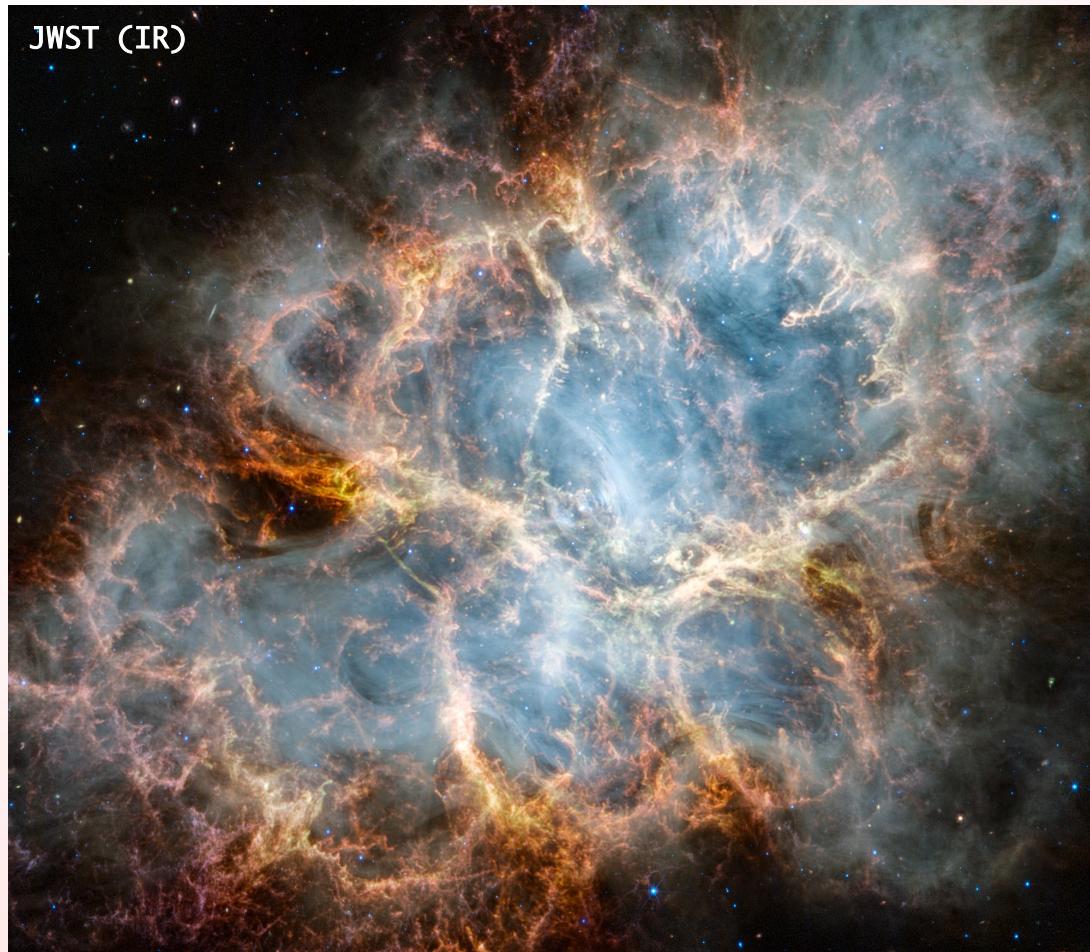


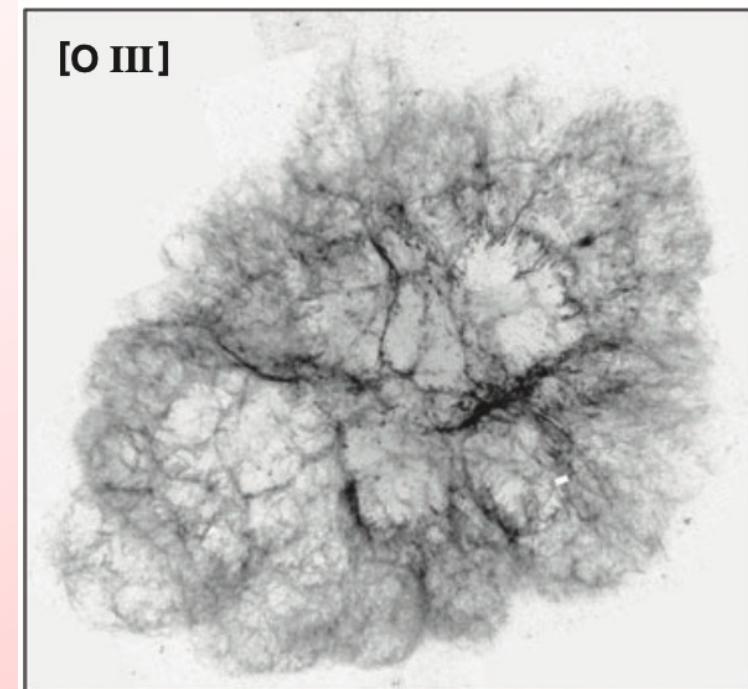
A Self-regulating Stochastic Acceleration Model of Pulsar Wind Nebulae

Shuta J. Tanaka
(Aoyama Gakuin Univ.)
with Wataru Ishizaki

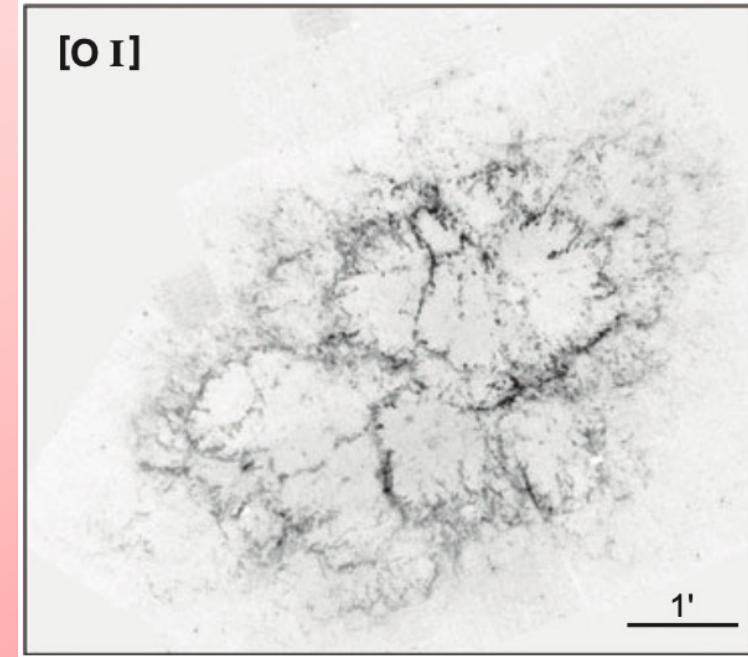
JWST (IR)



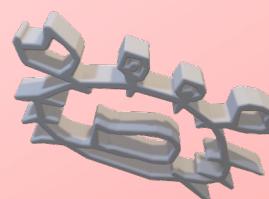
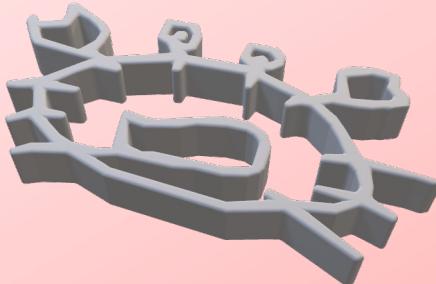
[O III]



[O I]



Introduction



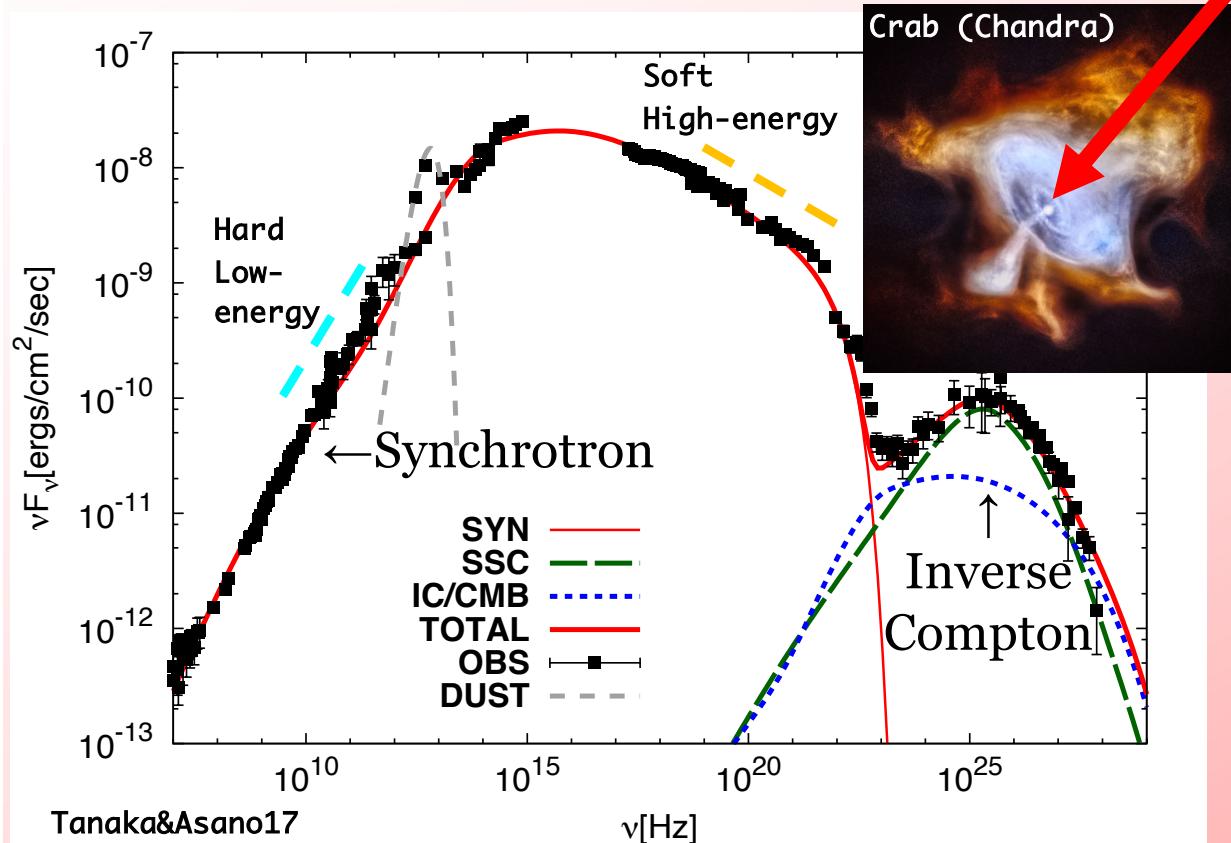
1'

2

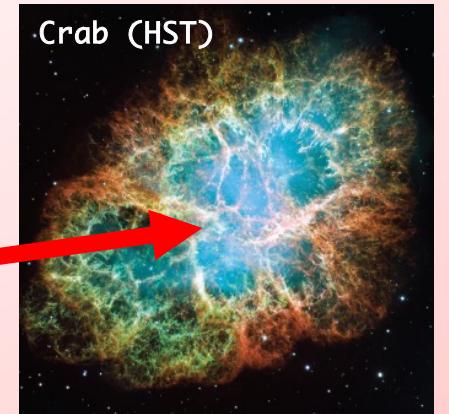
PWN as Particle Accelerator

pulsar wind nebula

Broadband spectrum from sub- μ eV through PeV



Crab Pulsar

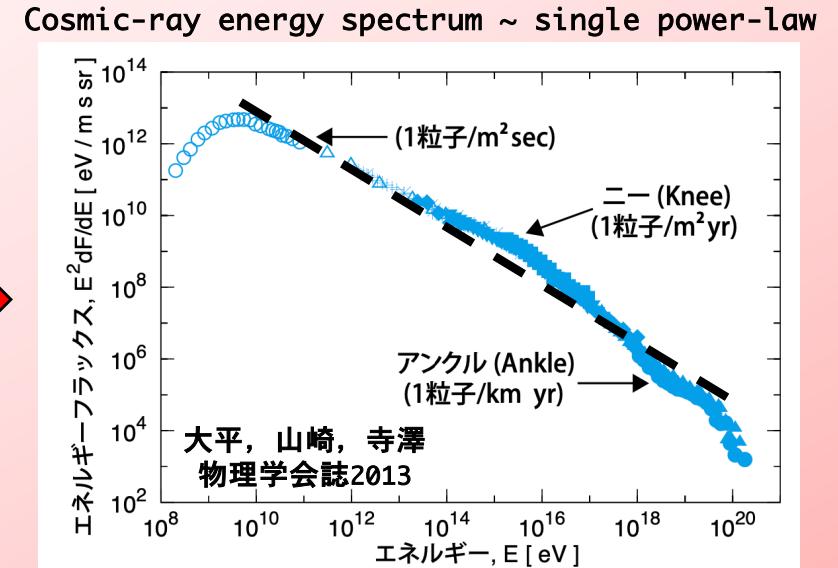
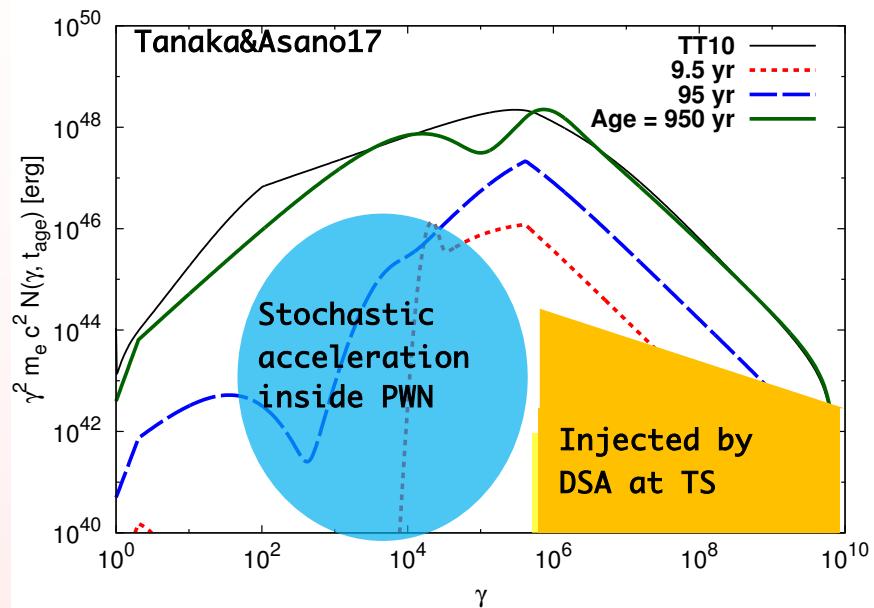


Crab (HST)

- PeV photons from the Crab
- LHAASO Collab. 21 Sci
- Particle accelerator > PeV
- Optical peak ← synchrotron from TeV particles.
- Relativistic ($T_{\text{PWN}} \sim \text{TeV}$) and rarefied ($n_{\text{PWN}} < 10^{-6} \text{ cm}^{-3}$) magnetized plasma cloud

Closest relativistic object

Origin of Radio Emission



- A single power-law particle distribution is predicted from particle acceleration theory.
- Low energy component has a harder spectrum than diffusive shock acceleration
- Low energy component is dominant in total particle number. (κ -problem)

our stochastic accel.
model

Tanaka&Asano17

Two different accel.
mechanisms.

+

Particles supplied from
outside PWNe

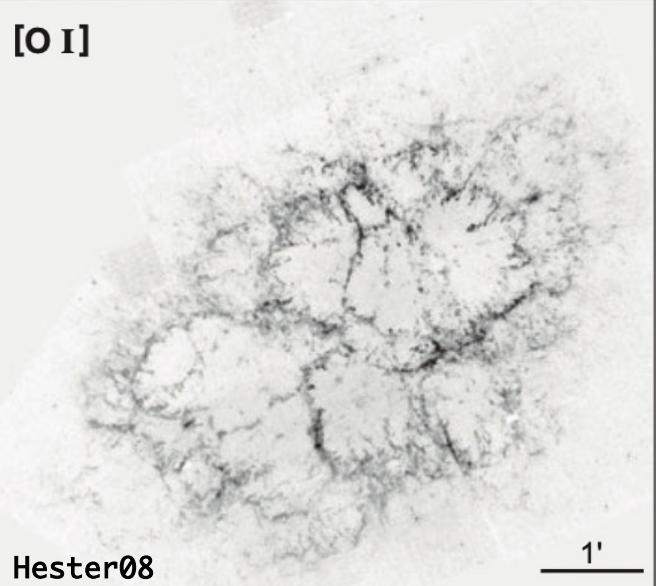
Injection from Ejecta Filaments

Table 4
Gas-phase Elemental Ion Fractions for the Best-fit Clumpy Model vi

Owen&Barlow15

[O III]

Species	Neutral	1 ⁺	2 ⁺	3 ⁺	4 ⁺	5 ⁺
Hydrogen	0.130	0.870	?	?	?	?
Helium	0.332	0.630	3.77×10^{-2}	?	?	?
Carbon	1.01×10^{-2}	0.730	0.248	2.08×10^{-2}	2.01×10^{-6}	1.04×10^{-10}
Nitrogen	1.04×10^{-2}	0.708	0.237	5.39×10^{-3}	1.17×10^{-6}	2.34×10^{-9}
Oxygen	0.144	0.721	0.107	2.75×10^{-3}	1.05×10^{-6}	7.37×10^{-8}
Neon	0.114	0.772	0.113	3.72×10^{-4}	4.10×10^{-6}	9.93×10^{-9}
Sulphur	0.198	0.440	0.299	7.05×10^{-3}	3.34×10^{-5}	5.66×10^{-8}
Argon	2.31×10^{-5}	0.116	0.702	0.178	2.31×10^{-3}	4.25×10^{-5}



Photoionization of neutrals in filaments
 $d \sim 0.1 \text{ pc}$
 $L \sim 3 \text{ pc}$
 $(n_{\text{fil}}, T_{\text{fil}}) \sim (10^3 \text{ cm}^{-3}, 10^3 \text{ K})$

$$\dot{N}_{\text{fil}} \lesssim \pi d l c_s n_{\text{fil}} \\ \approx 10^{46} \text{ s}^{-1}$$

Only a tiny ($\sim 10^{-5}$) fraction to be accelerated.

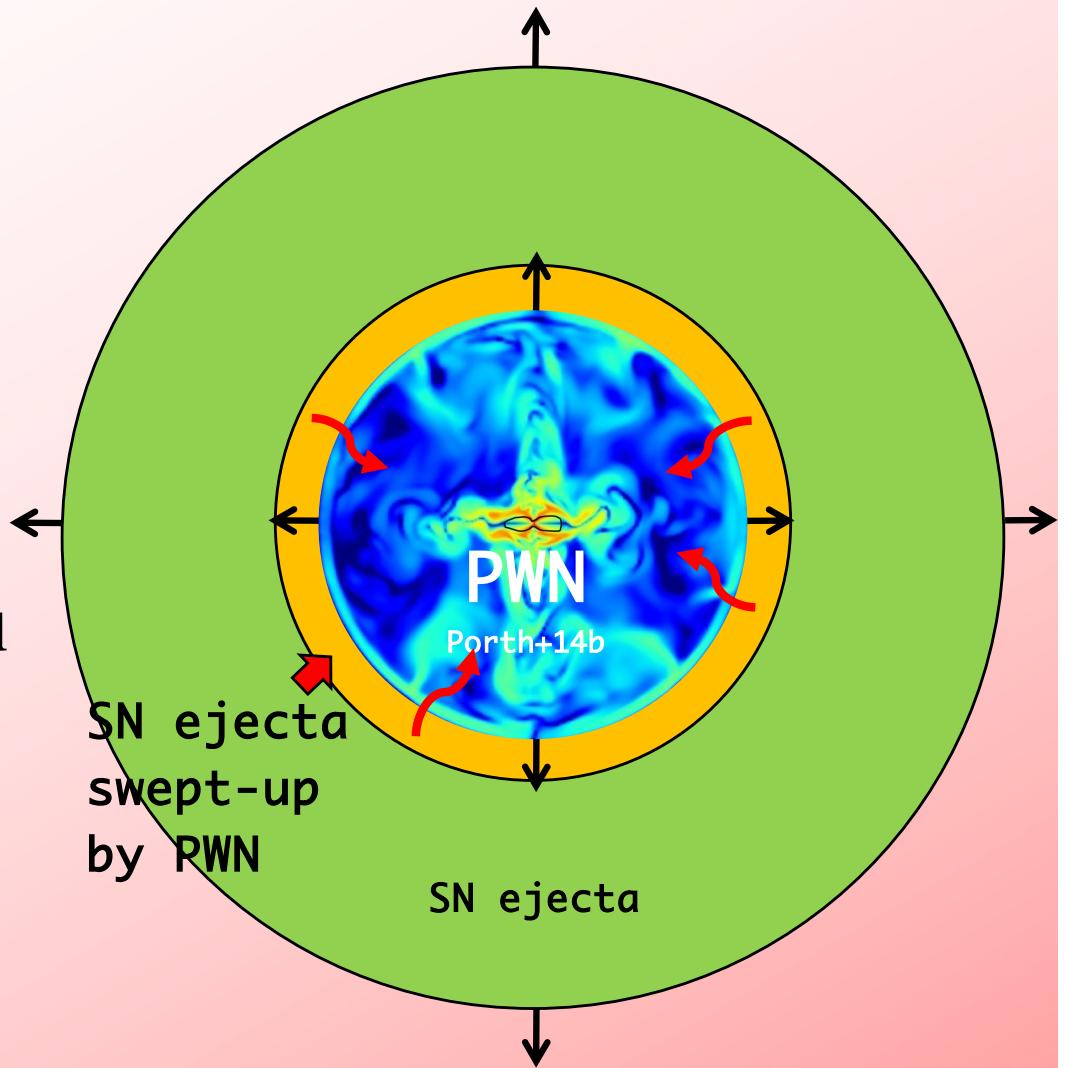
c.f. $\dot{N}_{\text{GJ}} \approx 10^{34} \text{ s}^{-1} L_{\text{spin,38}}^{1/2}$

Model

Energetics

One-zone approx. for PWN

- Uniform PWN expanding inside SN ejecta. e.g., Gelfand+09, Bandiera+20
- Rela. particles (e^\pm) and B-field supplied from PSR e.g., Pacini&Salvati73, Kennel&Coroniti84
- Non-rela. particles supplied from SN ejecta (photoionization) Tanaka&Asano17
- Turbulent energy for stochastic accel. is supplied from PSR and decreases by accelerating particles (backreaction).



$$\frac{4\pi}{3} R_{\text{PWN}}^3(t) \frac{B^2(t)}{8\pi} = \eta_B \int_0^t L_{\text{spin}}(t') dt'$$

$$L_{\text{spin}} = (\eta_e + \eta_B + \eta_{\text{turb}}) L_{\text{spin}}$$

Tanaka&Takahara10, Tanaka&Asano17, Tanaka&Kashiyama23

Particle Distribution

$$\frac{\partial}{\partial t} N(\gamma, t) + \frac{\partial}{\partial \gamma} \left[\underbrace{\left(\dot{\gamma}_{\text{cool}}(\gamma, t) - \gamma^2 D_{\gamma\gamma}(\gamma, t) \frac{\partial}{\partial \gamma} \frac{1}{\gamma^2} \right)}_{\text{cooling effects}} N(\gamma, t) \right] = \underbrace{Q_{\text{PSR}}(\gamma, t)}_{\text{from pulsar}} + \underbrace{Q_{\text{ext}}(t)}_{\text{Extra shock accel.}}$$

Tanaka&Ishizaki24

$$D_{\gamma\gamma} = \frac{\gamma^2}{2t_{\text{acc}}}, \quad t_{\text{acc}}(t) = \tau_{\text{acc}} \frac{\eta_T E_{\text{rot}}}{E_T(t)}$$

- τ_{acc} : initial acceleration time
- t_{acc} : accel. time increases with decaying of turbulence
- $E_T(t)$: energy of turbulence

$$\frac{dE_T}{dt} = \eta_T L_{\text{spin}} - \frac{E_T}{t_{\text{adi}}(t)} - \left(\frac{\delta E_T}{\delta t} \right)_{\text{damp}}$$

decay by expansion decay by particle accel.

$$Q_{\text{ext}}(\gamma, t) = f_{\text{inj}} 4\pi R_{\text{PWN}}^2(t) v_{\text{PWN}}(t) n_{\text{ej}}(R_{\text{PWN}}(t)) \delta(\gamma - \gamma_{\text{inj}})$$

- f_{inj} : injection efficiency
 $f_{\text{inj}} \ll 1 (O(10^{-5}))$
- γ_{inj} : injection energy
 $\gamma_{\text{inj}} \sim 1$

Injection of hadrons!!

Neutrino Emission

$$t_{pp} = \frac{1}{\xi_{pp} n \sigma_{pp} c} \approx 10^7 \text{yr} \left(\frac{n}{10 \text{cm}^{-3}} \right)^{-1} \left(\frac{\xi_{pp}}{0.2} \right)^{-1}$$

target hadrons inelasticity

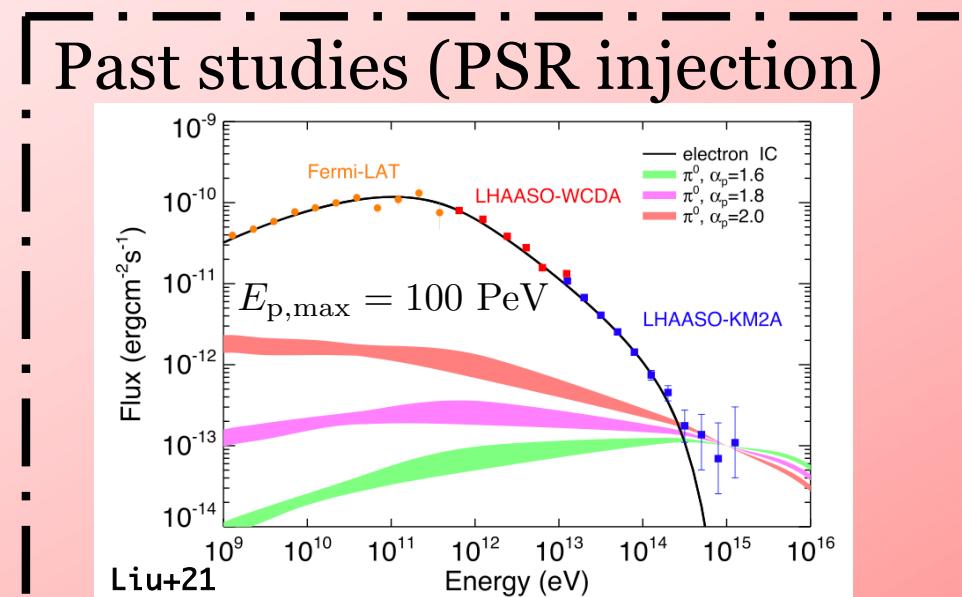
$$n_{\text{ej}} = 0.74 \text{ cm}^{-3} \left(\frac{M_{\text{ej}}}{9.5 M_{\odot}} \right) \left(\frac{R_{\text{SNR}}}{5 \text{pc}} \right)^{-3}$$

ejecta of SN 1054 condensate
as filaments inside the Crab.

Owen&Barlow15

$$n_{\text{ej}} = 17.2 \text{ cm}^{-3} \left(\frac{M_{\text{ej}}}{7.2 M_{\odot}} \right) \left(\frac{R_{\text{SNR}}}{1.6 \text{pc}} \right)^{-3}$$

$$\epsilon_{\nu}^2 \frac{dN_{\nu}}{dt} \approx \gamma^2 m_e c^2 \frac{dN_e}{d\gamma} \frac{1}{t_{pp}} \Big|_{\epsilon_{\nu}=0.1 \gamma m_e c^2}$$



$$\frac{dN_p}{dE_p} = N_0 E_p^{\alpha_p} \exp \left(-\frac{E_p}{E_{p,\max}} \right)$$

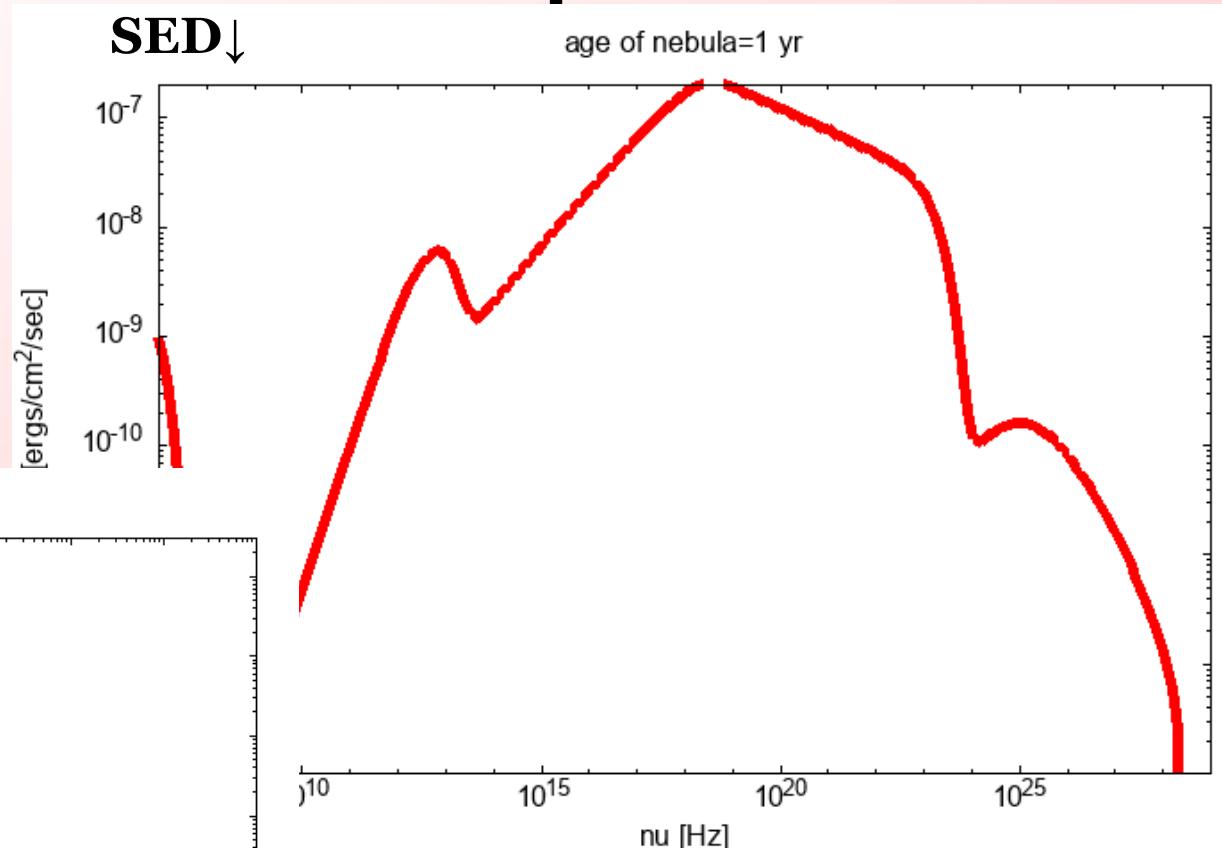
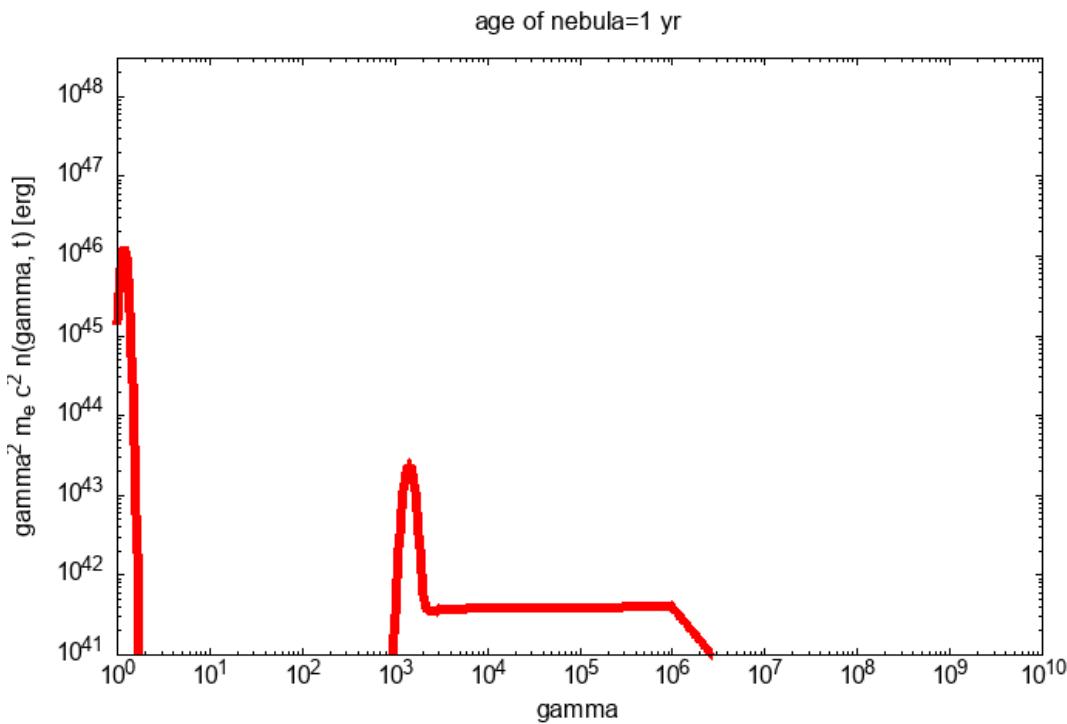
$\int E_p \frac{dN_p}{dE_p} dE_p = \eta_p \int dt L(t)$
 $\eta_p \sim 20 \%$

Results

Results: Broadband Spectrum

$$\tau_{\text{acc},0} = 10 \text{ yr}$$

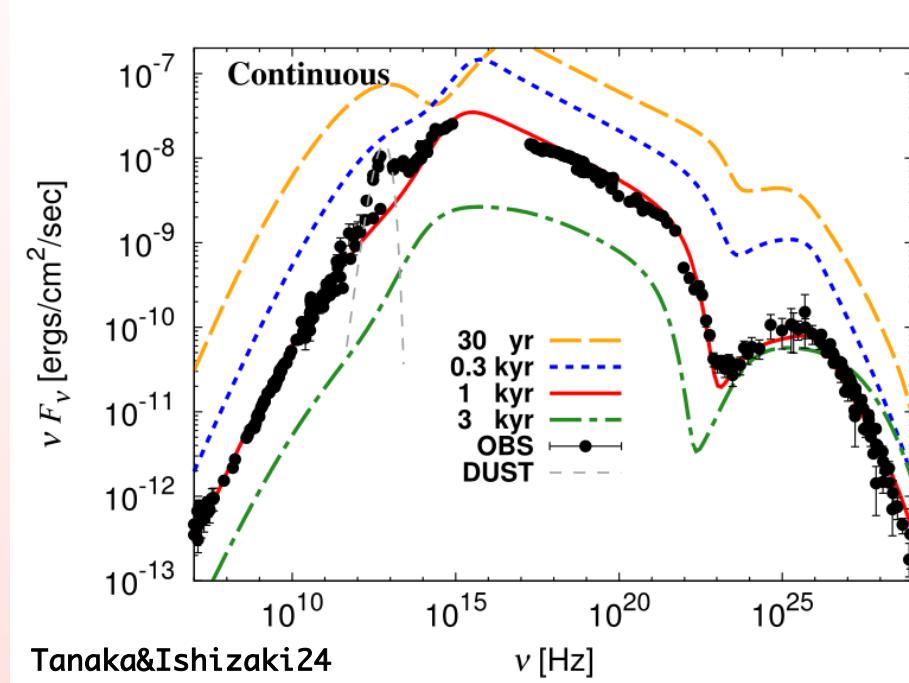
Electron Energy Spectrum ↓



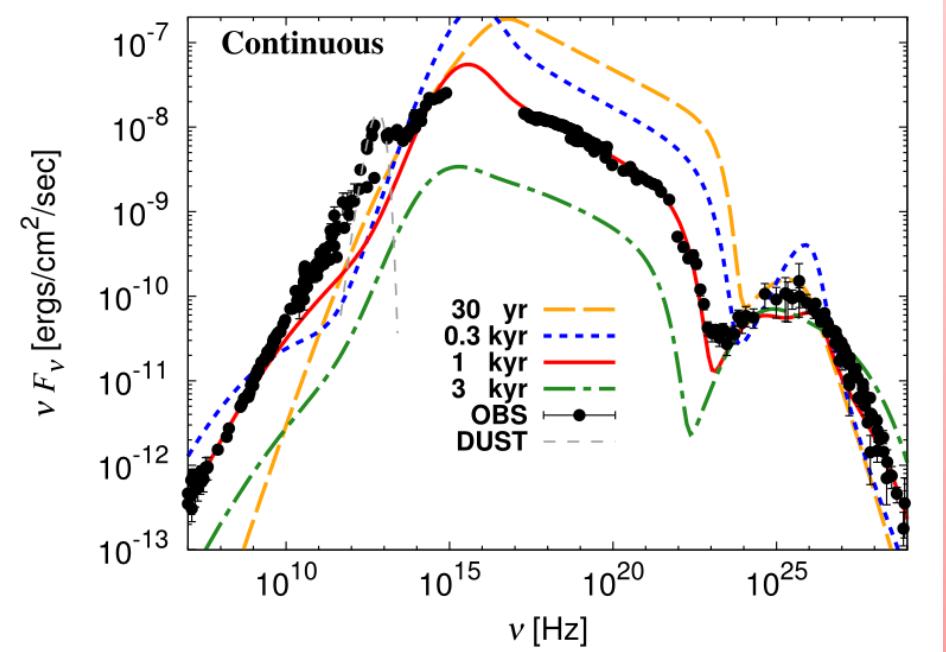
$$\frac{\partial}{\partial t} N(\gamma, t) + \frac{\partial}{\partial \gamma} \left[\left(\dot{\gamma}_{\text{cool}}(\gamma, t) - \gamma^2 D_{\gamma\gamma}(\gamma, t) \frac{\partial}{\partial \gamma} \frac{1}{\gamma^2} \right) N(\gamma, t) \right] = Q_{\text{PSR}}(\gamma, t) + Q_{\text{ext}}(t)$$

Results

$$\tau_{\text{acc},0} = 3 \text{ yr}$$



$$\tau_{\text{acc},0} = 30 \text{ yr}$$

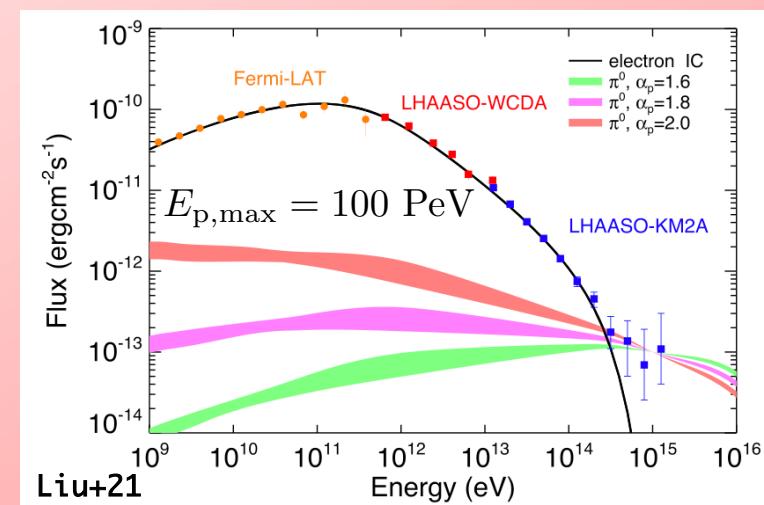
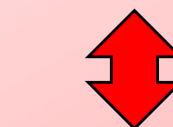
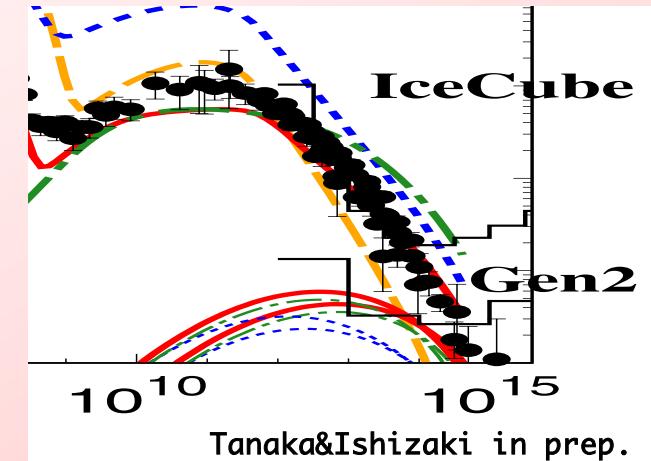


Acceleration time-scale should be < 10 yr

Conclusions

**Stochastic accel. model
including backreaction
to turbulent energy**

- **A physical model of the origin of radio-emitting particles (could be a solution to κ -problem).**
- **Neutrino signal which is different from previous studies is expected from our model!**



**JSPJ Grant-in-Aid 24Ho1816, 23K20038
Sumitomo foundation 210629**