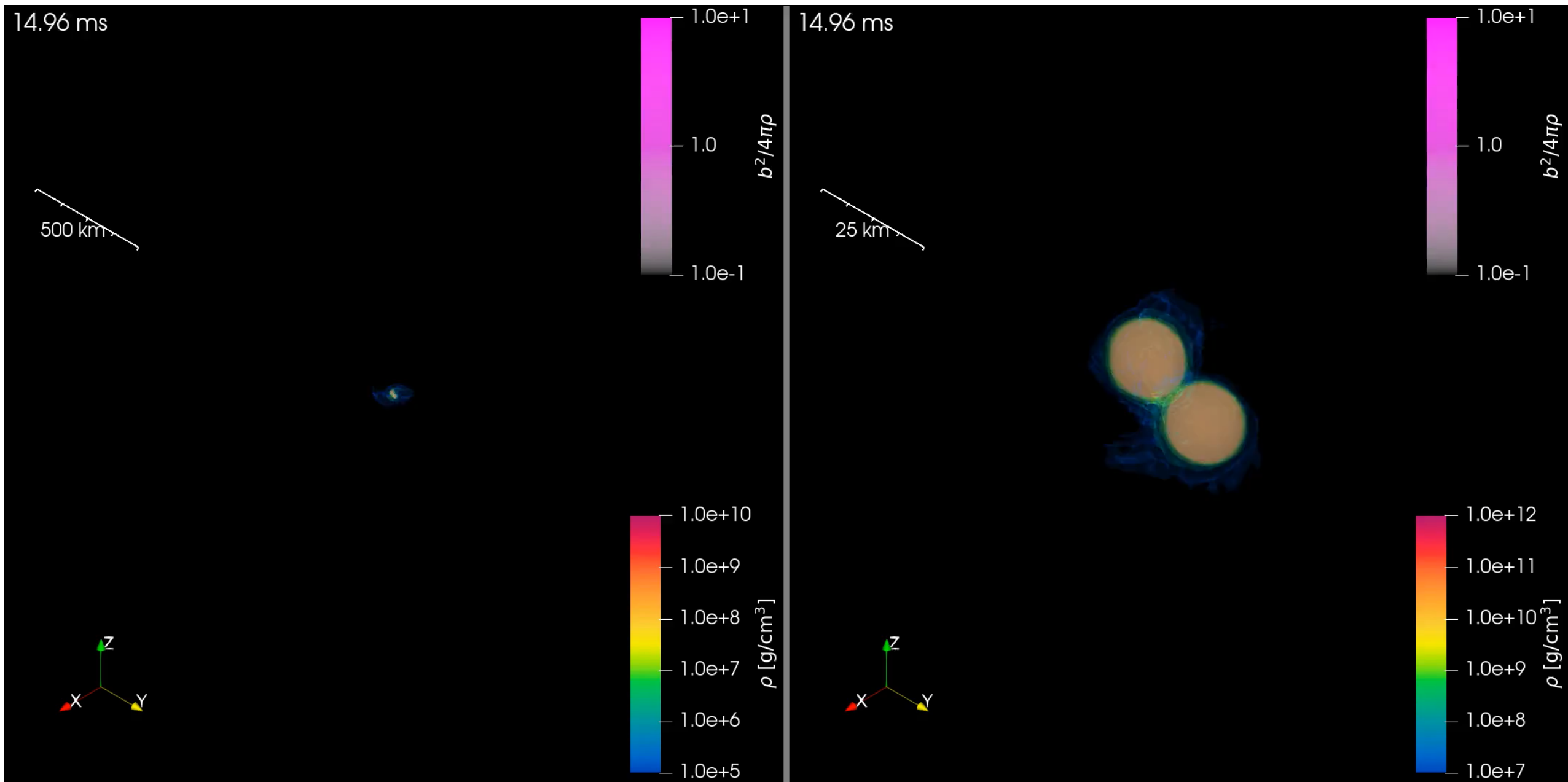
- ✓ maximum field strength is $\approx 10^{15}$ G

- **Grid set-up and timescale :**

- ✓ 13-level fixed mesh refinement
- ✓ finest grid resolution : $\Delta x = 150$ m enable to follow the fastest growing mode of magneto-rotational instability

$$\lambda_{\text{MRI}} \sim \frac{v_{\text{Alfven}}}{\Omega} \sim \frac{B}{\Omega \sqrt{4\pi\rho}}$$

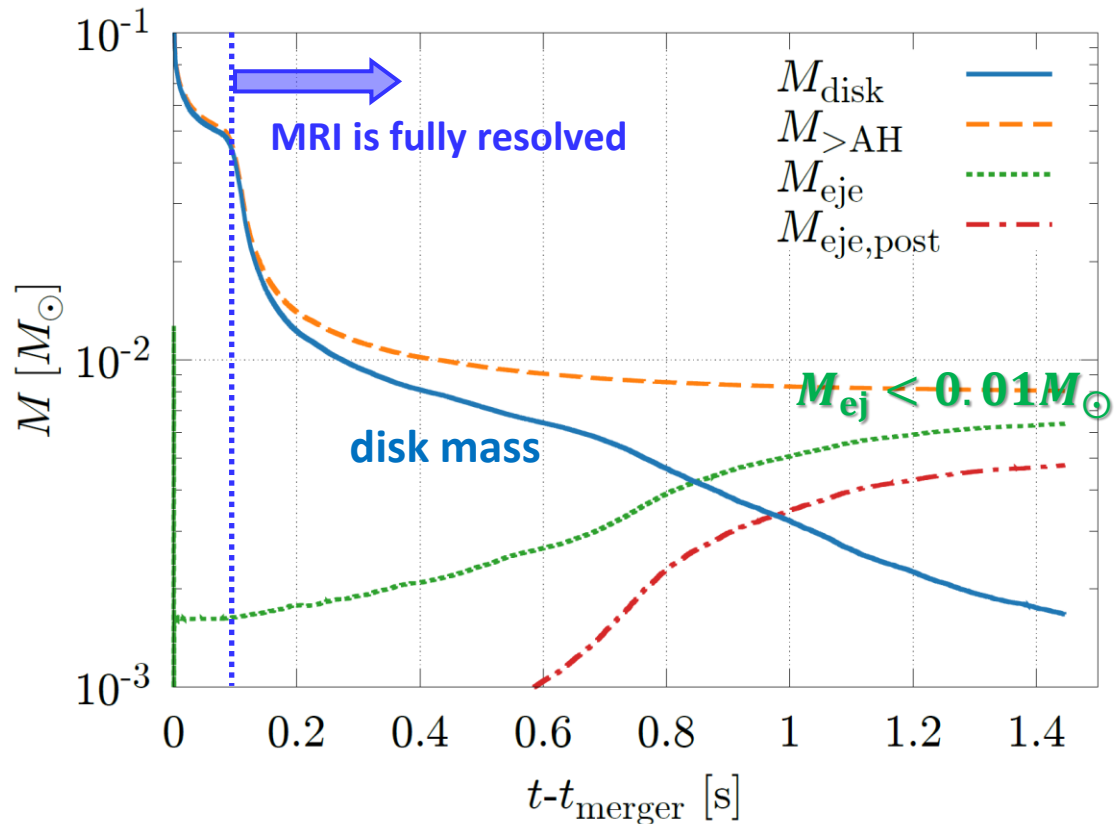
- ✓ Long-term (> 1 sec) simulation



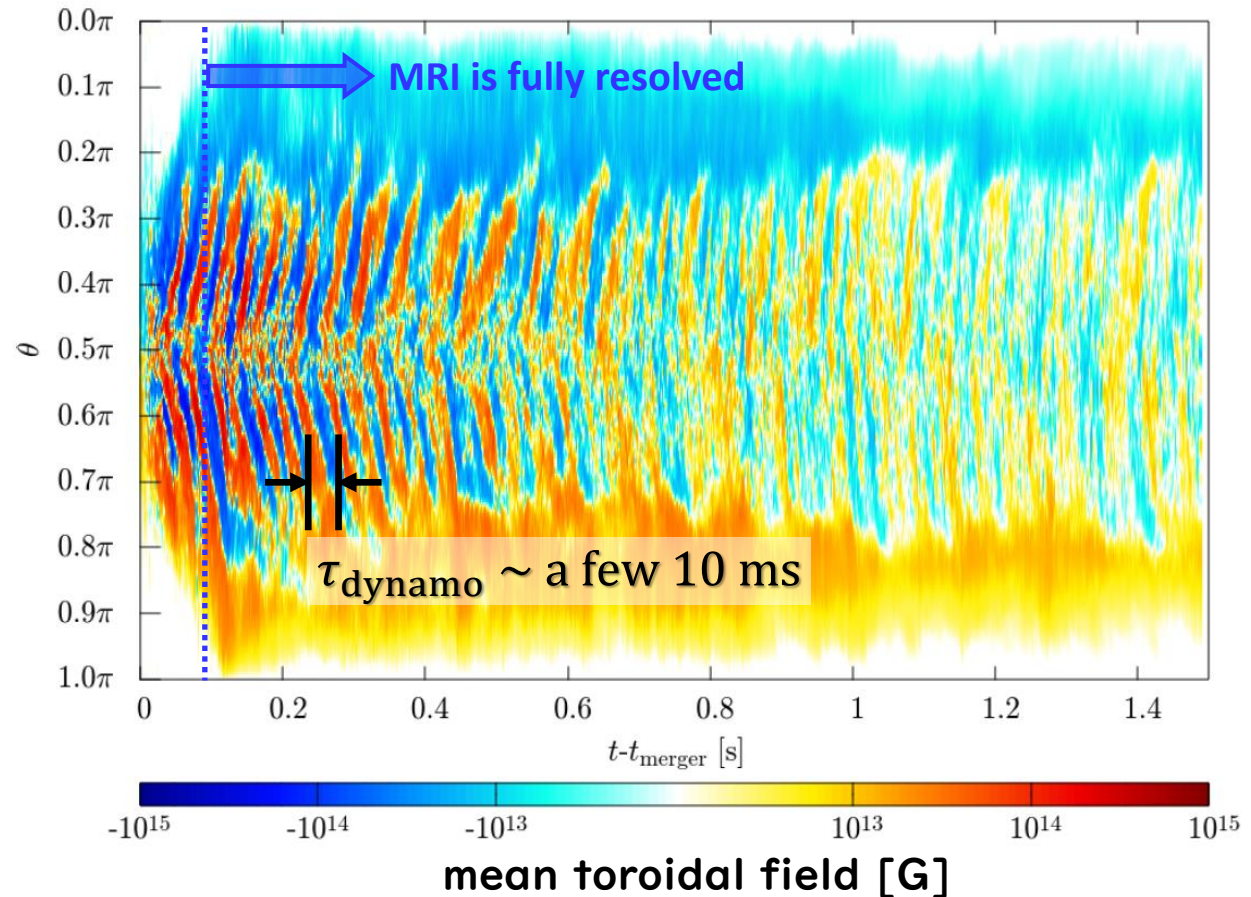
MRI induced viscosity and dynamo

The fastest growing mode is resolved : **partially** ($t - t_{\text{merger}} \gtrsim 10 \text{ ms}$), **fully** ($t - t_{\text{merger}} \gtrsim 100 \text{ ms}$)

- ✓ MRI driven turbulence induces effective viscosity and disk mass decays

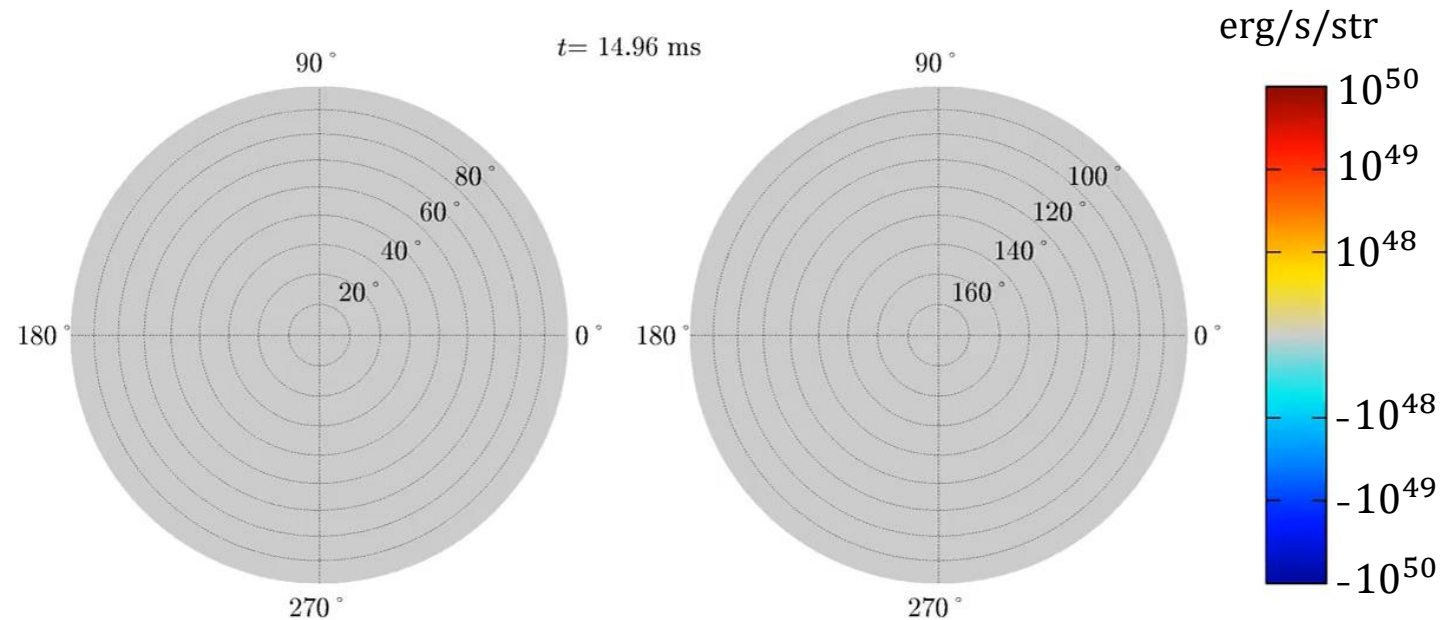
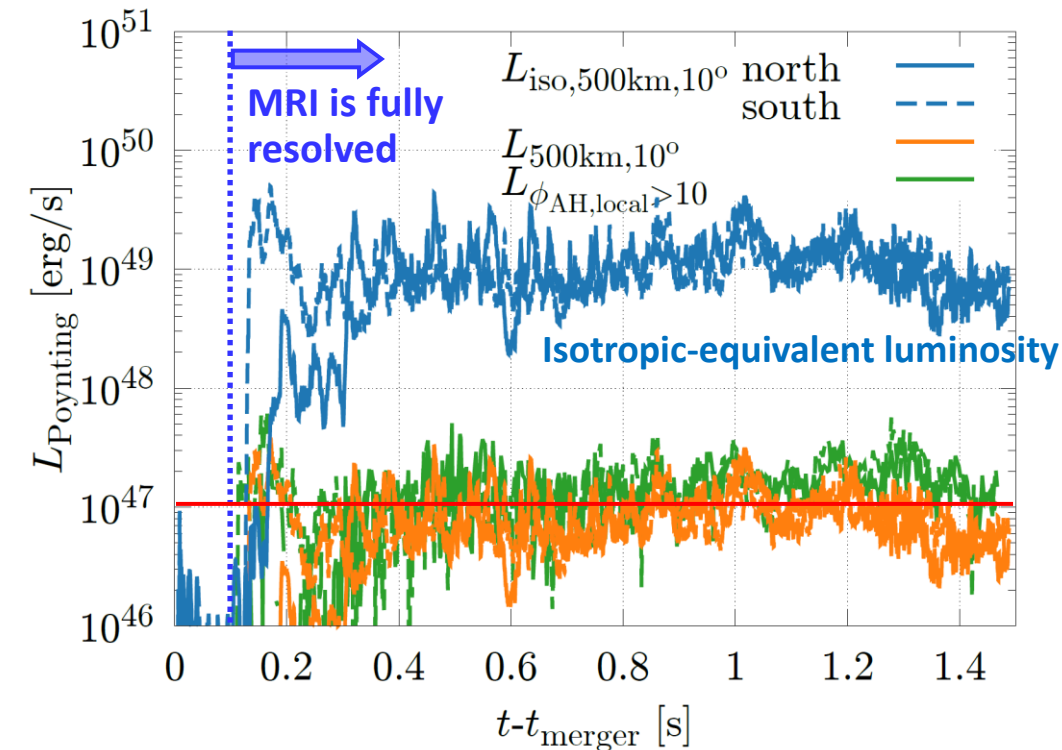


- ✓ MRI driven turbulence activates the dynamo cycle and coherent magnetic fields are formed



(Weak) Poynting flux dominated jet is launched

- ✓ Coherent magnetic fields accrete the BH and further amplified by winding
 - ✓ Note also that Prompt collapse to BH \Rightarrow density and ram pressure in the pole region is smaller
- ✓ Collimated ($\theta_{\text{jet}} \sim 10^\circ$), Poynting flux dominated jet launched with $L_{\text{Poy}} \sim 10^{47}$ erg/s
 - ✓ Jet angle gradually increases as matter pressure, which confines the jet, decreases



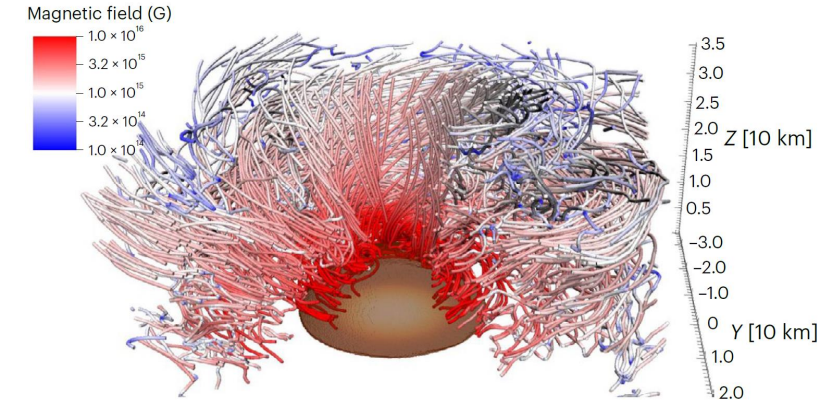
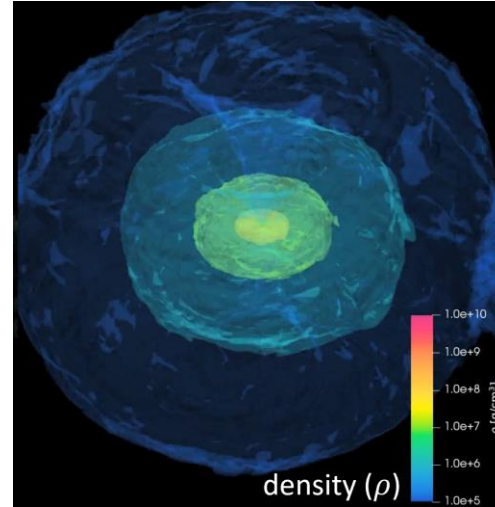
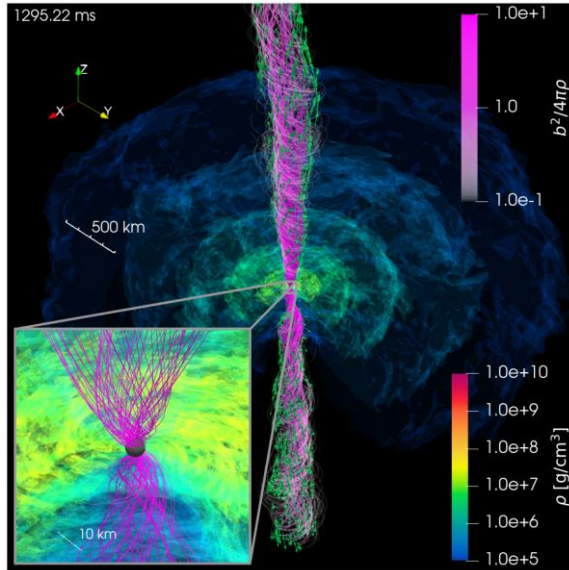
prompt collapse

$O(0.01 \text{ sec})$

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Remnant NS lifetime



✓ **Prompt BH formation**

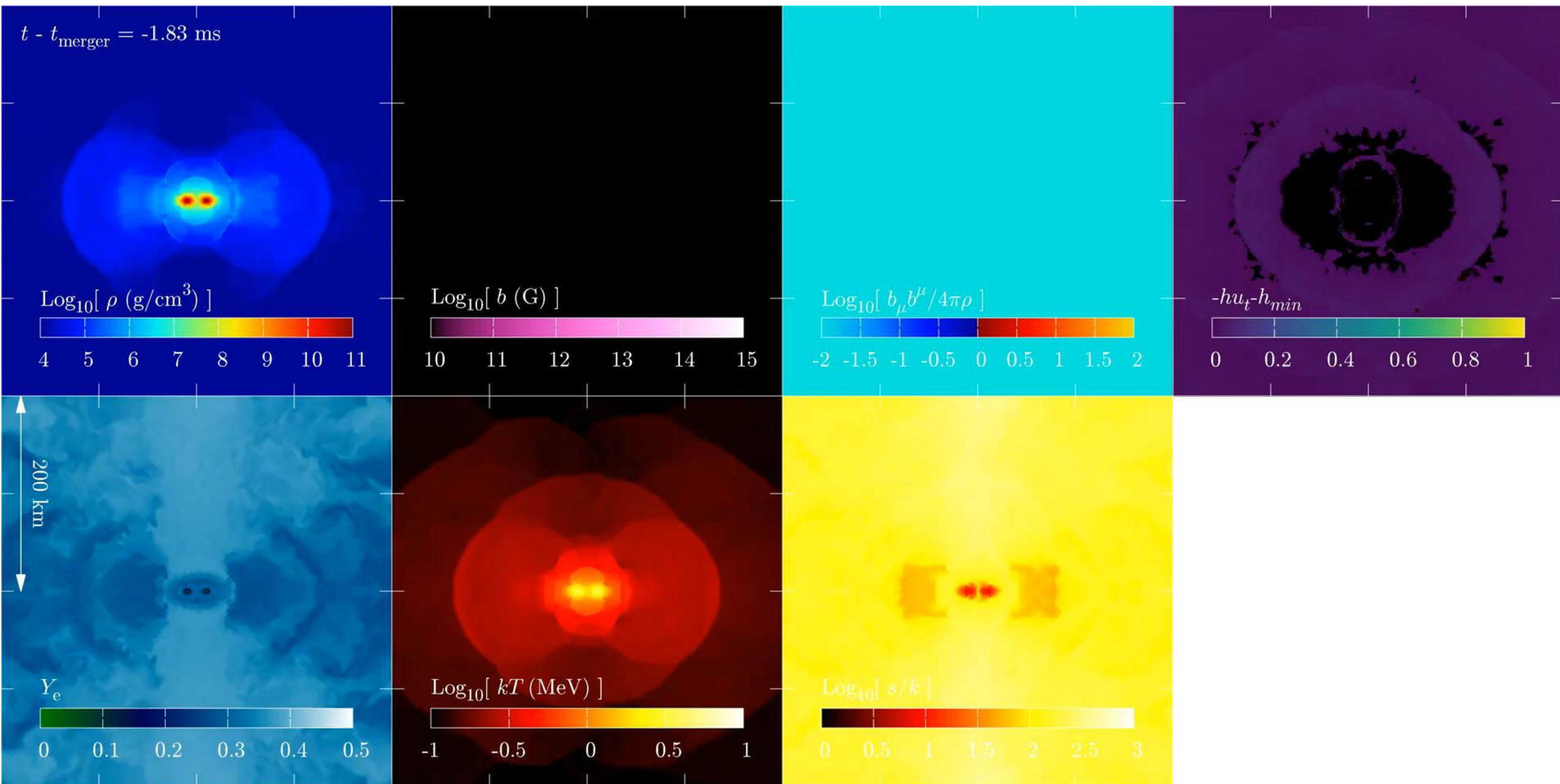
- ✓ mean field generation by MRI-induced dynamo
- ✓ ram pressure is smaller
- ✓ $L_{\text{Poy}} \sim 10^{47} \text{ erg/s}$, $\theta_{\text{jet}} \sim 10^\circ$
- ✓ $M_{\text{ej}} \lesssim 0.01 M_\odot$

✓ **delayed BH formation**

- ✓ MRI is resolved
- ✓ ram pressure is stronger than magnetic pressure
- ✓ Jet is not launched in 1 sec after the merger

✓ **Long-lived NS** : amplification

- by KH instability and MRI
- ✓ mean field generation by $\alpha\Omega$ dynamo
- ✓ $L_{\text{Poy}} \sim 10^{51} \text{ erg/s}$, $\theta_{\text{jet}} \approx 12^\circ$
- ✓ $M_{\text{ej}} \geq 0.1 M_\odot$ ($X_{n,\text{ave}} \sim 0.7$)



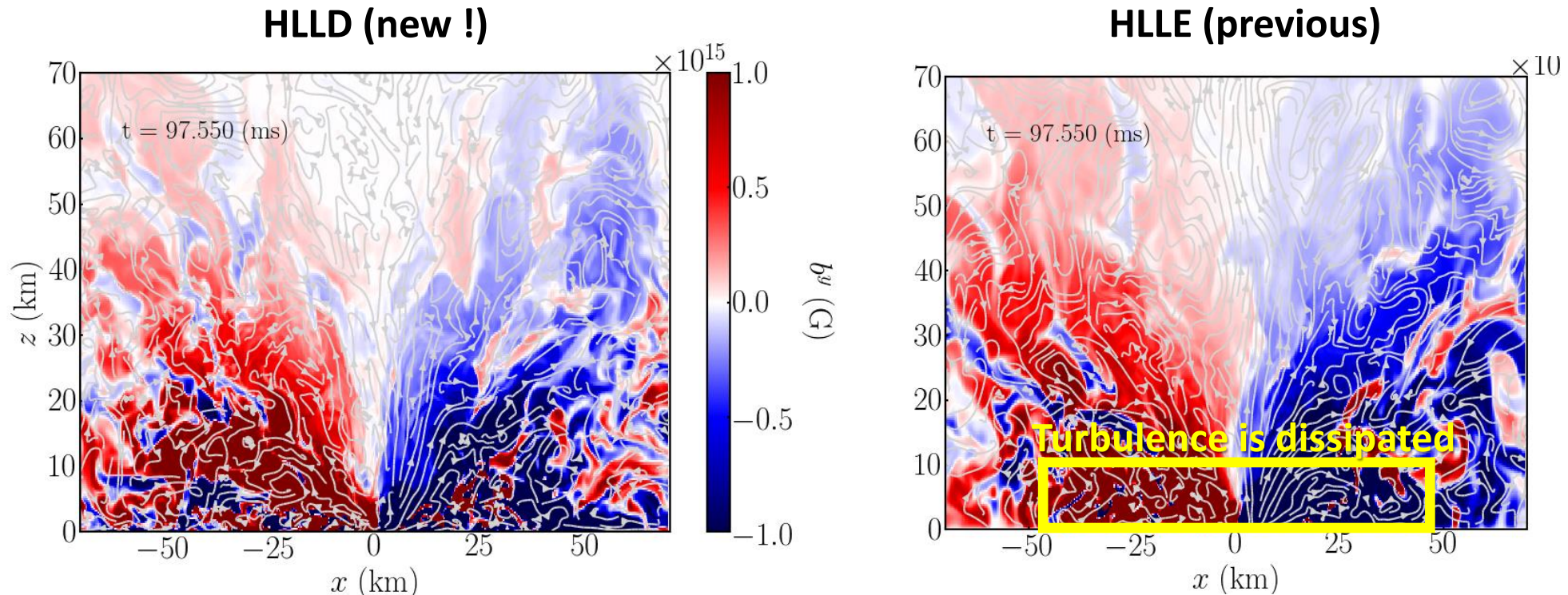
Technical issues to follow small-scale fields

- High resolution is required

- we adopted $\Delta x_{\text{finest}} = 12.5\text{m}$ (previously $\Delta x_{\text{finest}} = 150\text{m}$) to resolve the fastest growing mode of MRI and accurately follow the B-field amplification in KH instability

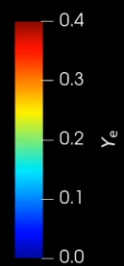
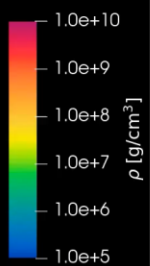
- Less dissipative MHD solver is advantageous

- we developed less dissipative HLLD solver (Kiuchi, YS+ 2022) in the framework of NR



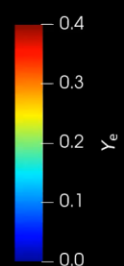
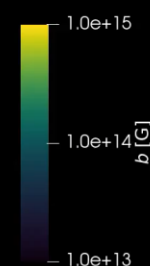
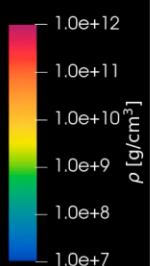
14.96 ms

500 km



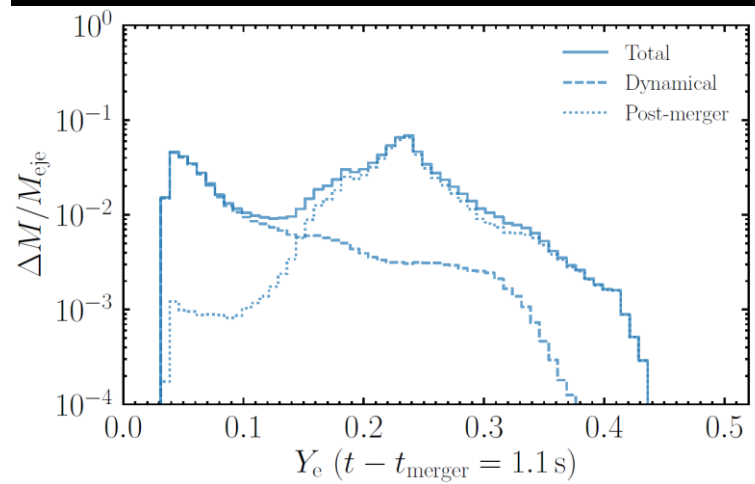
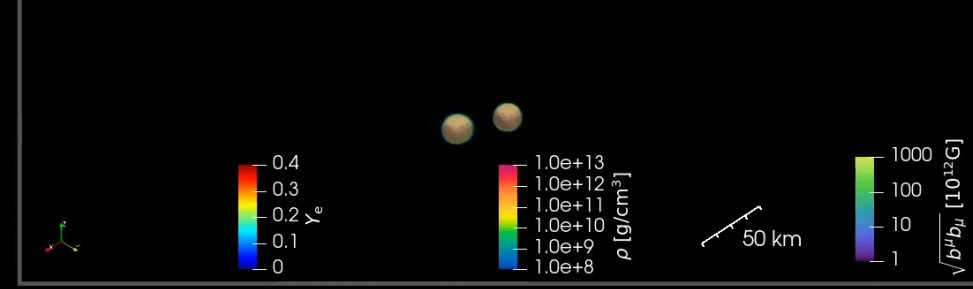
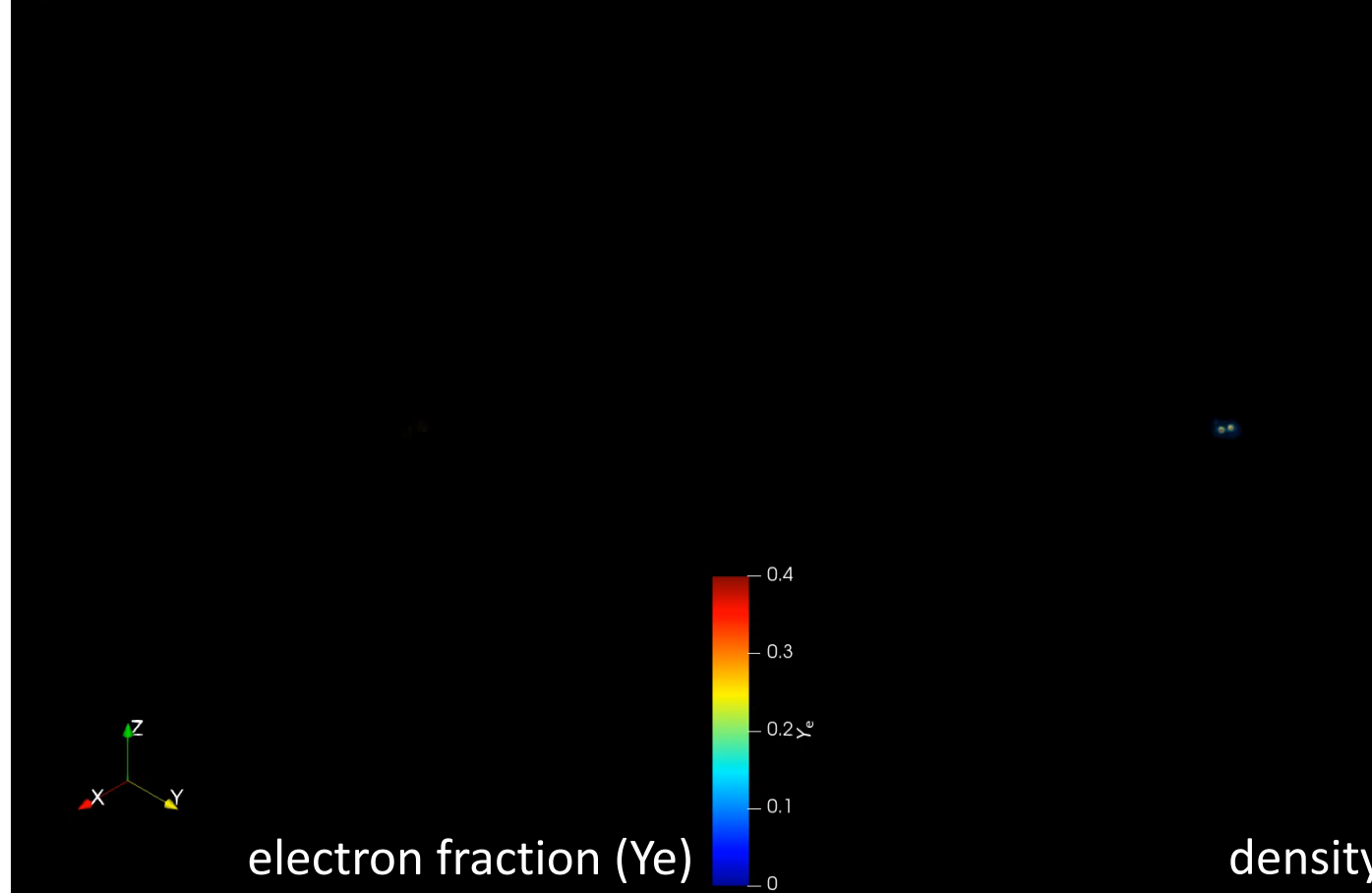
14.96 ms

25 km

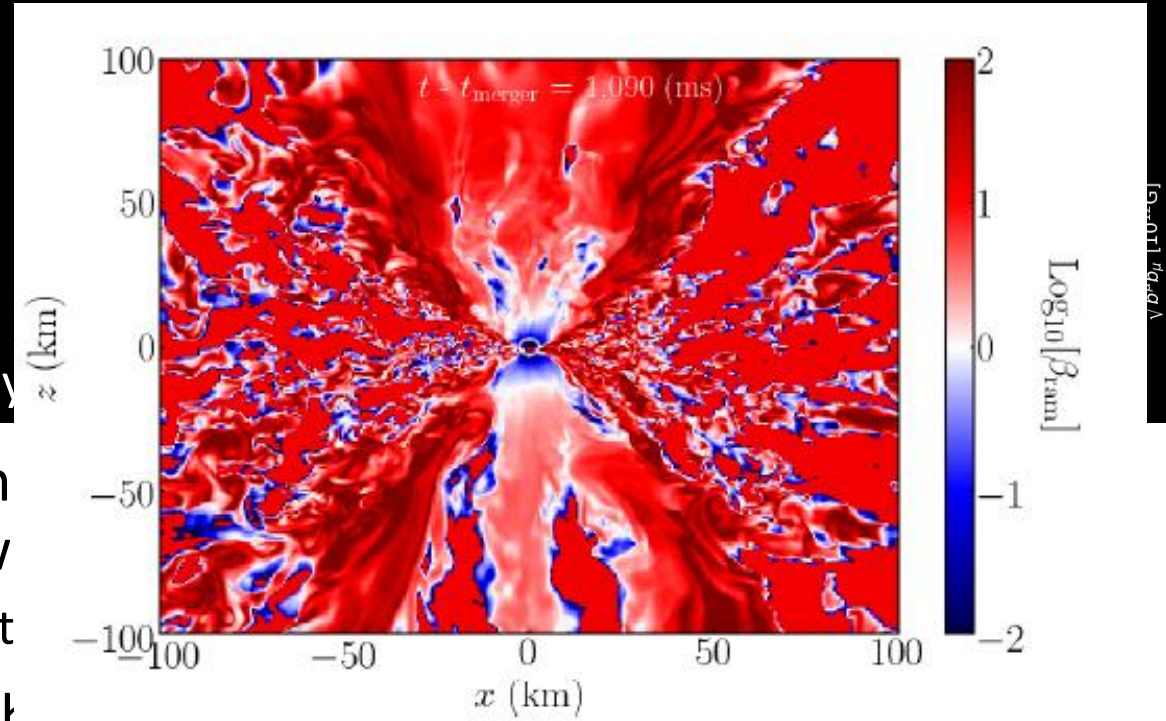


Time: 7.52 ms

Time: 7.52 ms



- ejecta properties (ch broadly consistent w
 - ejecta mass may not
- Absence of jet launch
 - **Insufficient coherent B-field ?**, too large ram pressure due to shorter simulation time ?



Result for $1.35-1.35M_{\odot}$ with DD2 EOS

