

Measurement of the depth of the maximum of air-shower profiles above $10^{17.7}$ eV at the Pierre Auger Observatory

Nikolas Denner on behalf of the Pierre Auger Collaboration

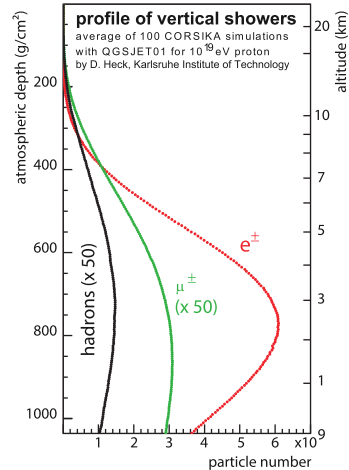
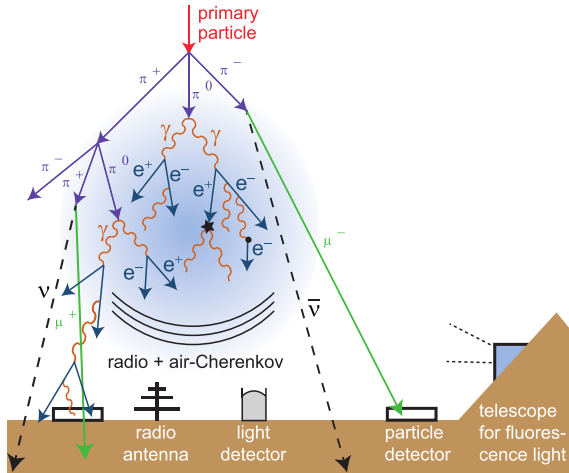
Institute of Physics of the Czech Academy of Sciences, Prague

3 September 2026

Longitudinal profile of cosmic ray showers

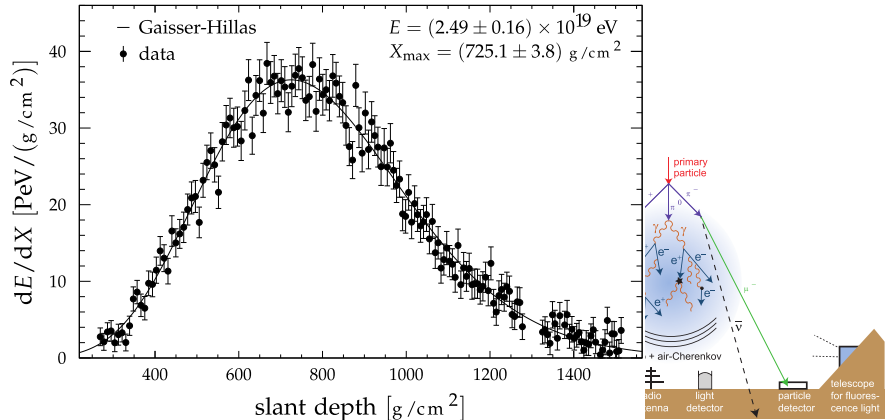
- Cosmic-ray showers are caused by impact of primary cosmic rays with our atmosphere

[F. G. Schröder, Prog. Part. Nucl. Phys. 93 (2017) 1–68]



Longitudinal profile of cosmic ray showers

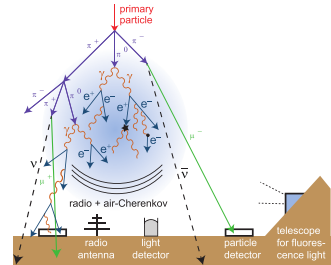
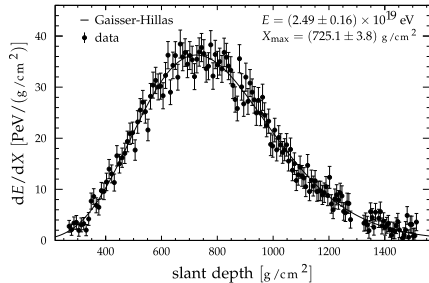
- Cosmic-ray showers are caused by impact of primary cosmic rays with our atmosphere
- Longitudinal profile = energy deposited by cosmic-ray shower as a function of slant depth



Longitudinal profile of cosmic ray showers

- Cosmic-ray showers are caused by impact of primary cosmic rays with our atmosphere
- Longitudinal profile = energy deposited by cosmic-ray shower as a function of slant depth
- Slant depth X = mass per unit area traversed by a particle / cosmic-ray shower

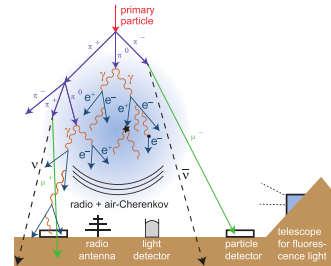
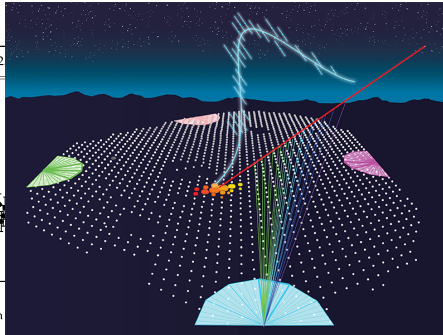
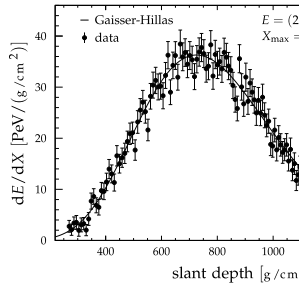
$$X(z) = \int_z^\infty \rho(r) dr$$



Longitudinal profile of cosmic ray showers

- Cosmic-ray showers are caused by impact of primary cosmic rays with our atmosphere
- Longitudinal profile = energy deposited by cosmic-ray shower as a function of slant depth
- Slant depth X = mass per unit area traversed by a particle / cosmic-ray shower
- Detectable through UV light emitted isotropically by excited nitrogen atoms

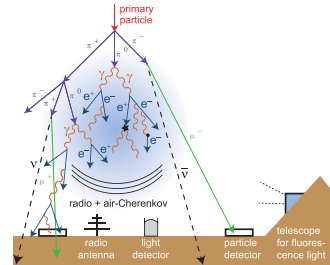
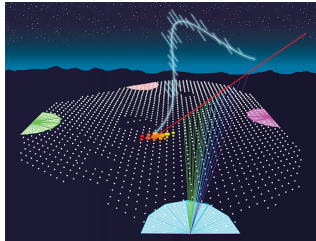
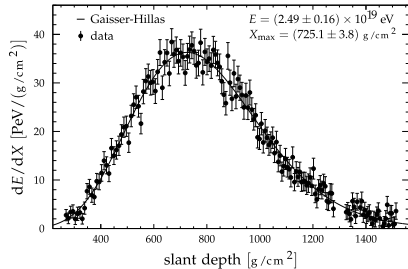
$$X(z) = \int_z^{\infty} \rho(r) dr$$



Longitudinal profile of cosmic ray showers

- Cosmic-ray showers are caused by impact of primary cosmic rays with our atmosphere
- Longitudinal profile = energy deposited by cosmic-ray shower as a function of slant depth
- Slant depth X = mass per unit area traversed by a particle / cosmic-ray shower
- Detectable through UV light emitted isotropically by excited nitrogen atoms
- Depth of the shower maximum $X_{\max}$ **correlated with** mass number A of primary particles

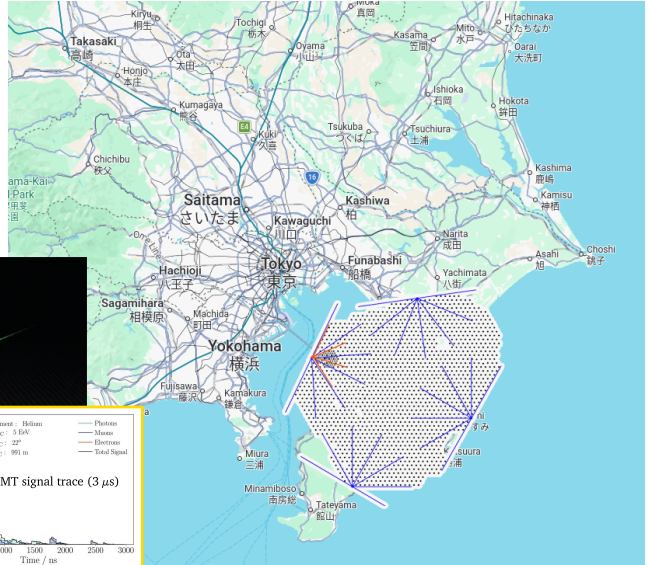
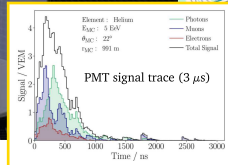
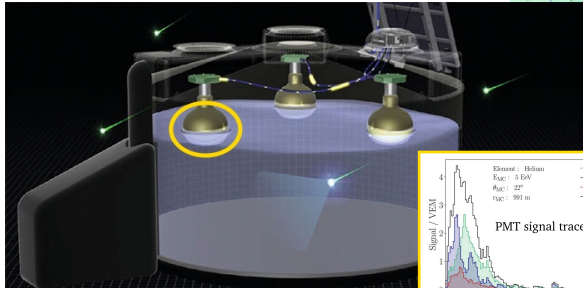
$$X(z) = \int_z^{\infty} \rho(r) dr$$



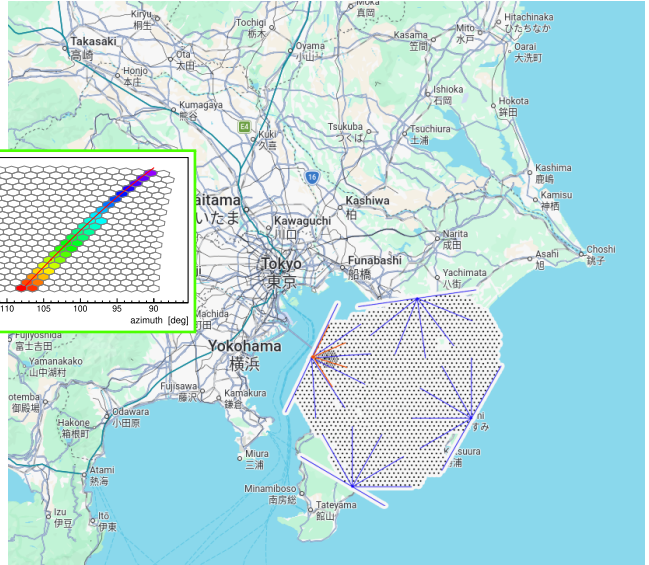
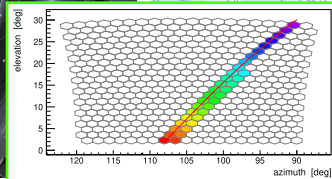
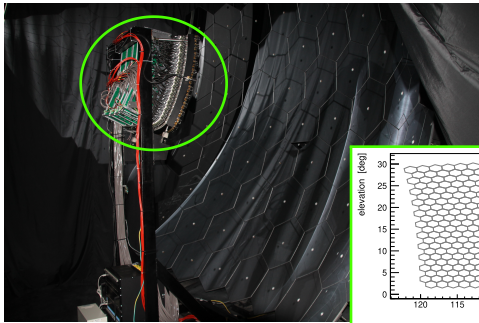
Hybrid Detector of the Pierre Auger Observatory

Surface detector (SD)

- Over 1600 water-Cherenkov stations
- 3 arrays: 1500 m, 750 m, 433 m
- Detecting **shower front** on the surface
- $\sim 100\%$ duty cycle



Hybrid Detector of the Pierre Auger Observatory

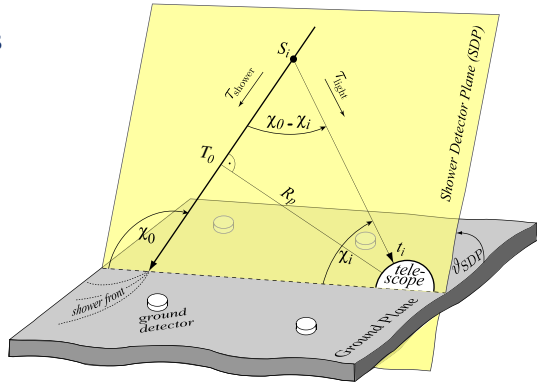


Fluorescence detector (FD)

- 24+3 fluorescence telescopes
- Overlooking SD array from 4 locations
- Detecting **longitudinal profile**
- $\sim 15\%$ duty cycle

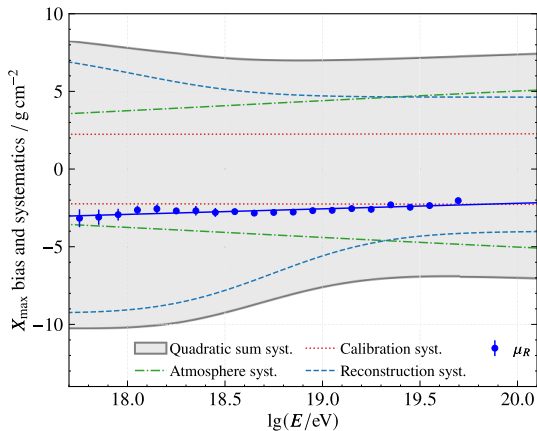
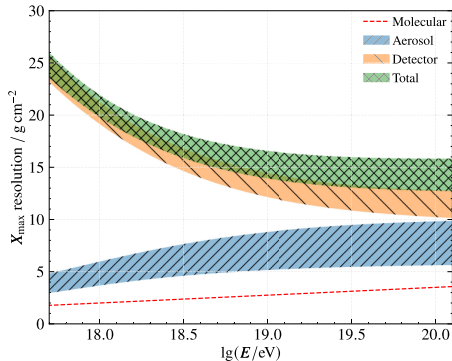
Latest measurement of $X_{\max}$ [Phys. Rev. D 114 (2026) 043016]

- Analysis of hybrid events **above $10^{17.7}$ eV from 2004 to the end of 2021**
- Hybrid events = events detected by the FD and SD-1500 together
- Shower geometry estimated from the FD and SD-station closest to the shower axis



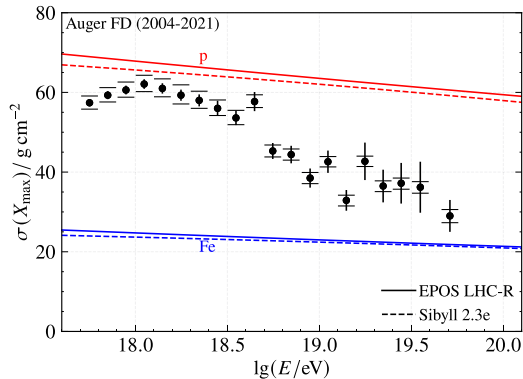
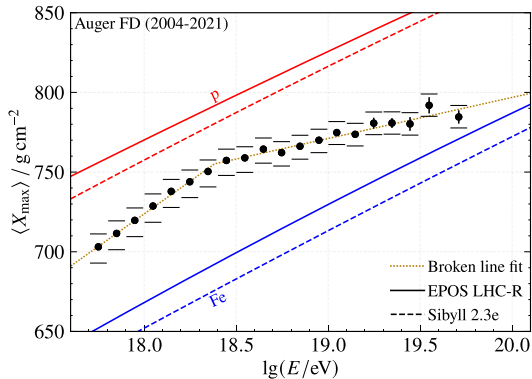
True $X_{\max}$ distributions and systematics

- To get the true $X_{\max}$ distributions - take into account detector response and acceptance
- Sibyll 2.3d simulations used to obtain $X_{\max}$ bias, resolution and acceptance
- Scale uncertainty - effects of calibration, reconstruction and atmosphere



Mean and standard deviation of $X_{\max}$

- Elongation rate $D_{10} = \frac{d\langle X_{\max} \rangle}{d \lg E}$ breaks at $\lg(E_{\text{break}}/\text{eV}) = 18.38 \pm 0.02(\text{stat})_{-0.01}^{+0.01}(\text{sys})$
 $\Rightarrow$ energy dependence of $\langle X_{\max} \rangle$ leans towards heavier elements
- Decreasing $\sigma(X_{\max}) \Rightarrow$ transition to predominantly heavy primaries at ultra-high energies



Transition from $X_{\max}$ to $\ln A$

- Mean and standard deviation of $X_{\max}$ are related to the statistical moments of $\ln A$ distributions
- Recalculation (interpretation) depends on hadronic-interaction models

$$\langle \ln(A) \rangle = \frac{\langle X_{\max} \rangle - \langle X_{\max} \rangle_p}{f_E}$$

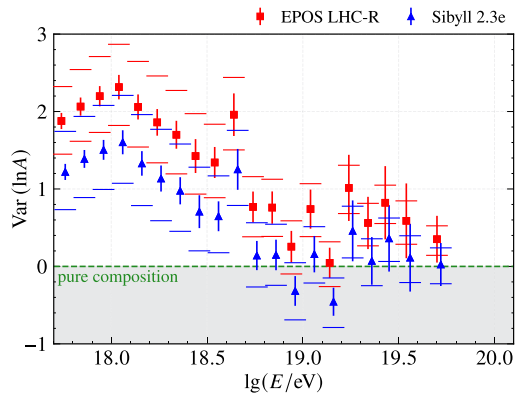
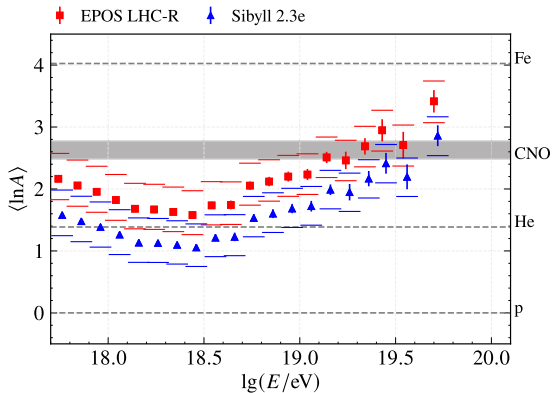
$$V(\ln(A)) = \frac{\sigma^2(X_{\max}) - \sigma_{\text{sh}}^2(\langle \ln(A) \rangle)}{b \sigma_p^2 - f_E^2}$$

$$f_E = \frac{\langle X_{\max} \rangle_{\text{Fe}} - \langle X_{\max} \rangle_p}{\ln(56)}$$

[JCAP 02 (2013) 026]

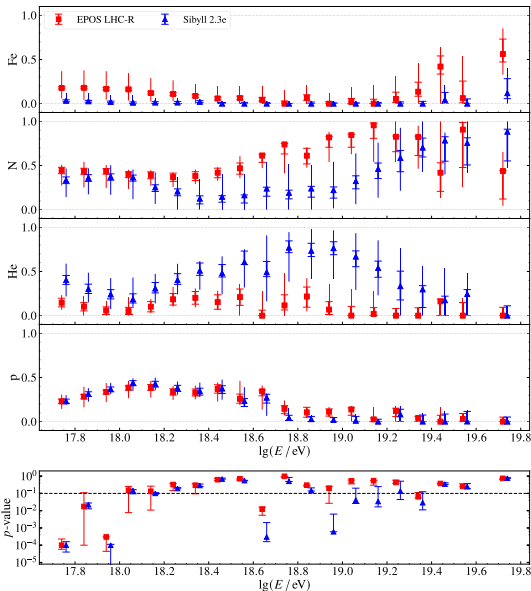
Mean and variance of $\ln A$

- Moments of $\ln A$ calculated from $X_{\max}$ using Sibyll 2.3e and EPOS LHC-R simulations
- About 0.5 g/cm^2 difference in $\langle \ln A \rangle$ between the two models



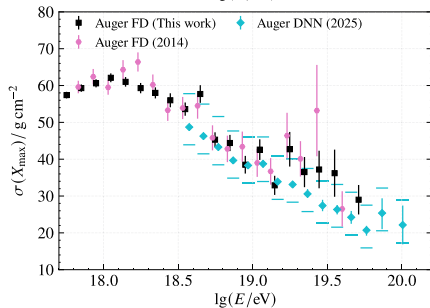
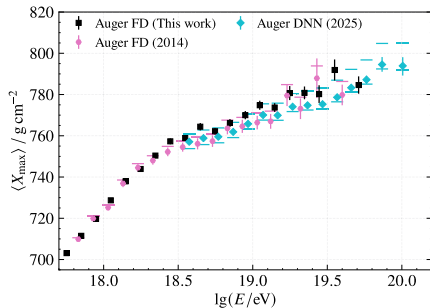
Primary particle fractions

- Fitting $X_{\max}$ distributions from data with $X_{\max}^{\text{MC}}$ distributions of four different nuclei
- Relative fractions used as fit parameter
- Sibyll 2.3e:
 - He dominant from $10^{18.5}$ eV to 10^{19} eV
 - after that, N dominates until the end
 - Fe consistent with zero at all energies
- EPOS LHC-R:
 - N dominant until Fe becomes significant
 - End of spectrum $\rightarrow \sim 50/50$ N/Fe



Comparison with previous results

- Auger DNN $\rightarrow X_{\max}$ from SD using deep learning
- FD 2014 $\rightarrow$ data from 2004 to the end of 2012
- Offset between new and old analyses caused by different treatment of aerosols
- Ongoing joint study comparing Auger and TA $X_{\max}$



Summary

Data-driven conclusions:

- Break in $X_{\max}$ as a function of energy $\Rightarrow$ **change of composition towards heavier nuclei**
- $X_{\max}$ distributions narrowing with **increasing energy** $\Rightarrow$ heavier and **less-mixed composition**

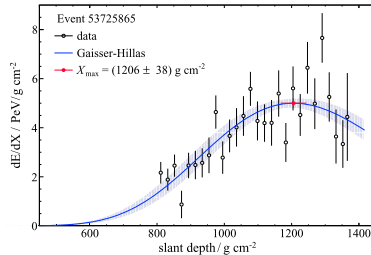
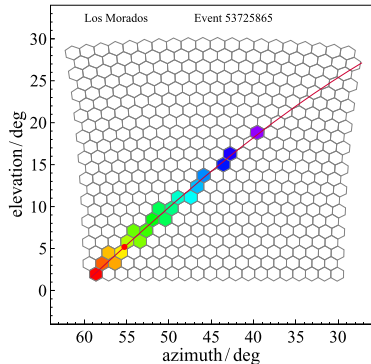
Model-dependent conclusions:

- Average mass around **CNO group at the highest energies**
- Mass fractions $\Rightarrow$ **p disappearing** above $\sim 10^{18.7}$ eV, composition mostly dominated by N
- **Data** for self-contained studies of mass composition and hadronic interactions
available on [zenodo](#)

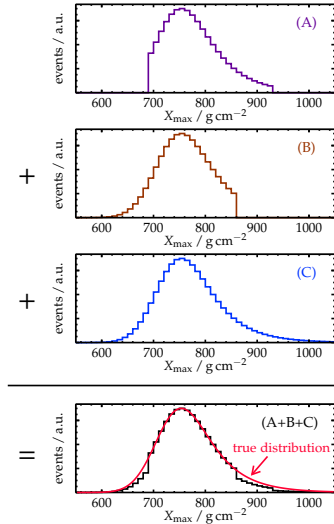
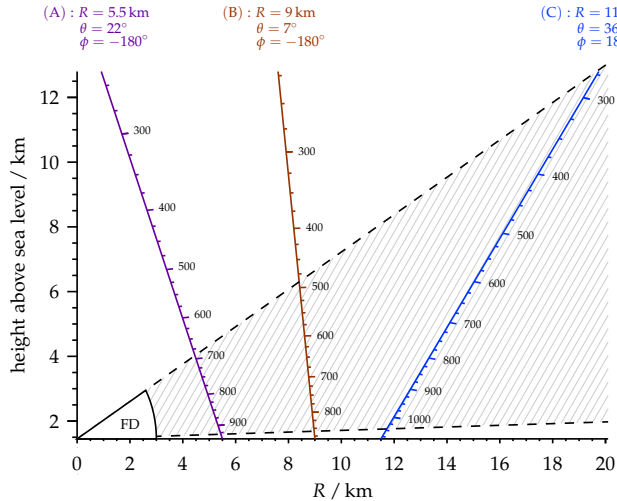
BACKUP SLIDES

Deepest event in the $X_{\max}$ dataset

- Event from June 2019, $X_{\max} = 1206 \pm 38 \text{ g/cm}^2$ in the FD site passing all selection cuts (LM)
- $E_{\text{tot}}^{\text{LM}} = 4.2 \pm 0.6 \text{ EeV}$
- Compatible with proton-initiated shower
- Heavier nuclei disfavored



$X_{\max}$ bias caused by the geometrical FOV of the telescopes



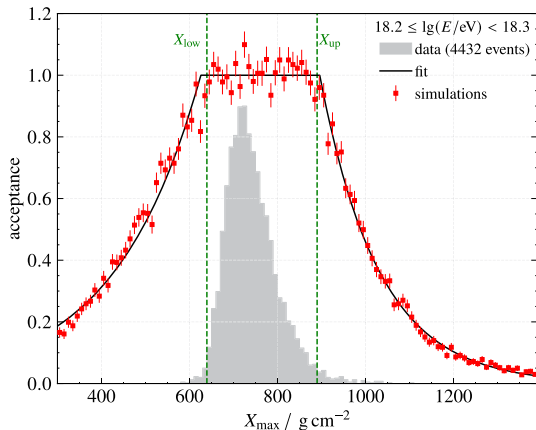
- Bias removed by the fiducial FOV cut - more details about this cut in the paper

$X_{\max}$ acceptance

- $X_{\max}$ acceptance ε as a function of $X_{\max}$
- Proton showers with Sibyll 2.3d across a wide range of $X_{\max}$
- Studying how many showers of given $X_{\max}$ are recorded and pass all the cuts

$$\varepsilon(X_{\max}) = \begin{cases} e^{+(X_{\max}-X_1)/\lambda_1}, & X_{\max} < X_1; \\ 1, & X_1 \leq X_{\max} \leq X_2; \\ e^{-(X_{\max}-X_2)/\lambda_2}, & X_2 < X_{\max}. \end{cases}$$

- Vertical lines = fiducial FOV range



$X_{\max}$ bias and resolution

- Distribution of the difference between reconstructed and generated $X_{\max}$ parametrized as sum of two Gaussian distributions

$$R(X_{\max}^{\text{rec}} - X_{\max}^{\text{gen}}) = fG(\mu, \sigma_1) + (1 - f)G(\mu, \sigma_2)$$

- μ is the $X_{\max}$ reconstruction bias
- Detector resolution of $X_{\max}$ is obtained as
$$\sigma(X_{\max}^{\text{rec}} - X_{\max}^{\text{gen}}) = \sqrt{f\sigma_1^2 + (1 - f)\sigma_2^2}$$
- Other contributions come from the alignment of the telescopes, from the aerosols and from the uncertainty in atmospheric density profile

