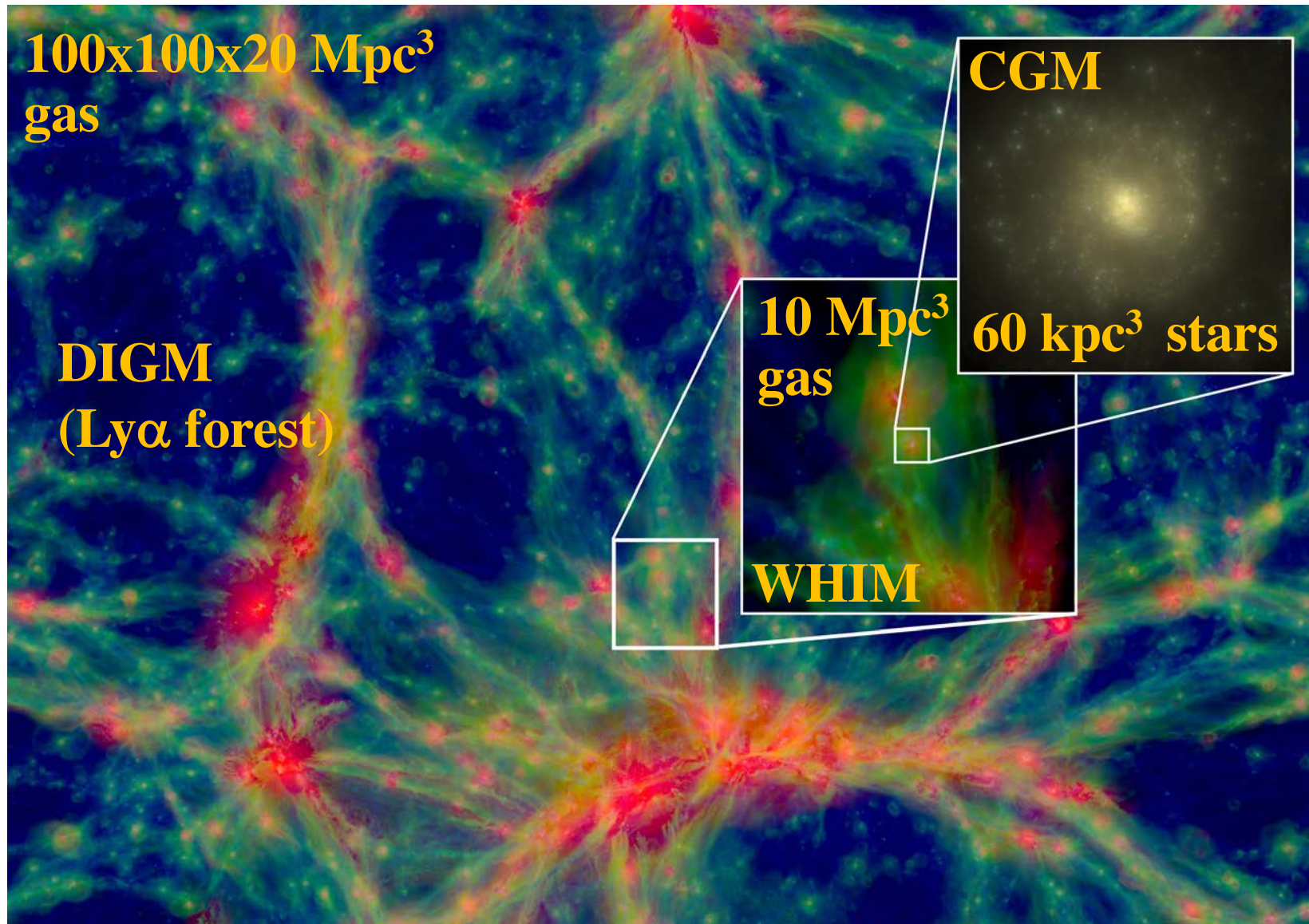


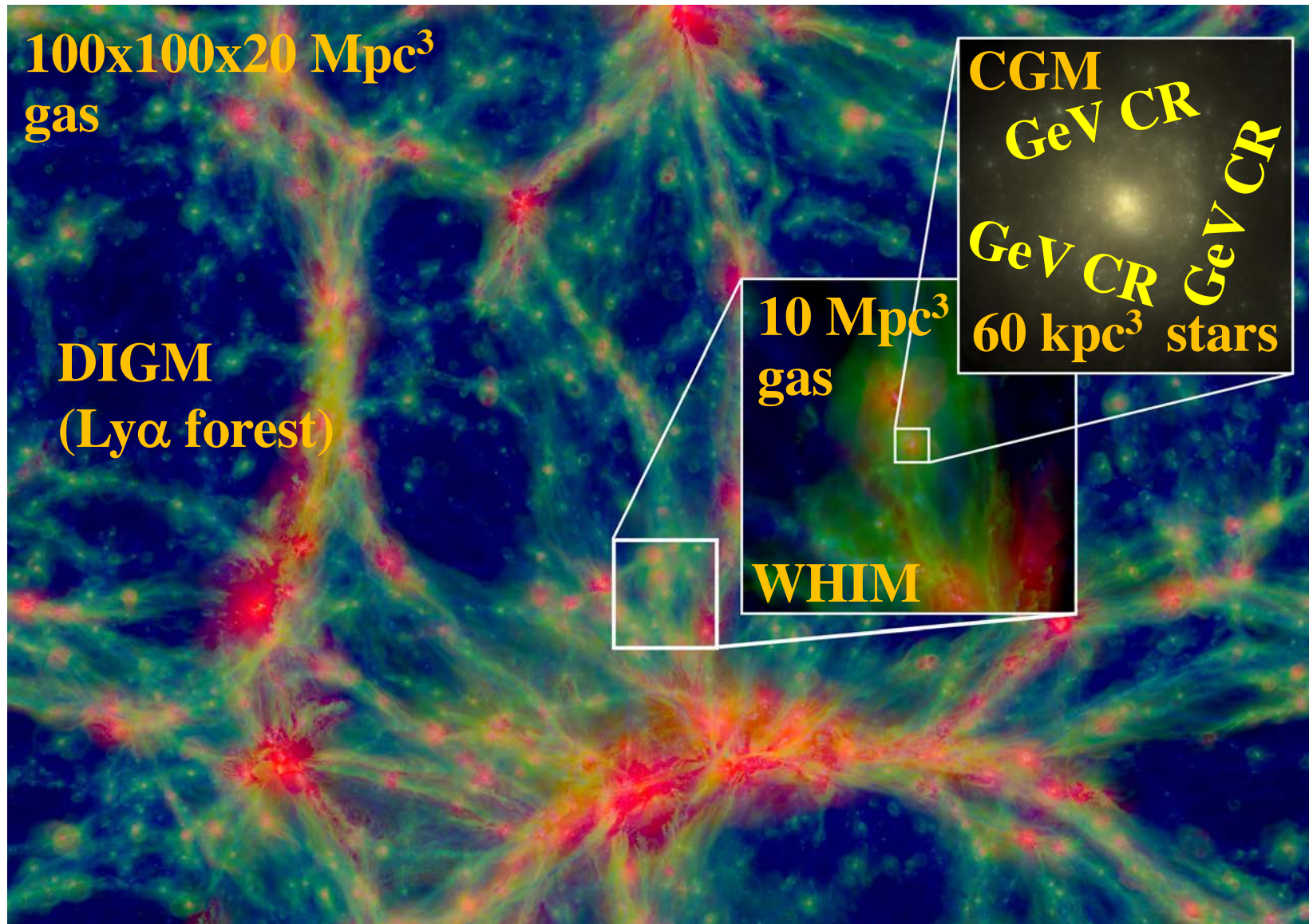
The role of TeV-PeV cosmic rays in the Milky Way and beyond

Susumu Inoue (Chiba)



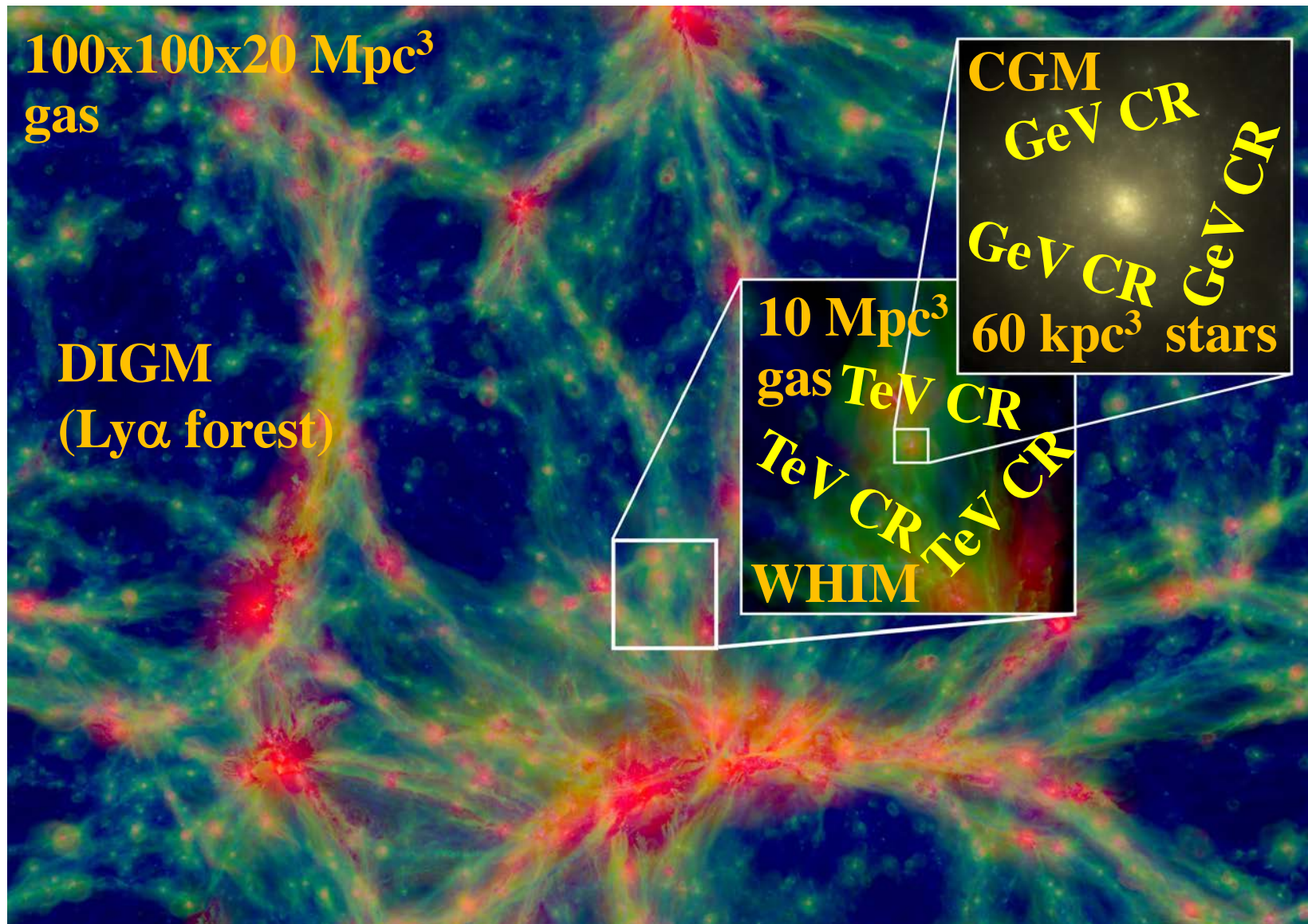
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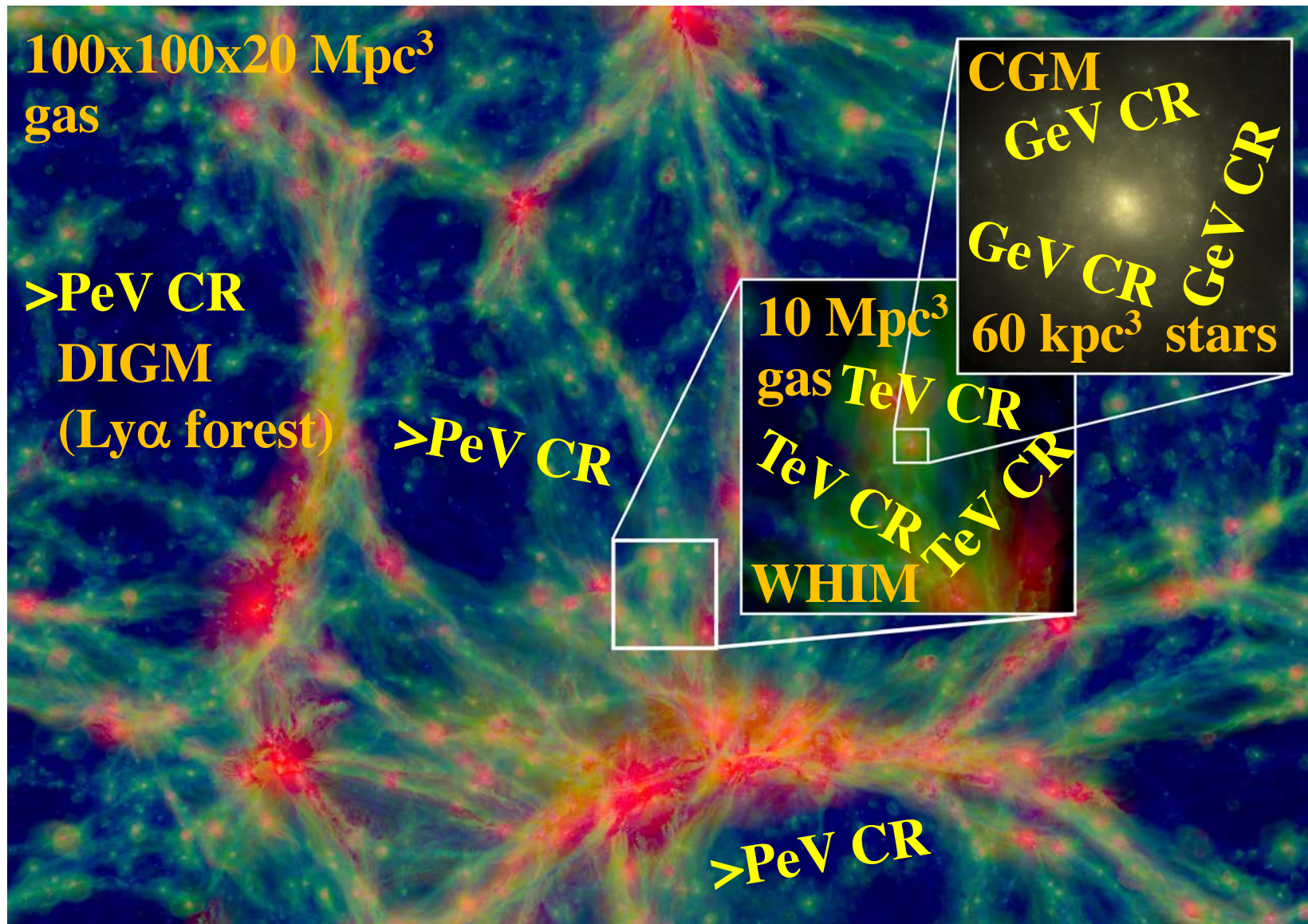
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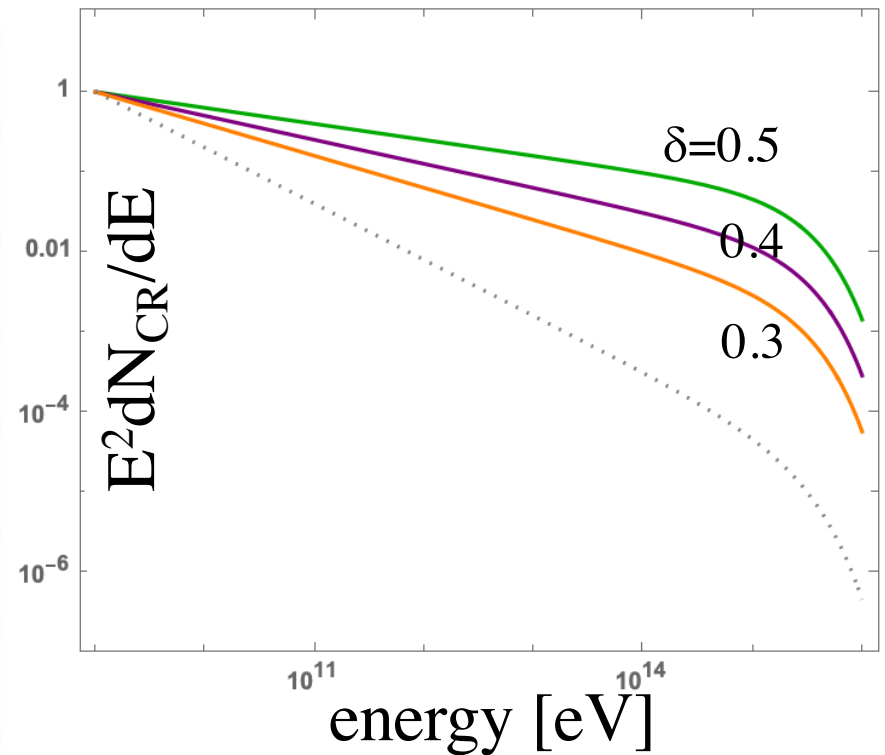
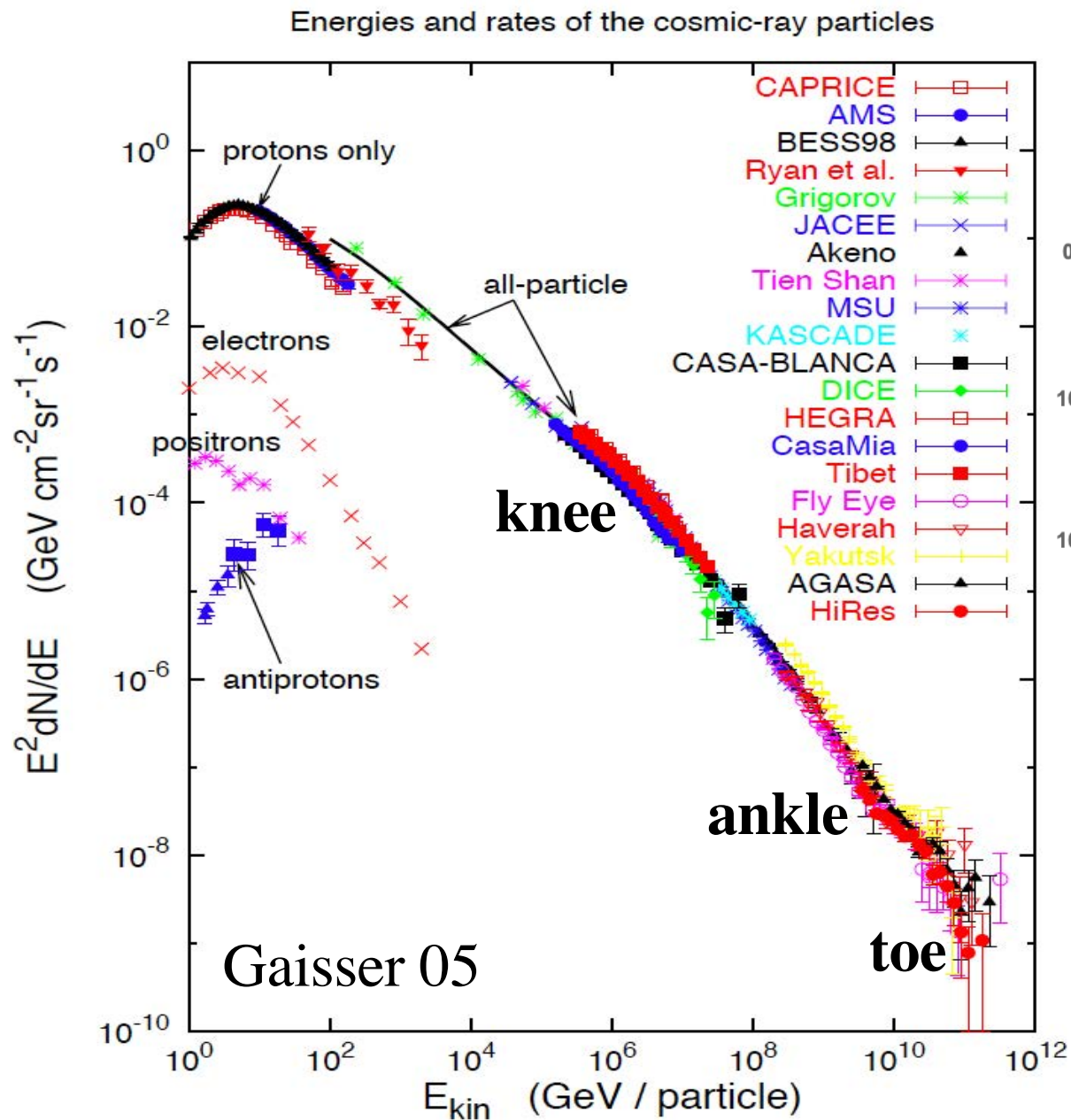
The role of TeV-PeV cosmic rays in the Milky Way and beyond

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EAGLE
Schaye+ 14

observed CR spectrum vs source energetics



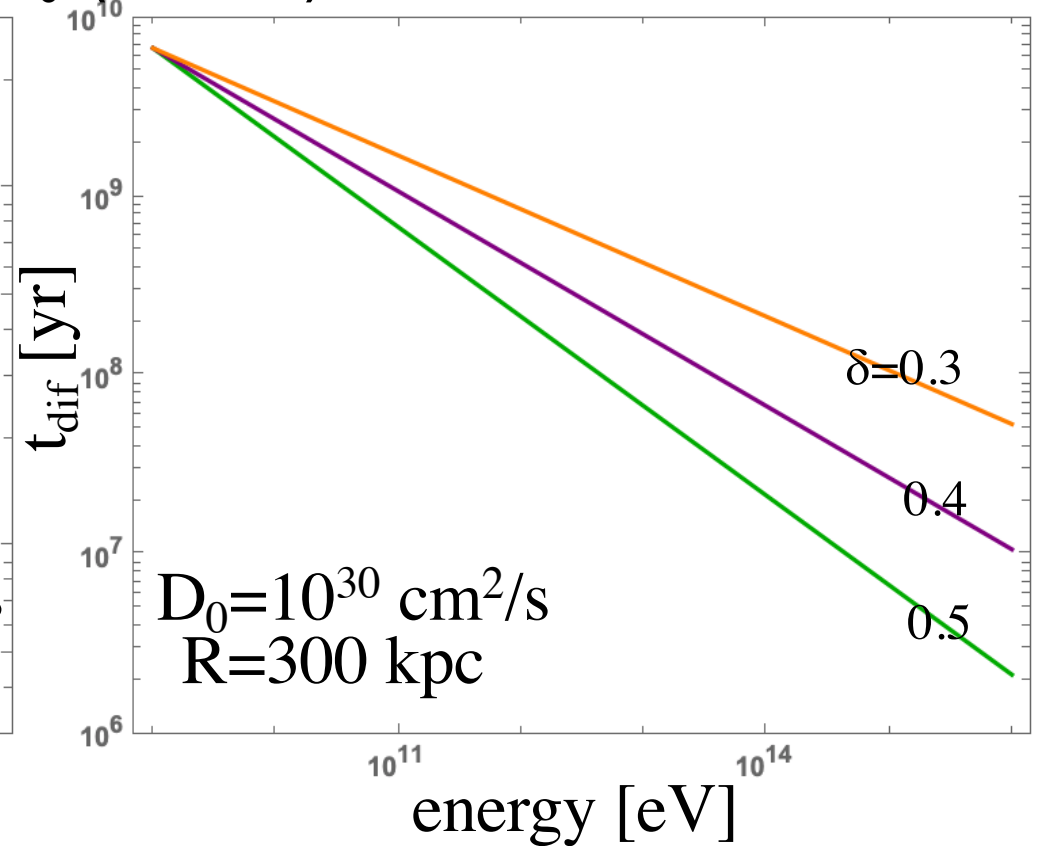
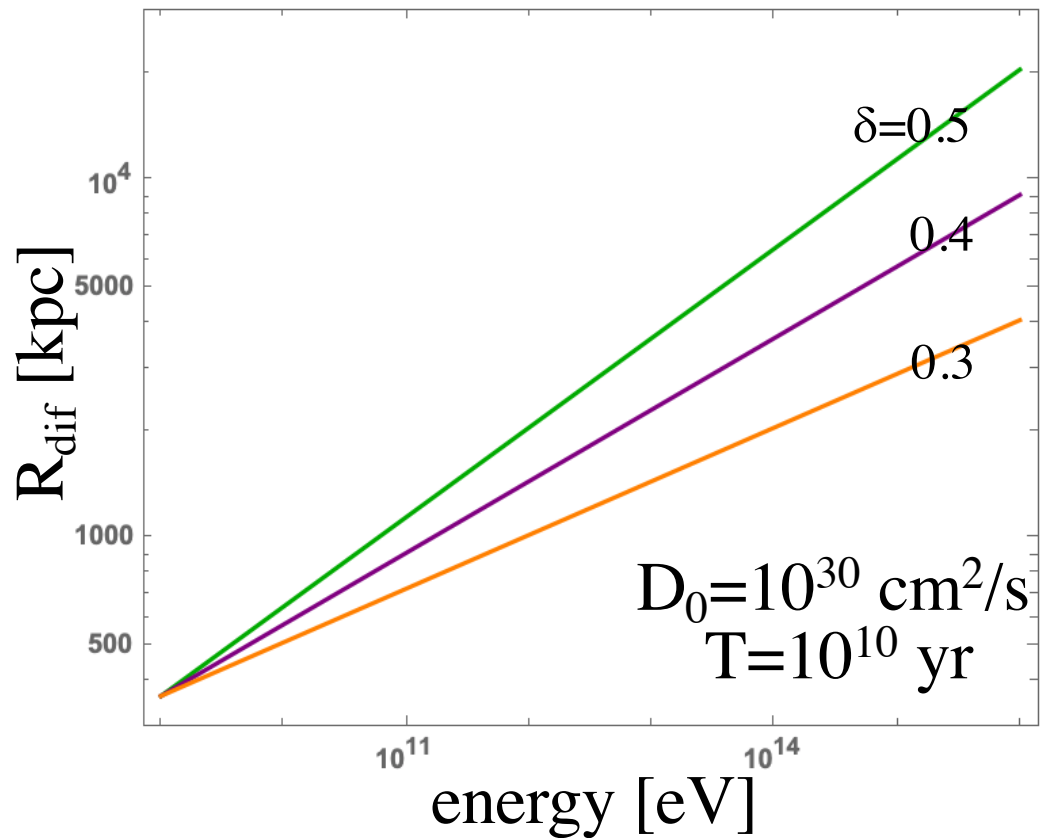
$D(E) = D_0 (E / \text{GeV})^\delta$
 spectrum in ISM
 source spectrum +
 steepening by E-dep.
 confinement

diffusion radius, diffusion time

$$R_{\text{dif}}(E) = (4 D(E) t)^{1/2}$$

$$t_{\text{dif}}(E) \sim \frac{R^2}{4 D(E)}$$

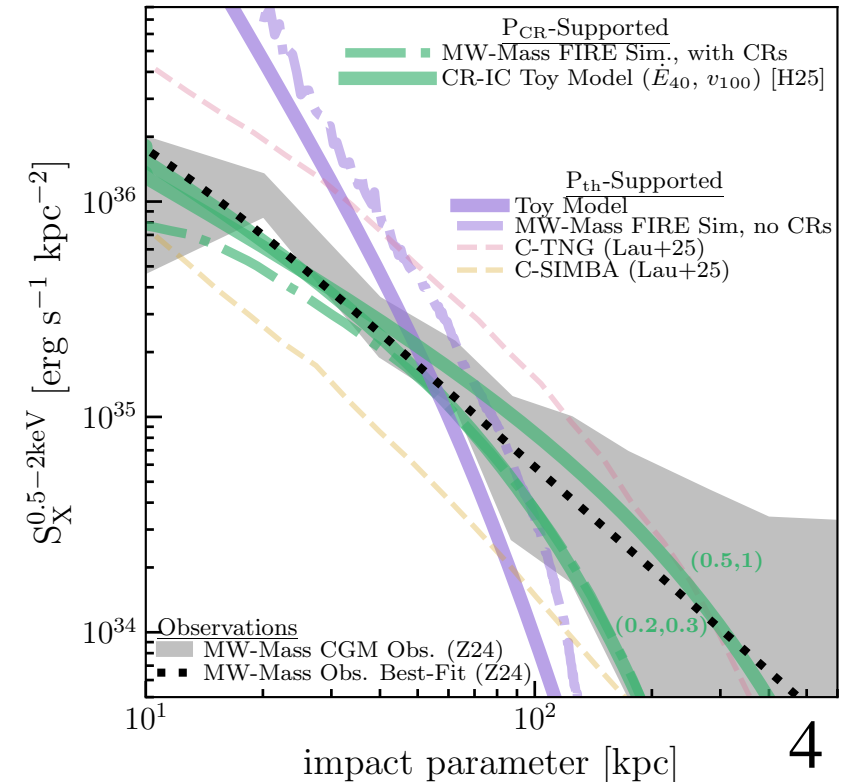
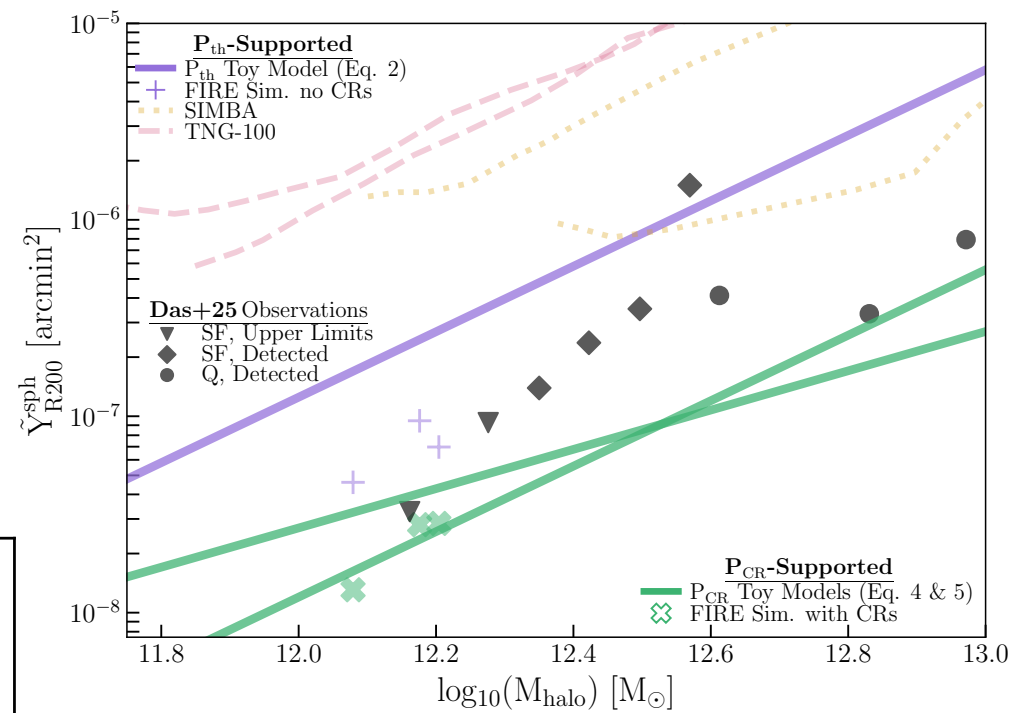
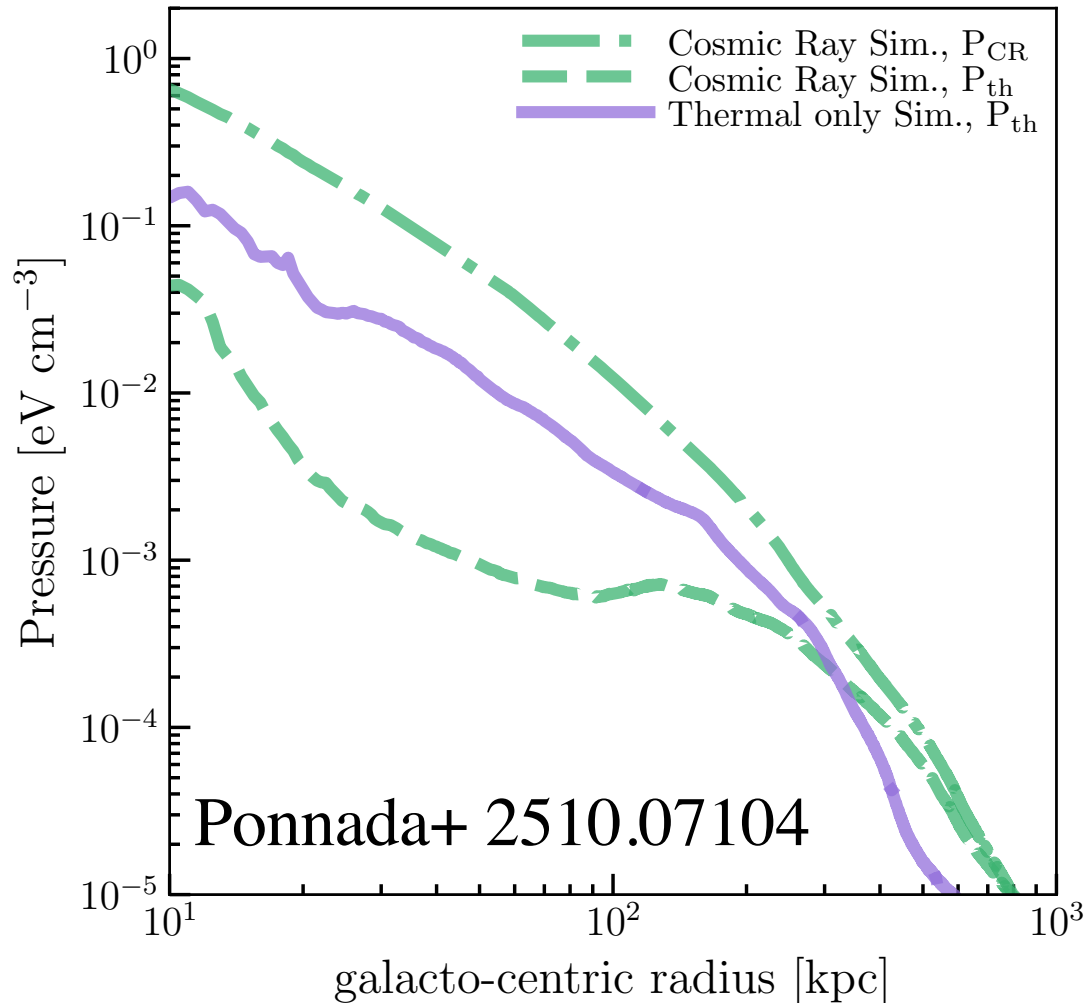
$$D(E) = D_0 (E / \text{GeV})^\delta$$



consider protons only
for nuclei $E \rightarrow E/Z$

CR-dominated halos: possible evidence

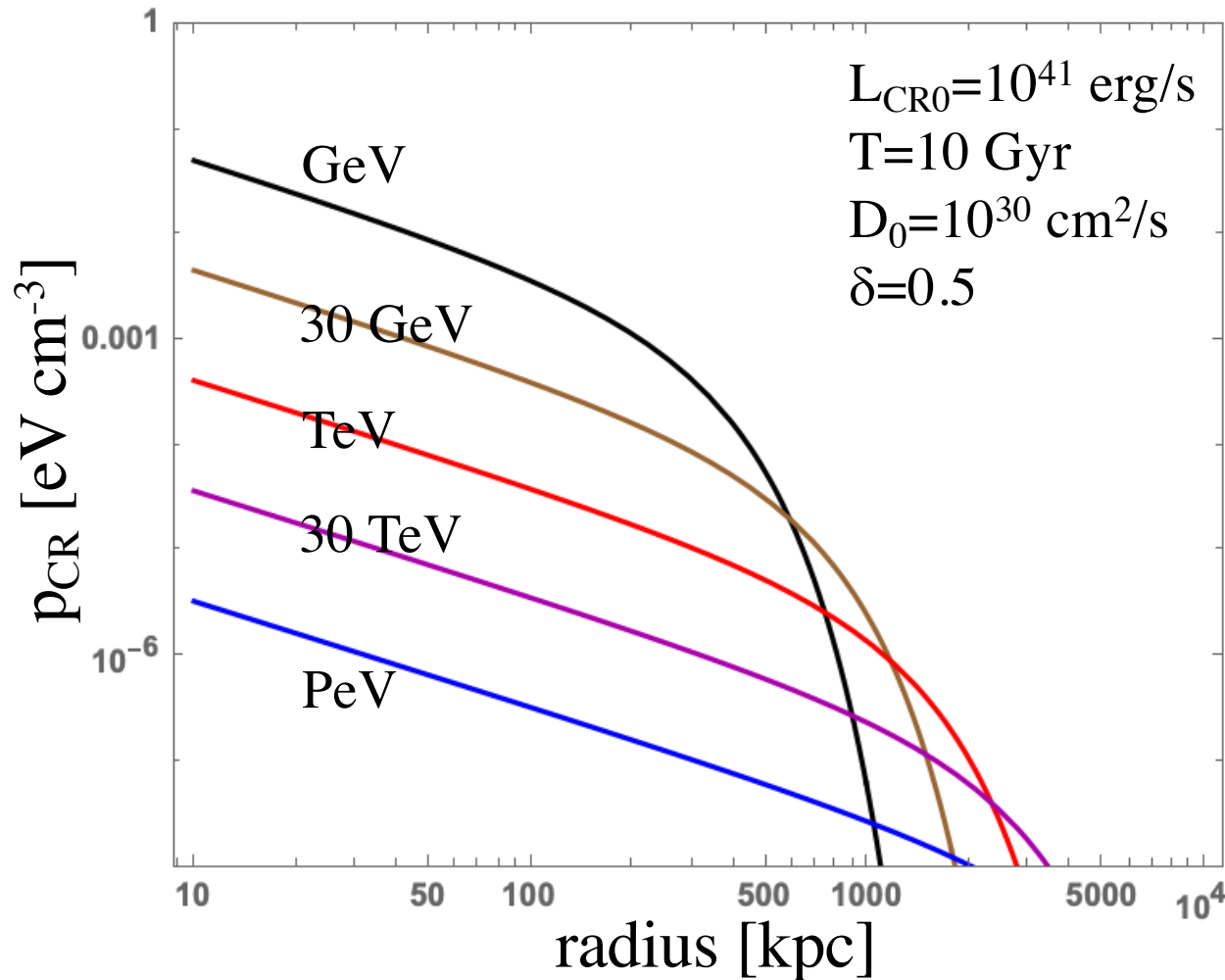
MW-like galaxies stacked sample
 Sunyaev-Zeldovich effect Das+ 25
 X-rays (eROSITA) Zhang+ 24
 $D(\text{GeV}, <100 \text{ kpc}) \sim 10^{30} \text{ cm}^2 \text{s}^{-1}$



simple diffusion solution

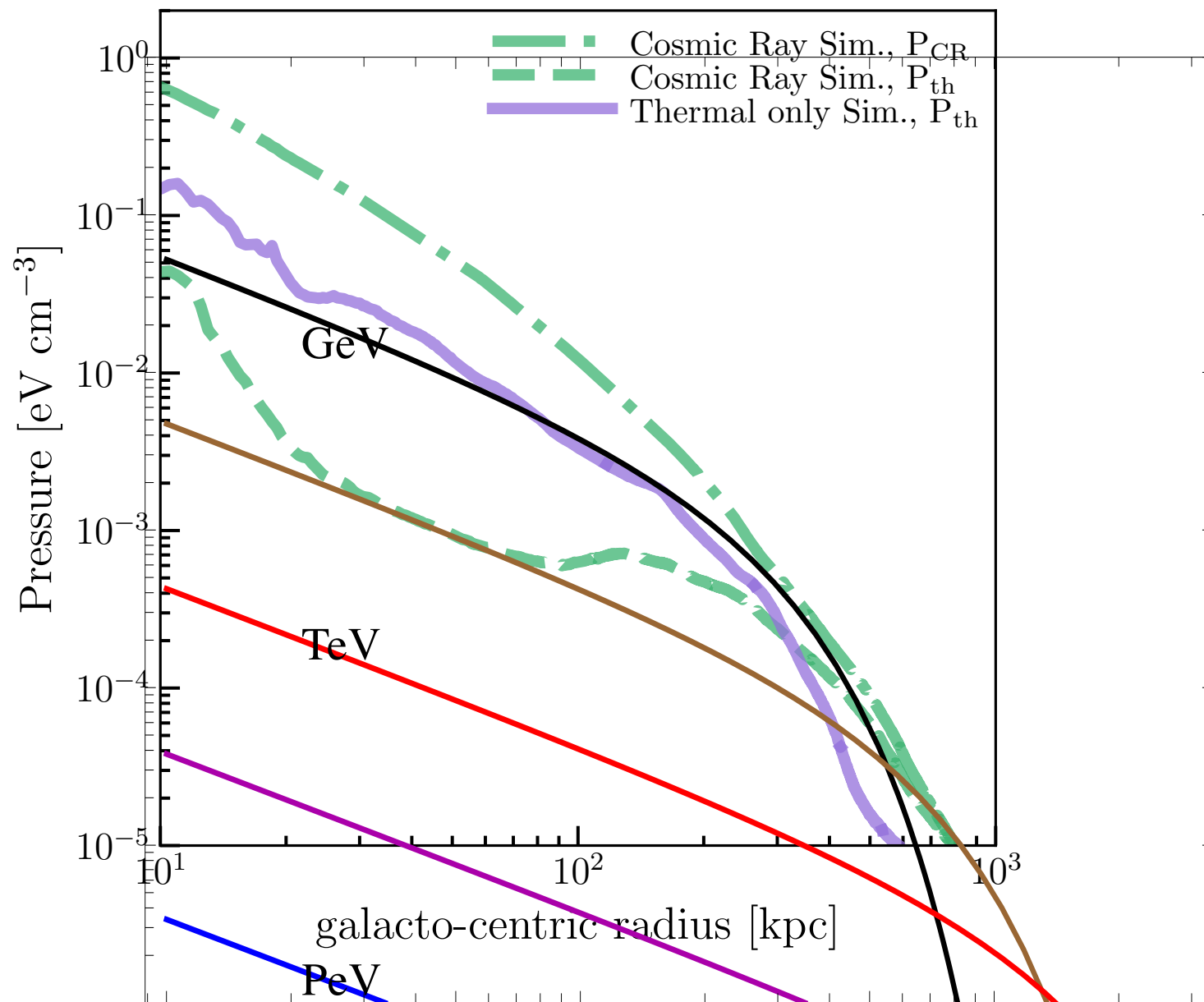
$$U_{\text{CR}}(E, r, t) = \frac{L_{\text{CR}}(E)}{4\pi D(E) r} \operatorname{erfc}\left(\frac{r}{(4D(E)t)^{1/2}}\right) \quad \operatorname{erfc}(x) = 1 - \frac{2}{\sqrt{\pi}} \int_0^x e^{-t^2} dt$$

$$p_{\text{CR}}(E, r, t) = \frac{1}{3} U_{\text{CR}}(E, r, t) \quad D(E) = D_0 (E / \text{GeV})^\delta \quad L_{\text{CR}0} (E / \text{GeV})^{-0.7+\delta}$$



c.f. Ponnada 26 and references therein

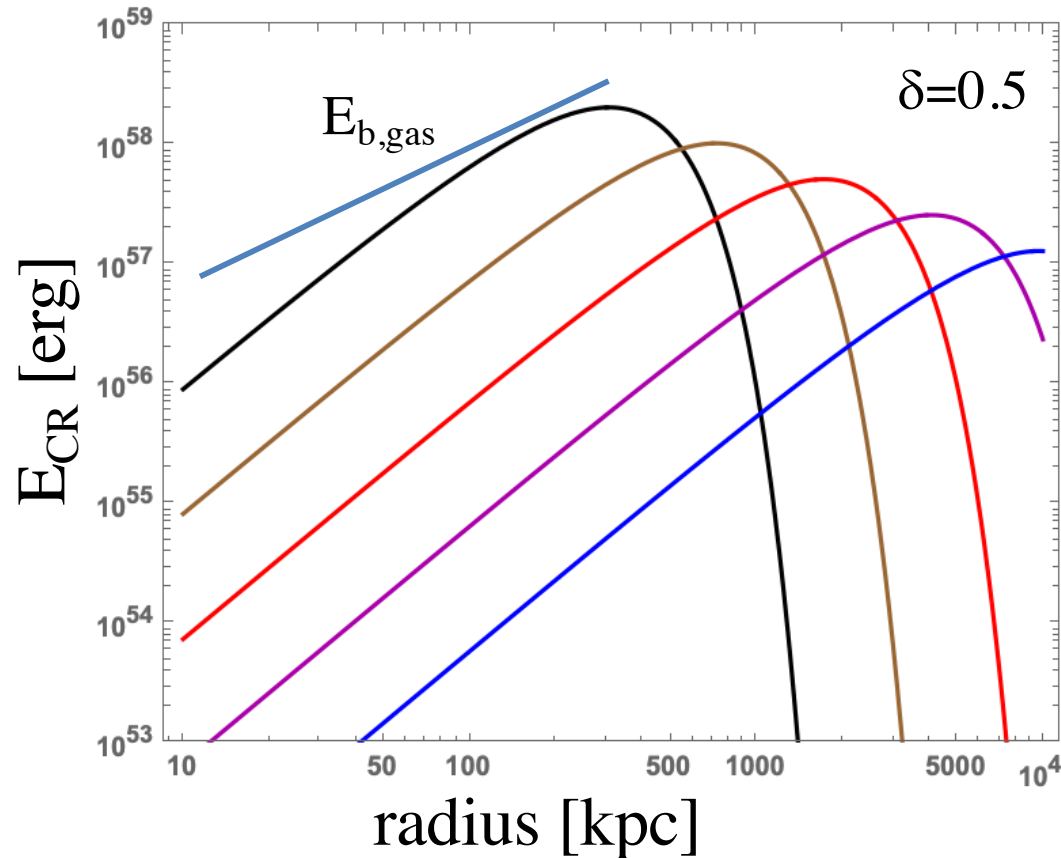
CGM vs CR pressure: Milky Way-like



$L_{\text{CR}}=10^{41} \text{ erg/s}$
 $T=10 \text{ Gyr}$
 $D_0=10^{30} \text{ cm}^2/\text{s}$
 $\delta=0.5$

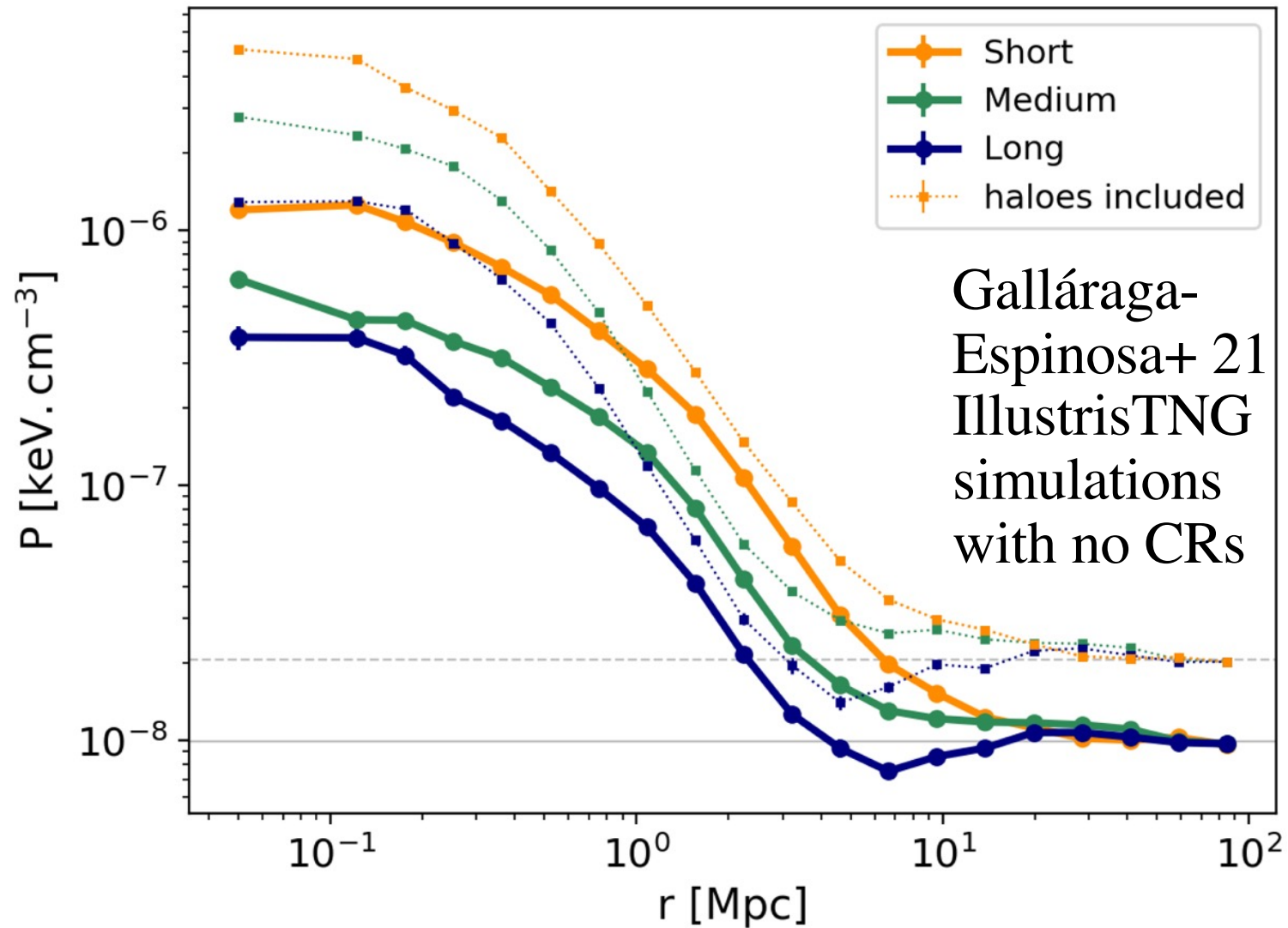
TeV-PeV CRs compared to GeV CRs

$E_{\text{CR}}(E, r, t) = 4 \pi r^3 U_{\text{CR}}(E, r, t)$ total CR energy per log radius



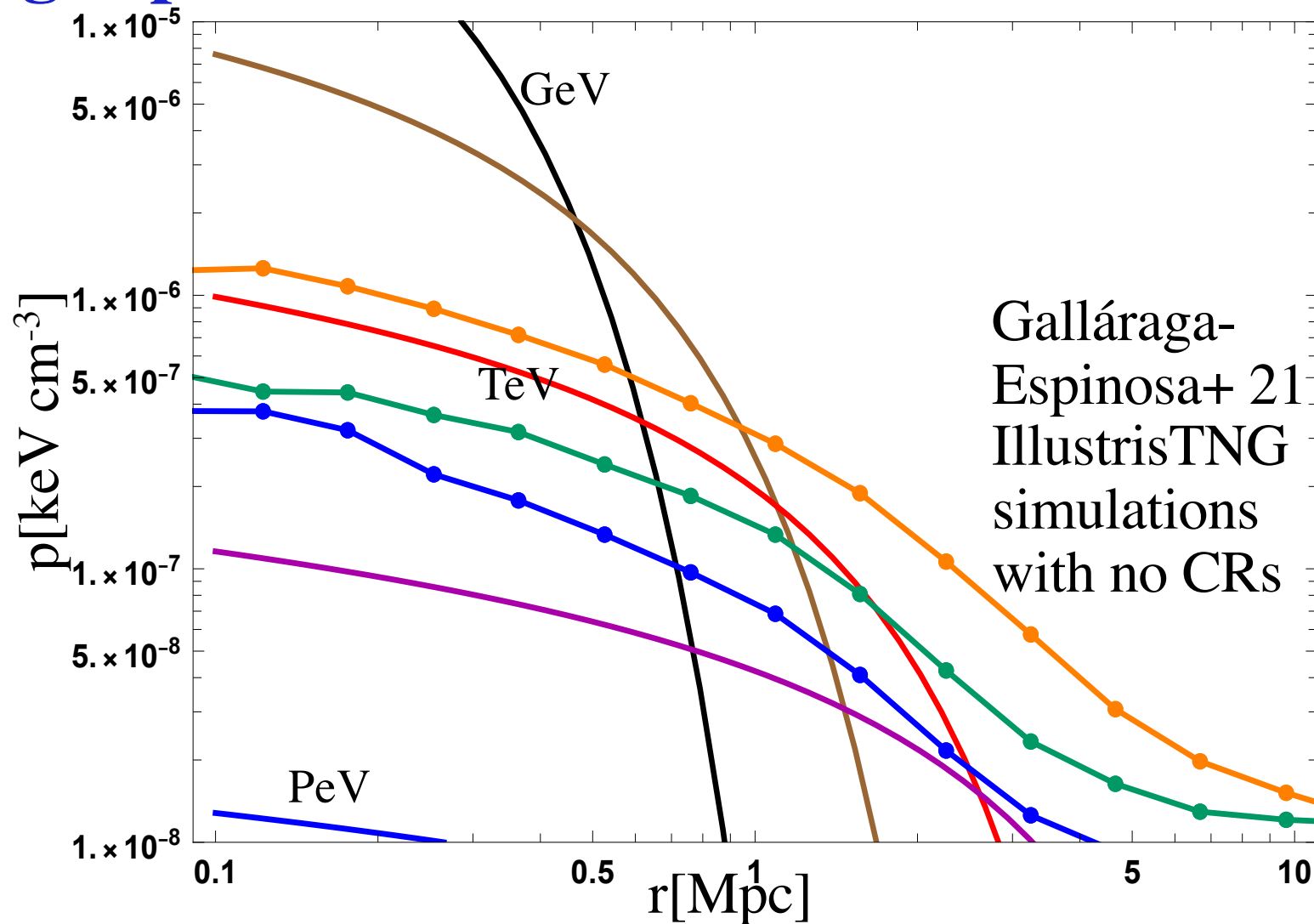
- total energy of escaped TeV-PeV CRs $\sim$ source energy
- gas inflowing from larger R
 - > TeV-PeV CRs affect earlier stages of accretion

gas pressure in filaments



short: $L < 9 \text{ Mpc}$
medium:
 $9 \text{ Mpc} < L < 20 \text{ Mpc}$
long: $L > 20 \text{ Mpc}$

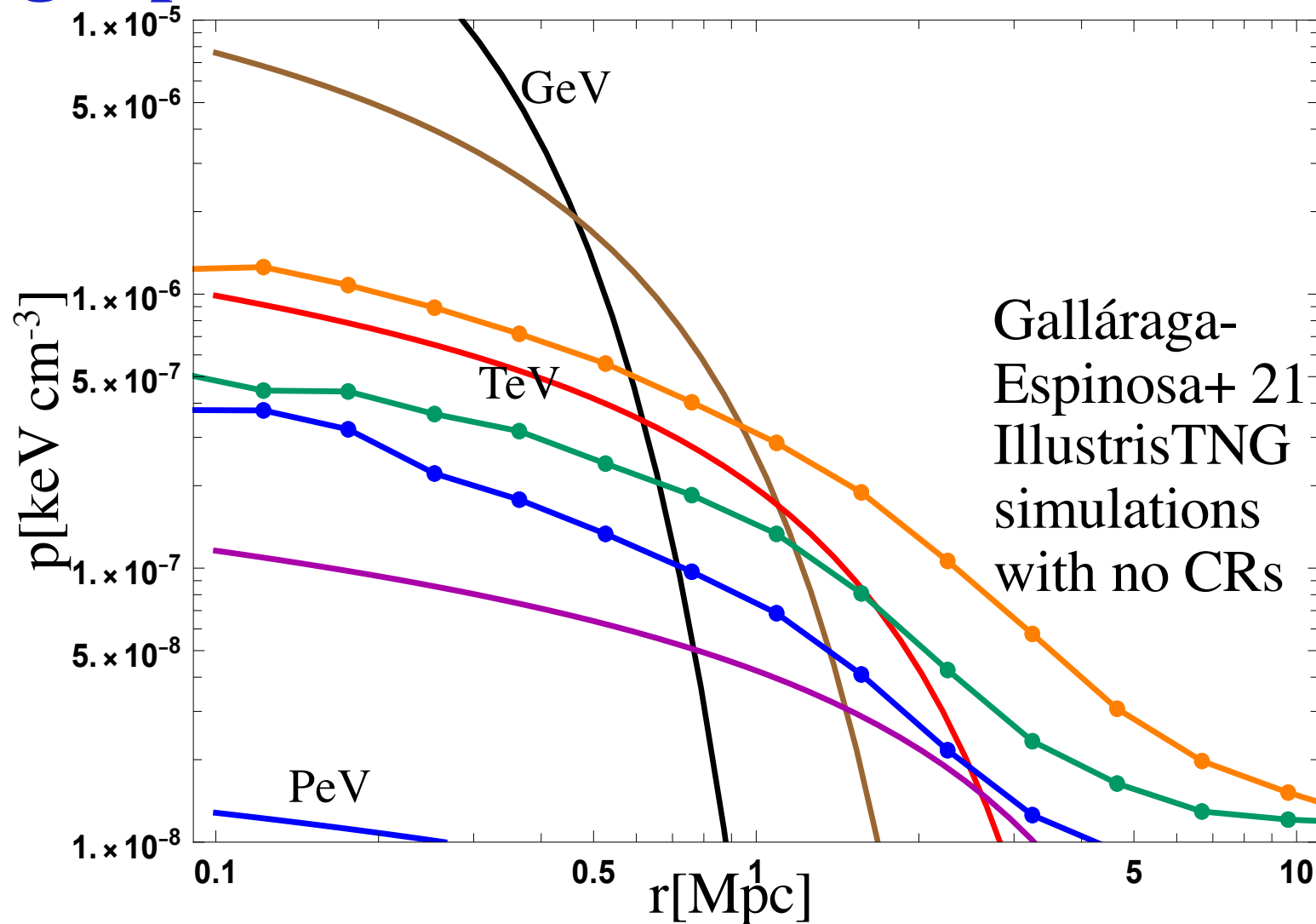
gas pressure in filaments vs CRs



cylindrical
symmetry
MW-like galaxy
population
 $L_{\text{CR}} = 6 \times 10^{42} \text{ erg/s}$
 $T = 10 \text{ Gyr}$
 $D_0 = 10^{30} \text{ cm}^2/\text{s}$
 $\delta = 0.5$

VHECRs can be energetically important vs WHIM!

gas pressure in filaments vs CRs

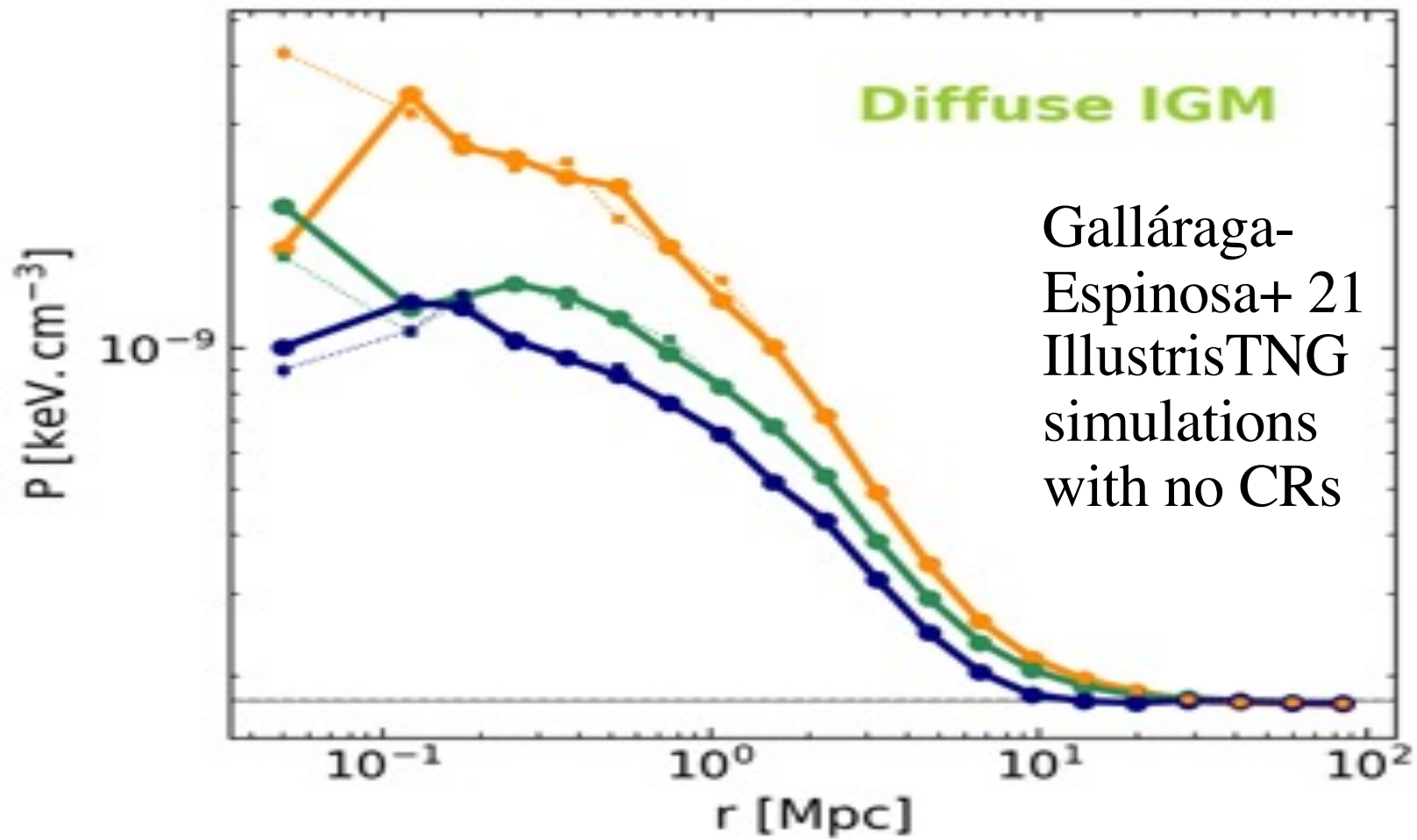


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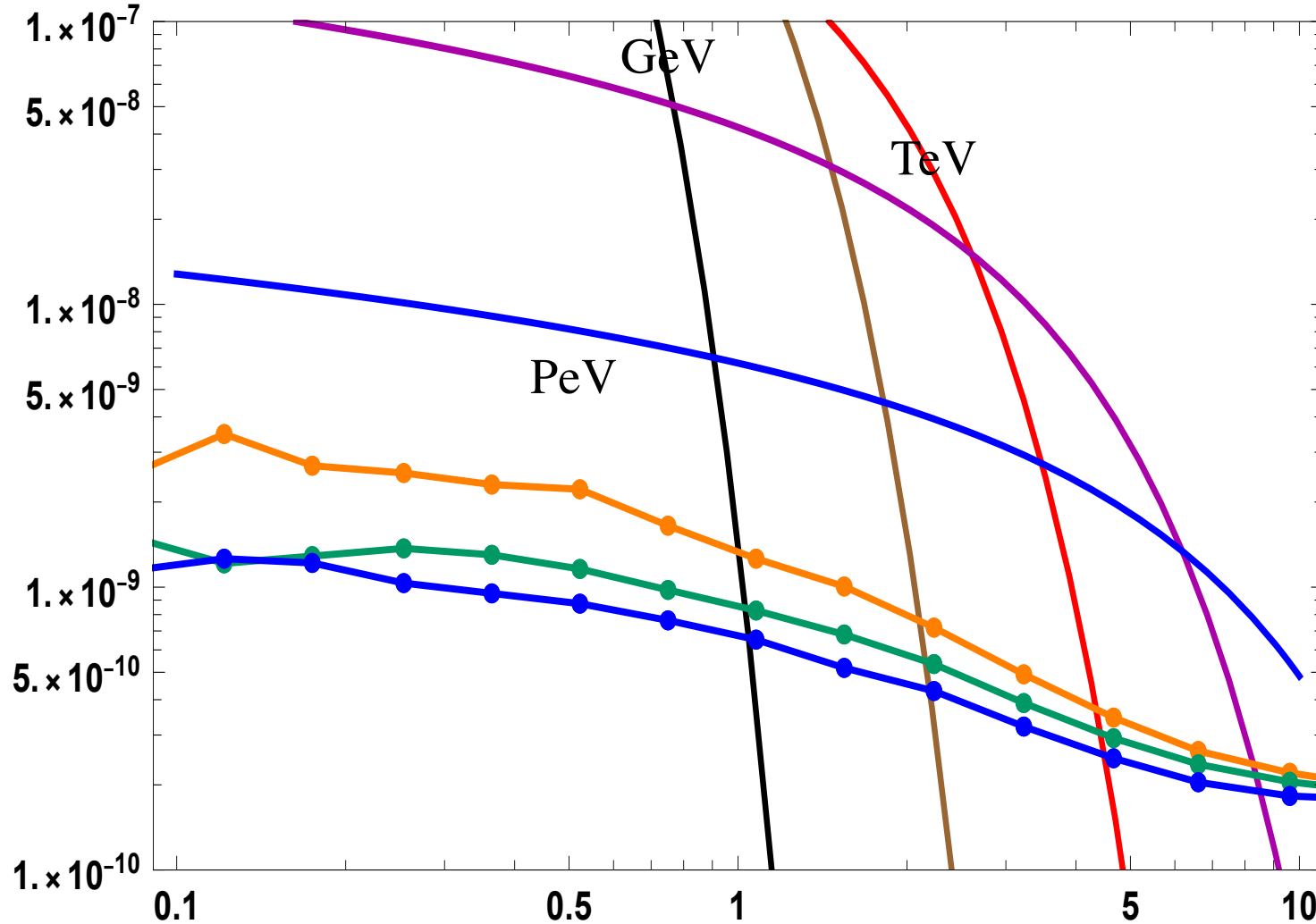
VHECRs can be energetically important vs WHIM!

- affect gas accretion onto halos, star formation rate, galaxy lum. func.
- allow cooler gas, affect observability of WHIM (missing baryons)
- affect measurements of cosmological parameter S_8

cool, diffuse IGM pressure



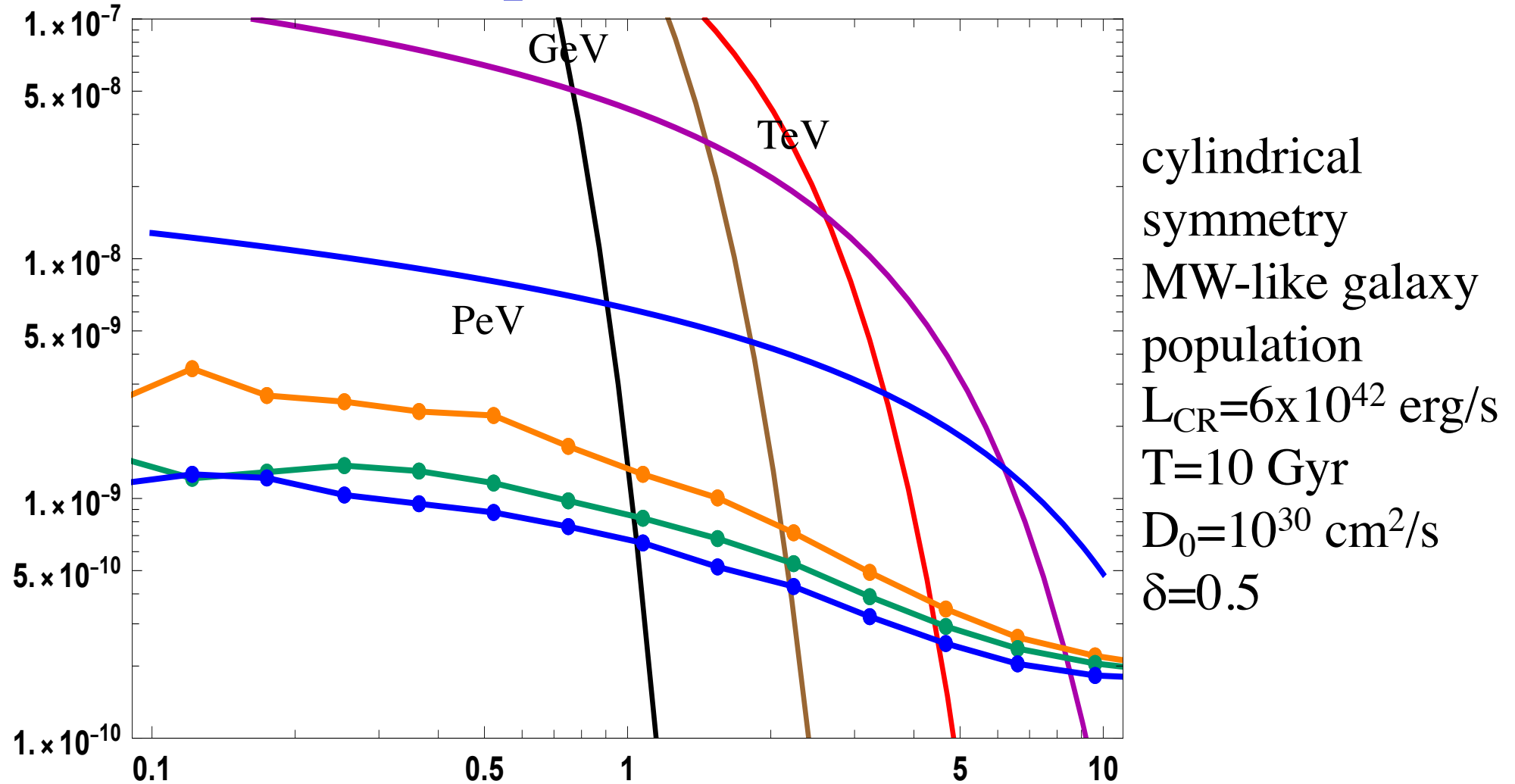
cool, diffuse IGM pressure



cylindrical
symmetry
MW-like galaxy
population
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VHE-(U)HECRs can be energetically important vs cool, diffuse IGM!

cool, diffuse IGM pressure

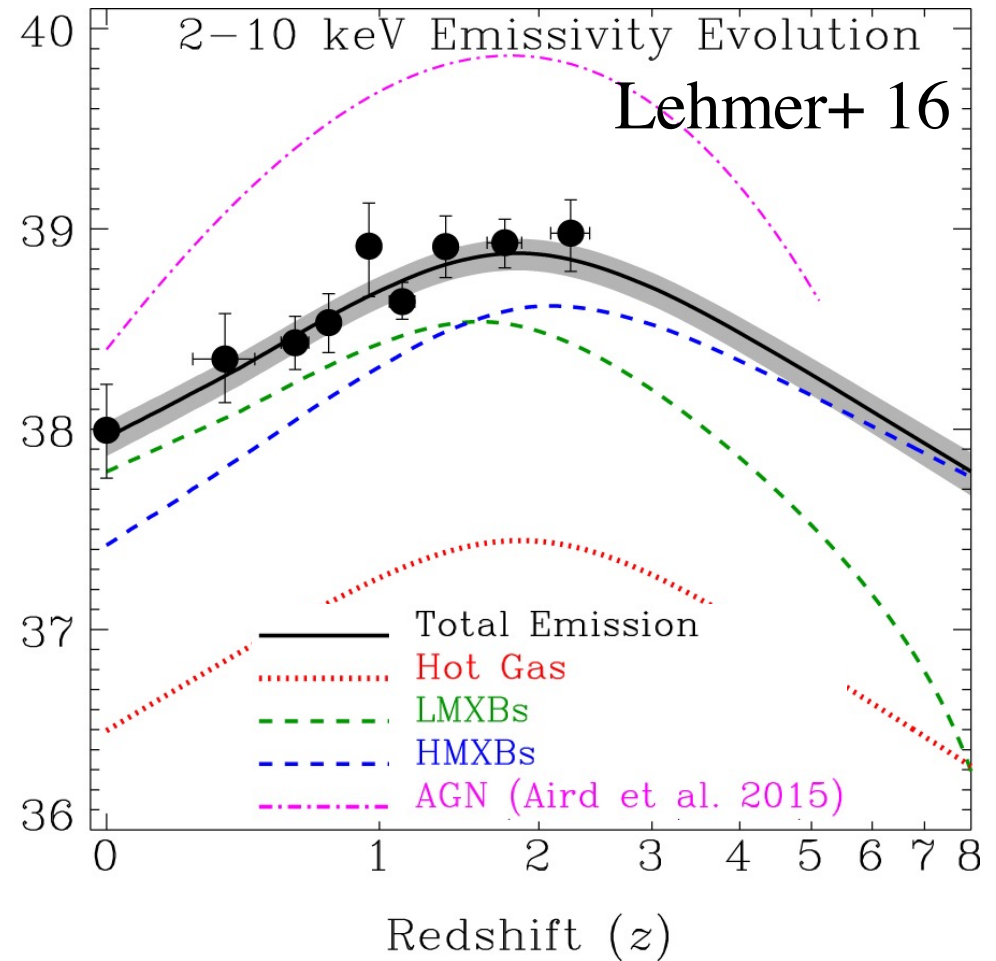
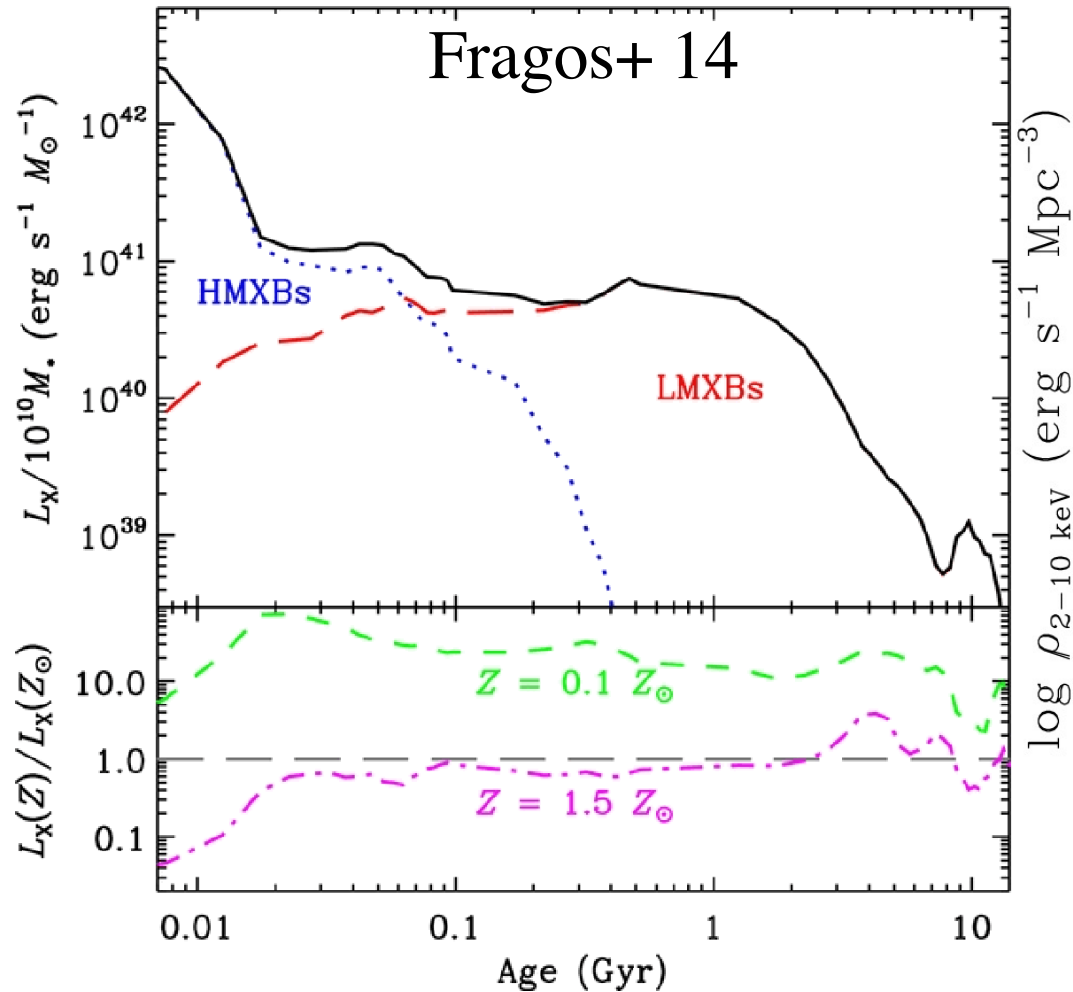


VHE-(U)HECRs can be energetically important vs cool, diffuse IGM!

- affect gas accretion onto filaments, star/structure formation
- relation to observed “anomalies” (excess heating, T inversion, etc)
- > potential constraints from existing data

potential dependence on pevatron identity

microquasars: HMXB vs LMXB: difference in t evolution



PeV CRs in cool IGM: z -dependence of effect on Ly α forest?

“large-scale” pevatrons: Fermi bubble, Galactic wind shock, accretion shock...

summary

- TeV-PeV CRs vs GeV CRs: energetically minor in MW disk, but possibly not so minor at outer CGM and beyond
- TeV CRs can reach the WHIM and be energetically important
 - > potential effect on gas accretion onto halos, WHIM observability, S8 measurement
- TeV-PeV CRs can reach cool diffuse IGM and be energetically important
 - > potential effect on gas accretion onto filaments, relevance for observed Ly α forest anomalies
- potential dependence on nature of source
 - Identifying pevatrons may also be important for understanding IGM evolution!
- (- observational tests: challenging)
 - c.f. PeV γ has some advantages for MW CGCRs, see SI et al. ICRC 2025; in prep.