

Collapse of rotating massive stars

Sho Fujibayashi (Tohoku U.)

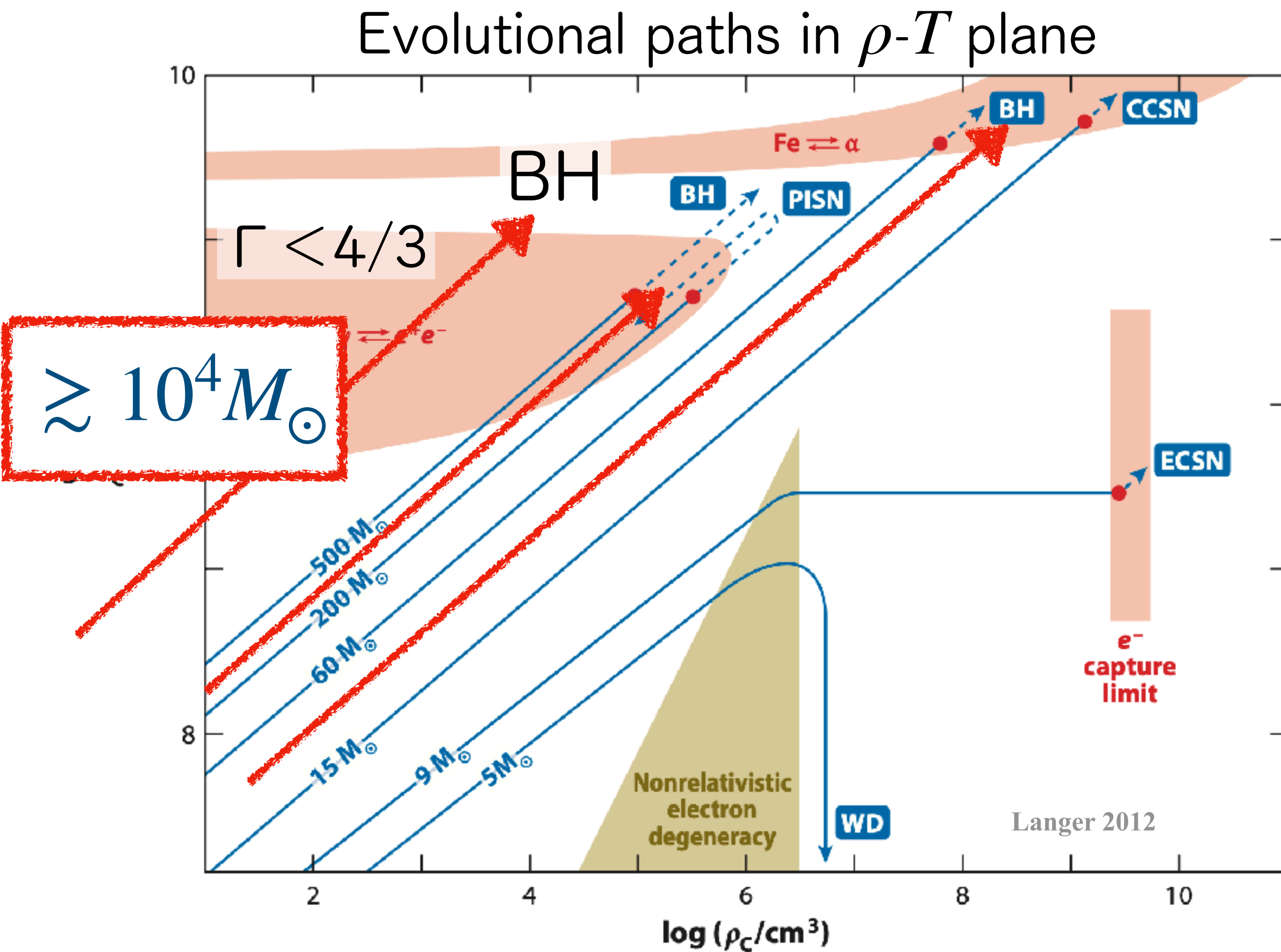
in collaboration with

C. Jockel, K. Ioka, K. Kawaguchi, A. T.-L. Lok, S. Wanajo, M. Shibata

Multi-messenger Annual Conference 2025, 2025.11.19, Naruko Kanko Hotel, Miyagi



Fate of massive stars



- ✓ $M_{\text{ZAMS}} \gtrsim (8 - 10)M_\odot$
 - Iron core formation → Gravitational collapse
 - Core-collapse SNe (possible BH formation)
- ✓ $M_{\text{ZAMS}} \gtrsim 130M_\odot$ (Very massive star; VMS)
 - Unstable by e^-e^+ production
 - Pair instability SN
 - or, $M_{\text{ZAMS}} \gtrsim 260M_\odot \rightarrow$ BH formation
- ✓ $M_{\text{ZAMS}} \gtrsim 10^4M_\odot$ (Supermassive star; SMS)
 - Collapse by GR instability
 - BH formation

Fate of massive stars ... with rotation

✓ $M_{\text{ZAMS}} \gtrsim (8 - 10)M_{\odot}$

- Iron core formation → Gravitational collapse
- Core-collapse SNe (possible BH formation)

✓ $M_{\text{ZAMS}} \gtrsim 130M_{\odot}$ (VMS)

- Unstable by $e^{-}e^{+}$ production
- Pair instability SN
- or, $M_{\text{ZAMS}} \gtrsim 260M_{\odot} \rightarrow$ BH formation

✓ $M_{\text{ZAMS}} \gtrsim 10^4M_{\odot}$ (SMS)

- Collapse by GR instability
- BH formation

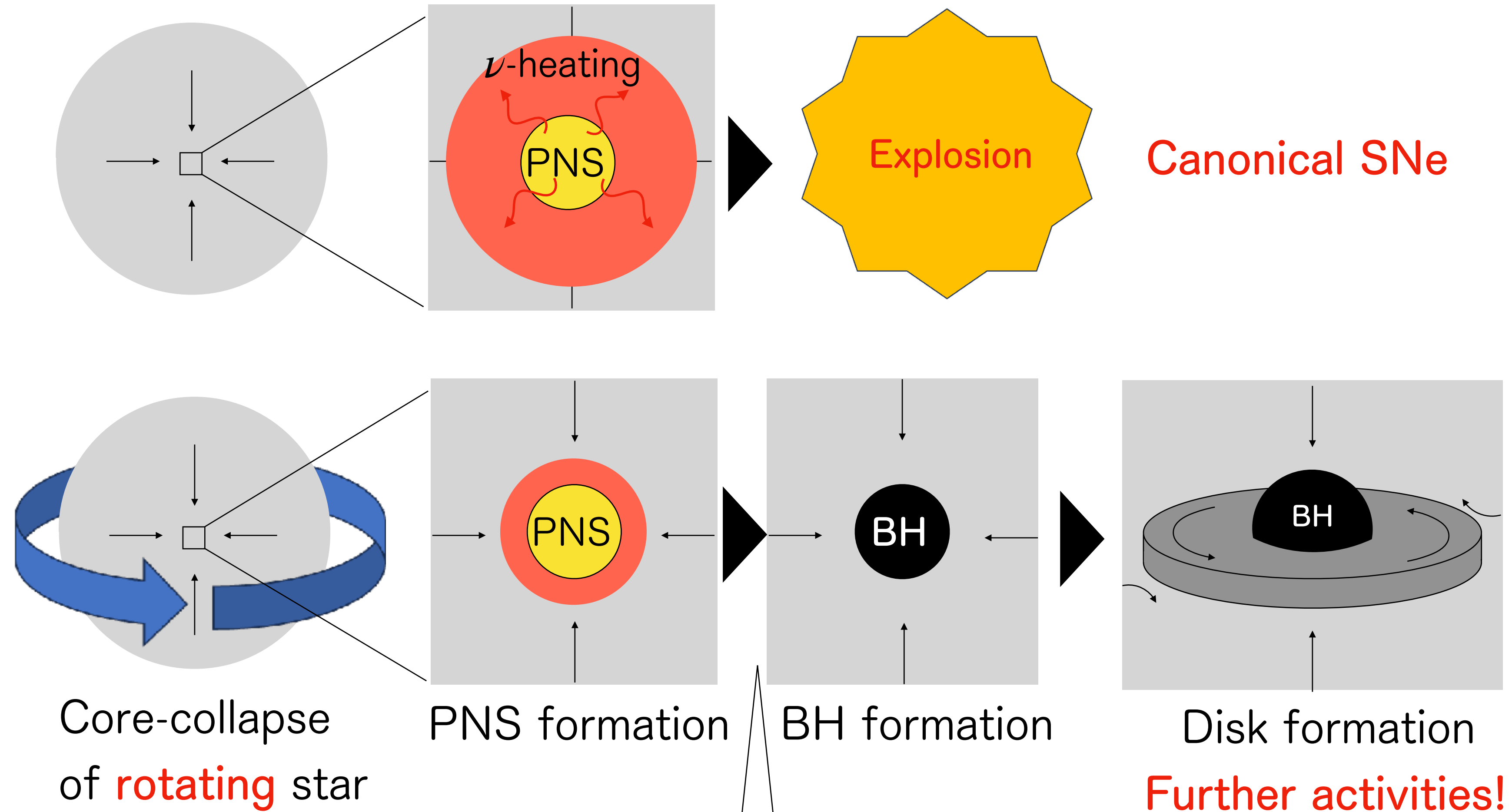
- Rapidly rotating NS (with B-field)
- Spinning BH with accretion disk

- Change PISN explosion feature?
- Spinning BH with accretion disk

- Spinning BH with accretion disk

Usual massive star: $M \sim 10M_{\odot}$

Collapsar



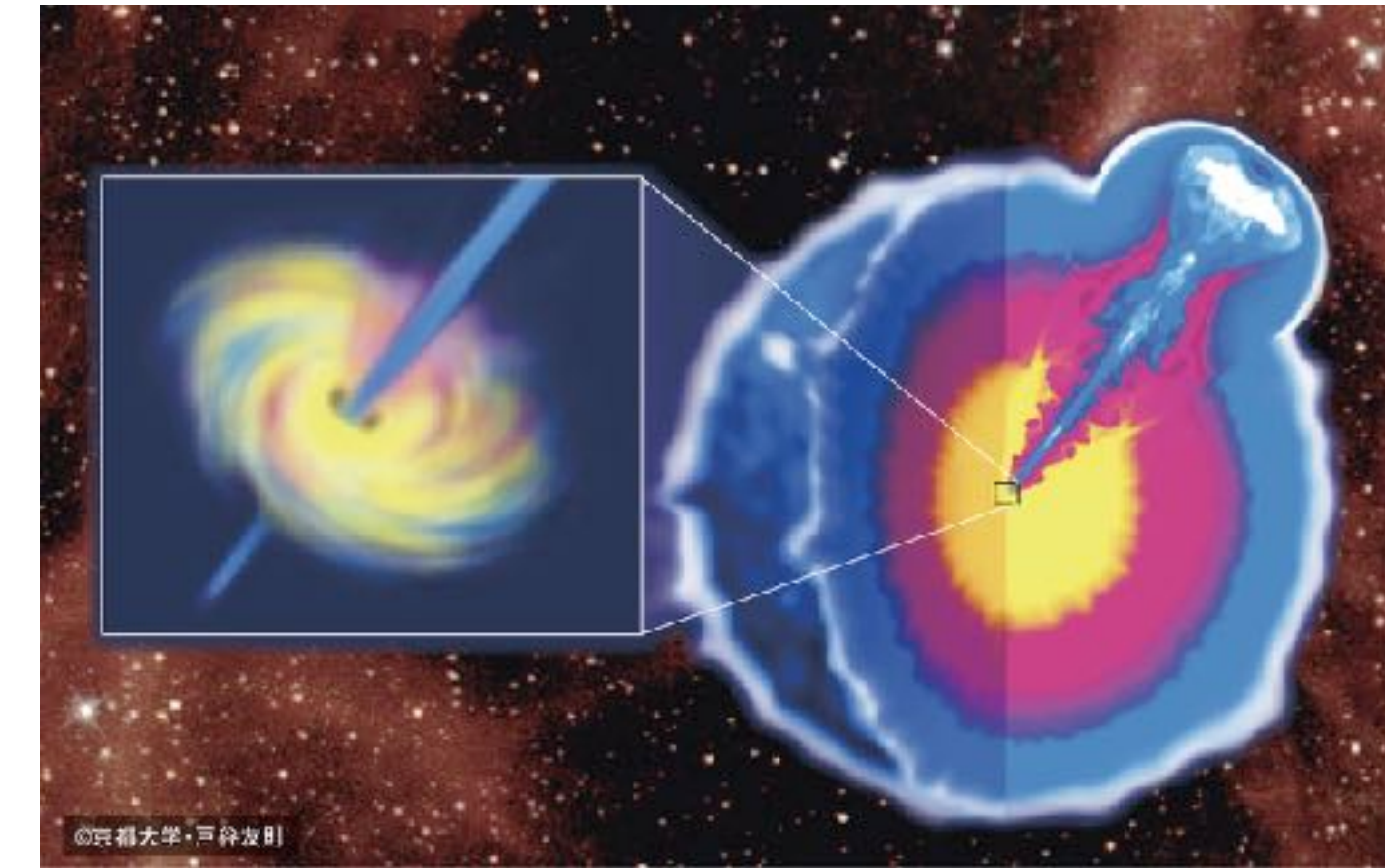
Explosion during PNS phase fails
if, e.g., the core compactness is too high.

Possible BH-disk activities

Gamma-ray bursts (GRBs)

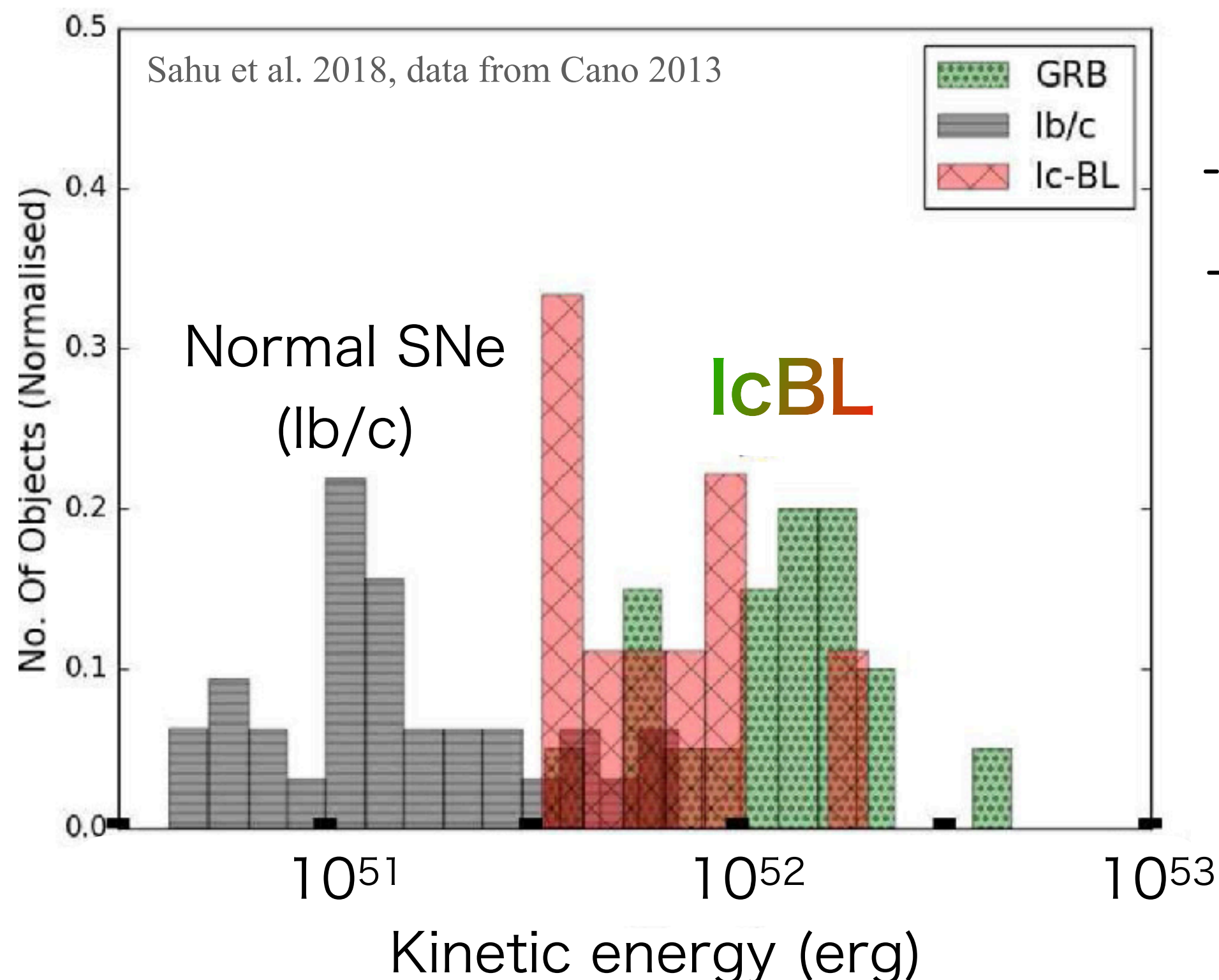
BH-disk is one of the promising central engines.

(e.g., Woosley et al. 1993...)



Broad-lined type Ic SNe (SNe Ic-BL; Hypernovae)

A high-energy class of supernovae



- Explosion (kinetic) energy: ~ 10 times higher than normal SNe
- ^{56}Ni mass : larger than normal SNe

They are sometimes associated with long-GRBs
Something to do with BH-disk activity.

BH-disk activities and GRB-SN

Broad-lined type Ic SNe (SNe Ic-BL; Hypernovae)

A high-energy class of supernovae

- Explosion (kinetic) energy: ~10 times higher than normal SNe
- ^{56}Ni mass : larger than normal SNe

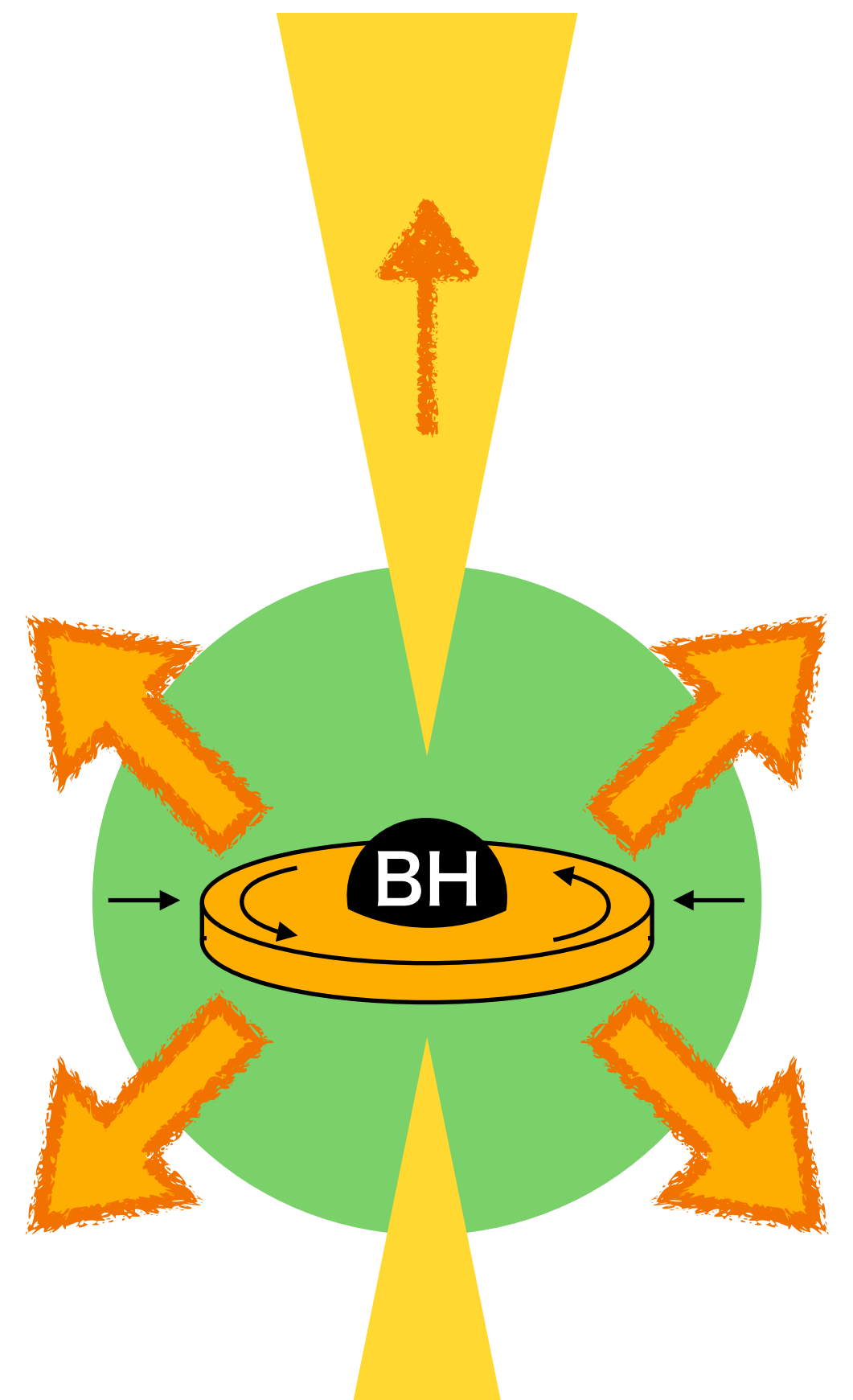
They are sometimes associated with long-GRBs
Something to do with BH-disk activity.

Disk outflow (MacFadyen & Woosley 1999)

Energy generated by viscous accretion:

$$\frac{GM_{\text{BH}}M_{\text{disk}}}{r_{\text{disk}}} \approx 3 \times 10^{52} \text{ erg} \left(\frac{M_{\text{BH}}}{10M_{\odot}} \right) \left(\frac{M_{\text{acc}}}{0.1M_{\odot}} \right) \left(\frac{r_{\text{disk}}}{10^7 \text{ cm}} \right)^{-1}$$

Viscosity-driven outflow from disk would naturally explain such SNe



MHD modeling of collapsar

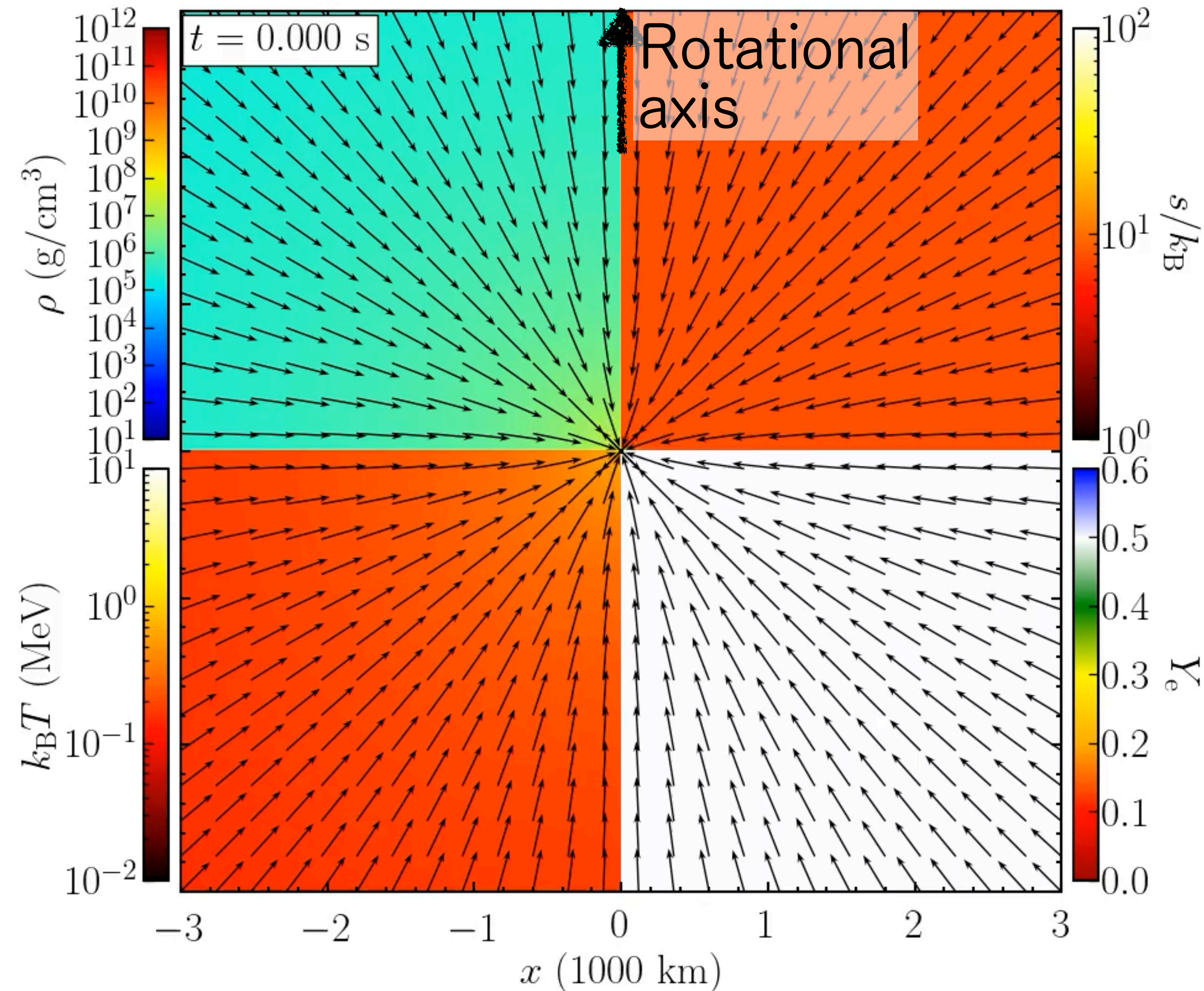
Shibata, SF+2025

$M_{\text{ZAMS}} = 35M_{\odot}$ star, starting from toroidal field

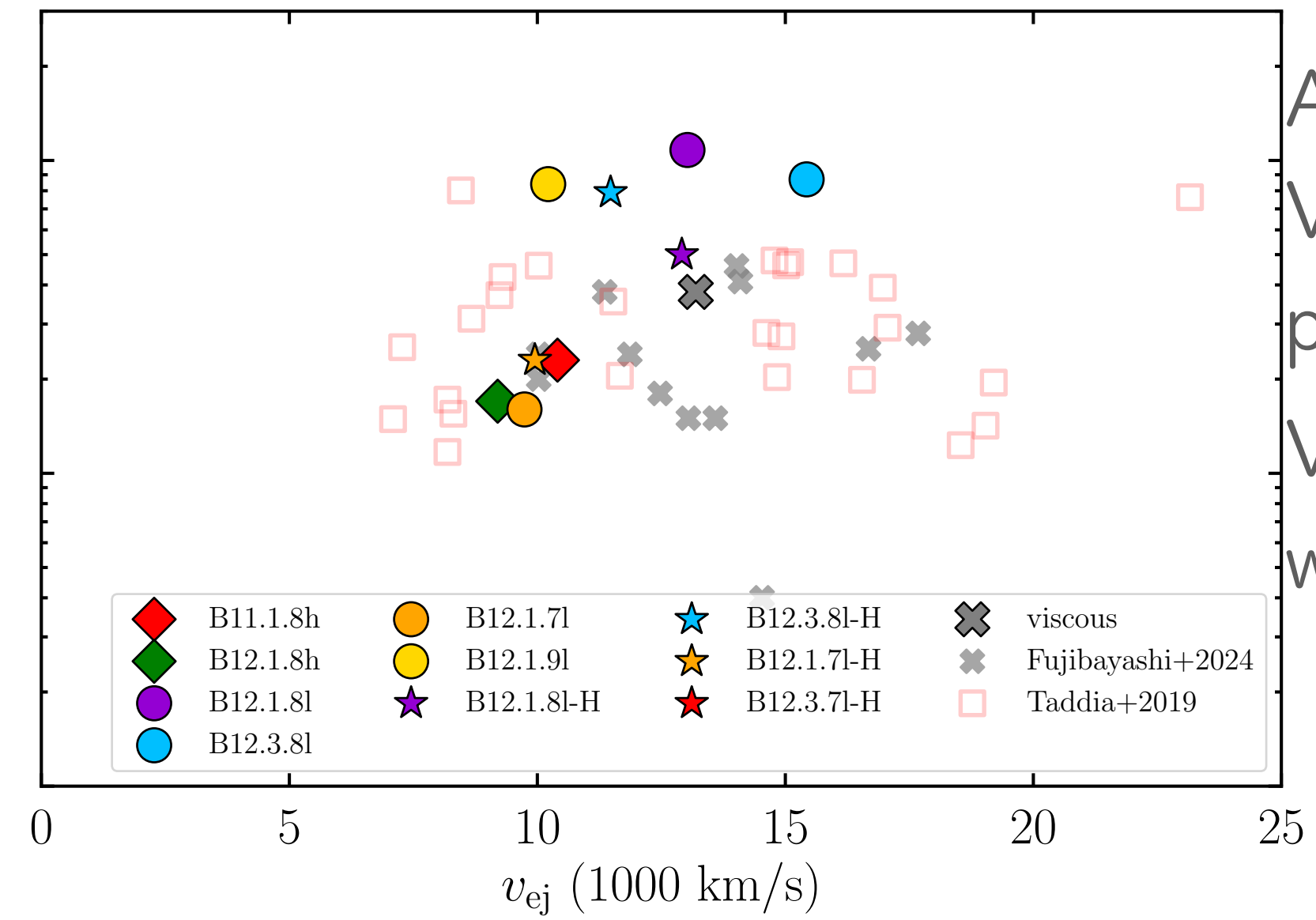
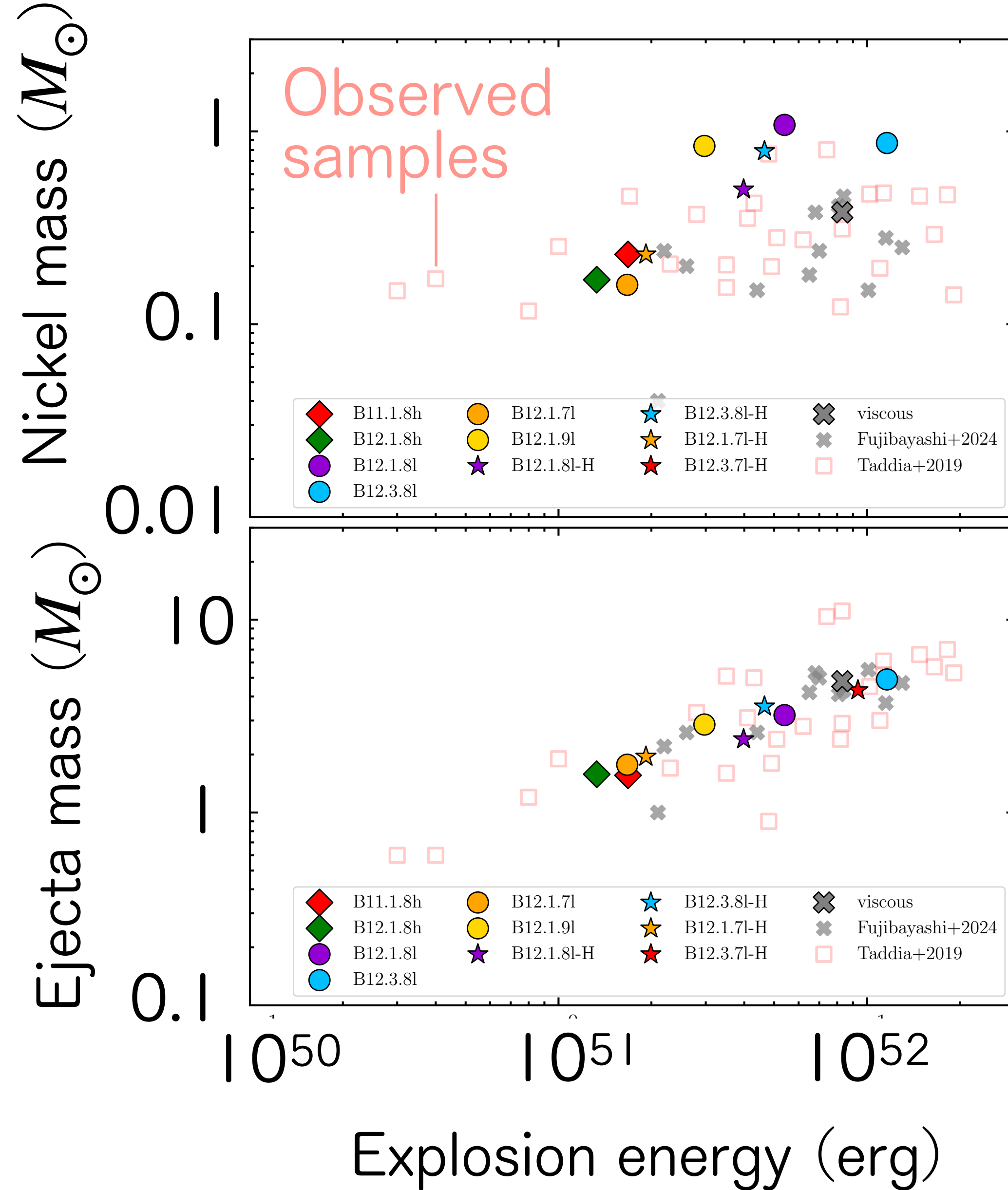
2D-axisymmetric simulation with solving

- ✓ Einstein's equation
- ✓ Neutrino radiation transfer equation
- ✓ Magneto-hydrodynamics equation

→ BZ-driven Jet
+ Spherical wind components



Observables



All 35Msun models
Variation is made by
parameters and resolution
Viscous models shown
with grey color.

$$E_{\text{exp}} = 10^{51} - 10^{52} \text{ erg},$$

$$M_{\text{Ni}} \sim 0.1 - 1 M_{\odot},$$

$$V_{\text{ej}} \sim 10000 \text{ km/s}$$

It may be a good model for SNe Ic-BL.

Global feature is similar to viscous models.

Very massive and Supermassive stars:

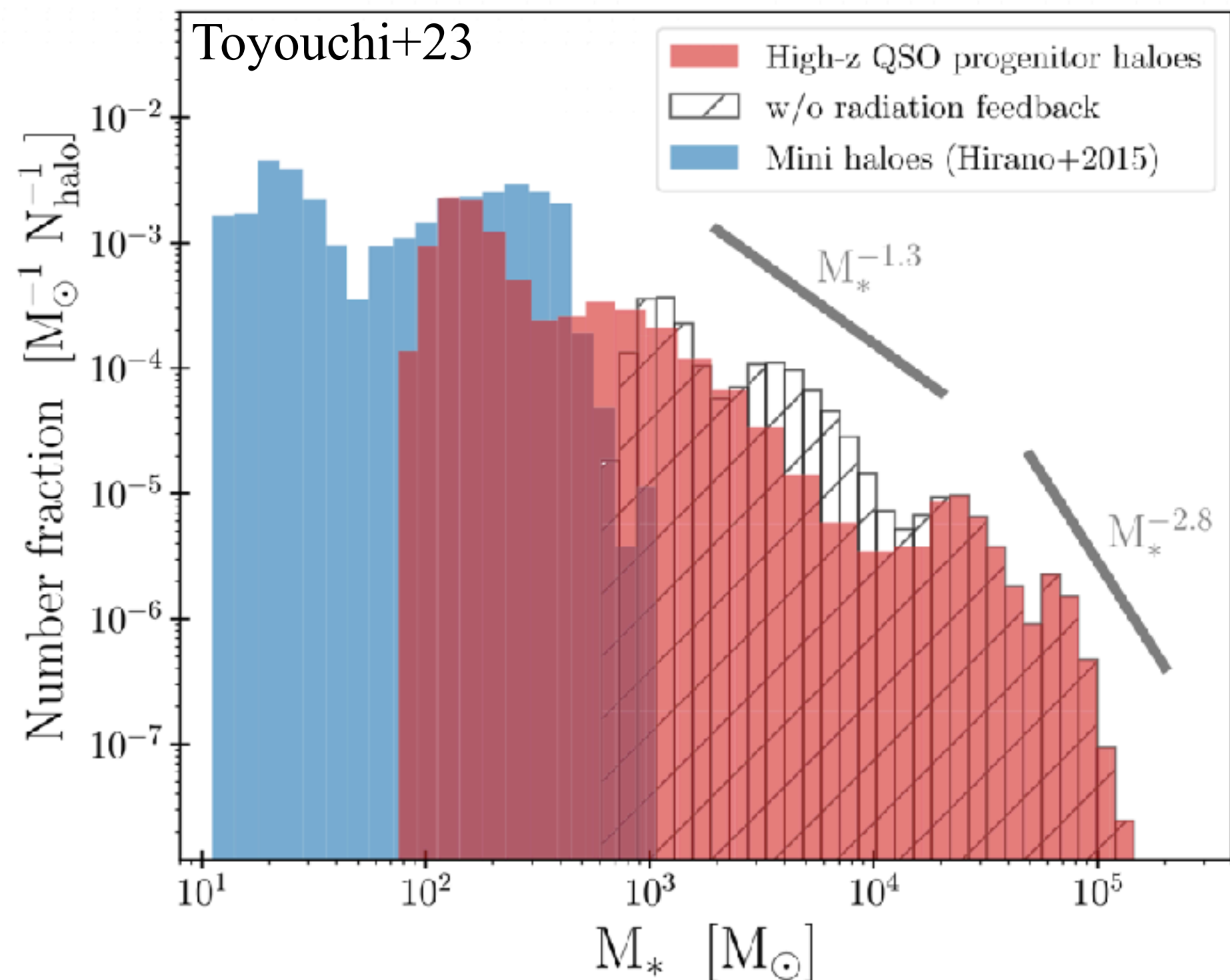
$$M \sim 10^3 - 10^6 M_{\odot}$$

Fate of rotating stars with $M \gtrsim 10^3 M_\odot$

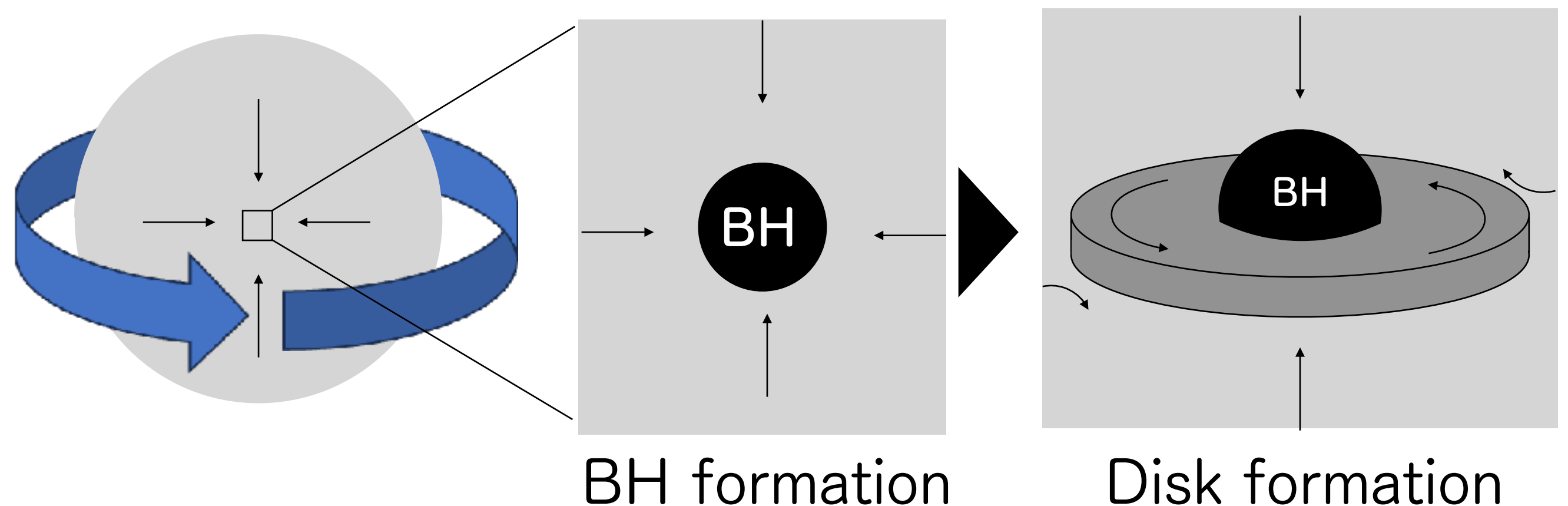
Hypothetical object in the early universe (potential seed of high-z SMBH)

$M \lesssim 10^4 M_\odot$: Very massive star (collapse via pair instability)

$M \gtrsim 10^4 M_\odot$: Supermassive star (collapse via GR instability)



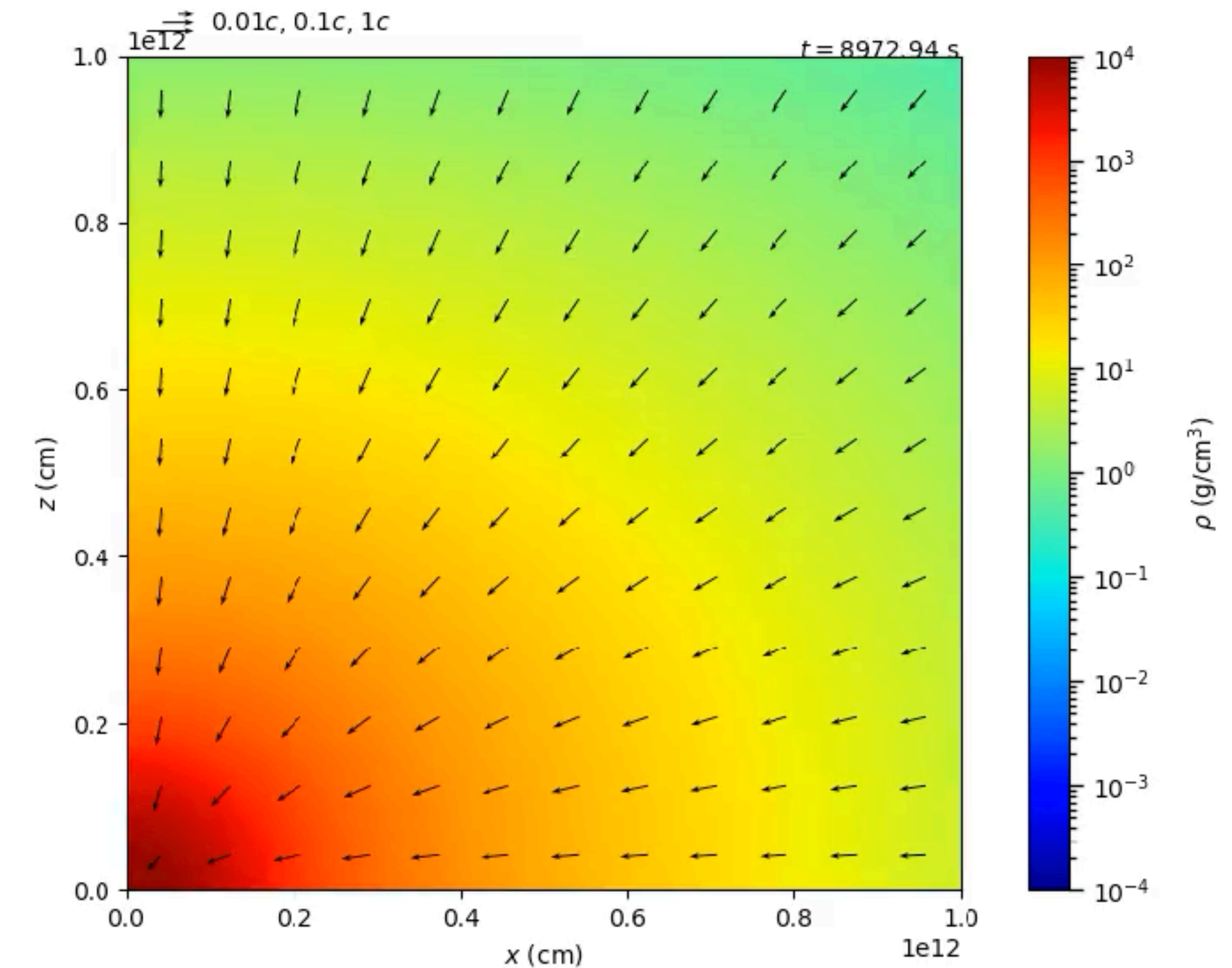
They just leave behind BHs....?
Something non-trivial occurs **with rotation**.



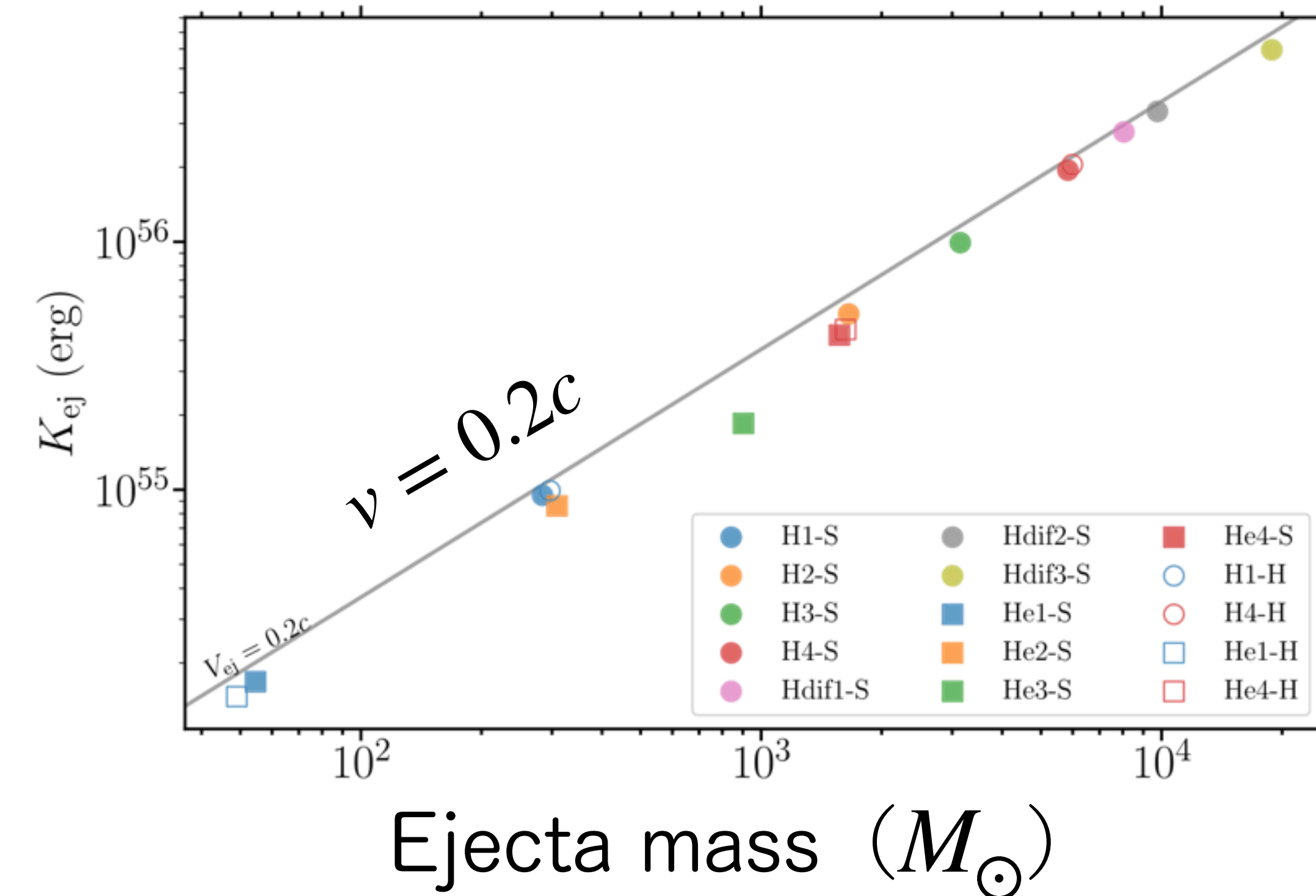
Disk bounce to mass ejection

Bounce of the disk just after the formation
 → A part of the disk becomes unbound.

Uchida et al. (2017)



Explosion energy (erg)



Similar average velocity $v = 0.2c$

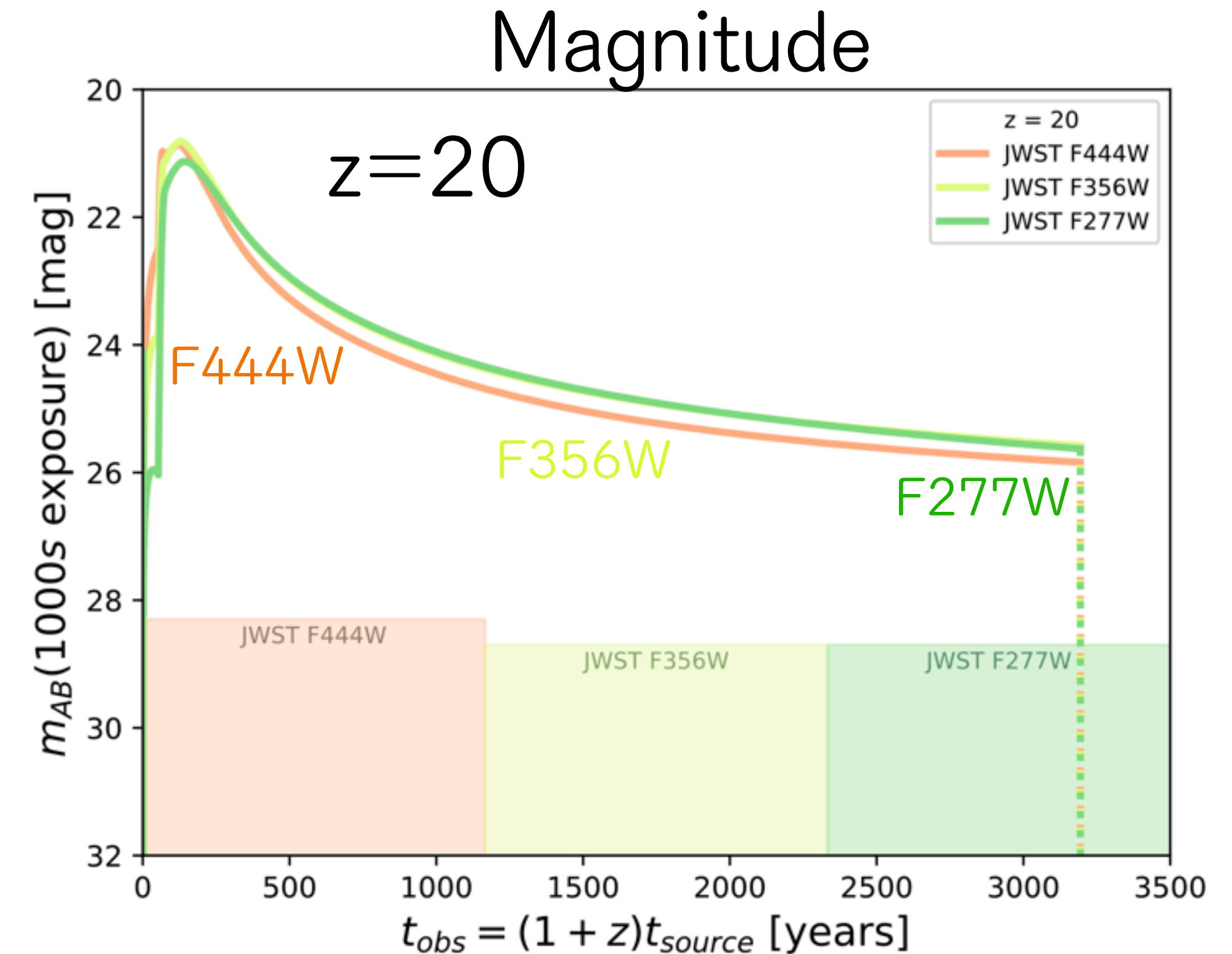
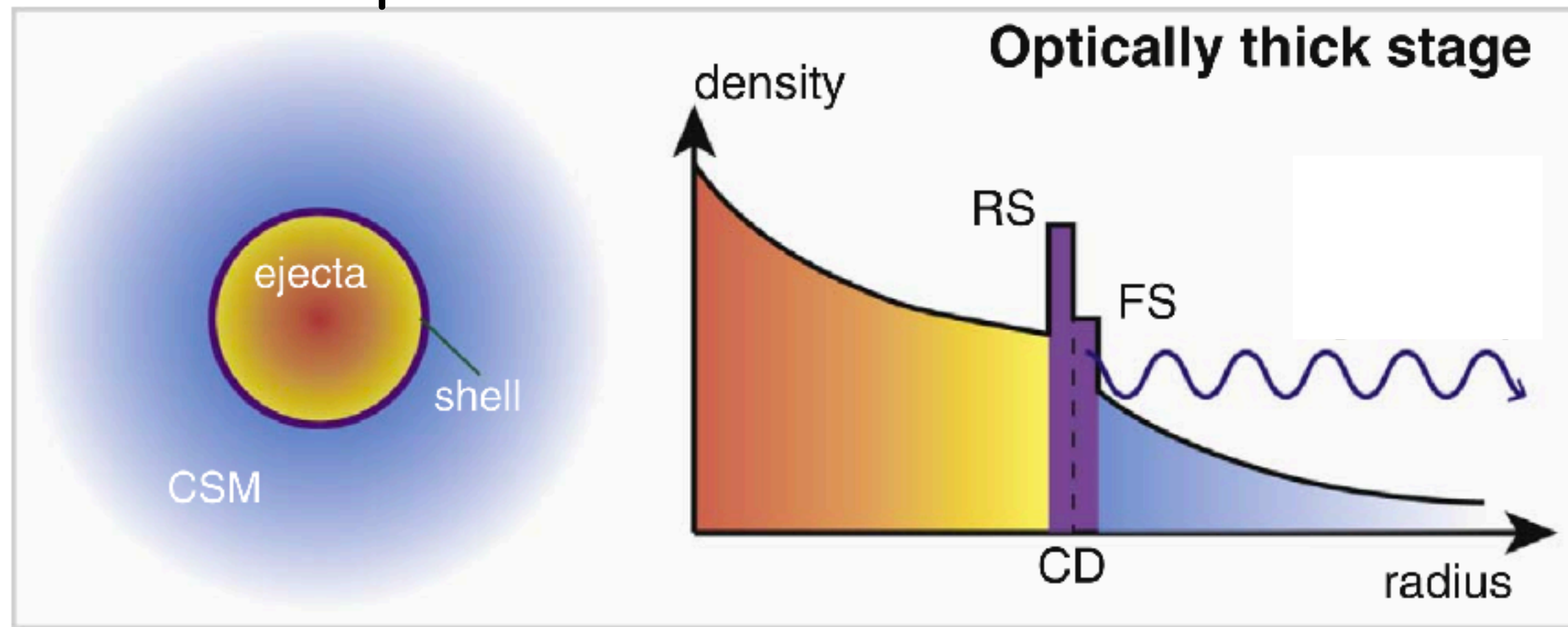
$$M_{\text{ej}} \approx 10^{-2}M, K_{\text{ej}} \approx 10^{-4}Mc^2$$

Similar escape vel. at disk formation (at \sim ISCO)
 & Similar structure ($n=3$ polytrope)
 & Coherent collapse

Observational properties

Jockel, Kawaguchi, SF, Shibata, arXiv:2507.15556 (in pres.)

Realistic picture



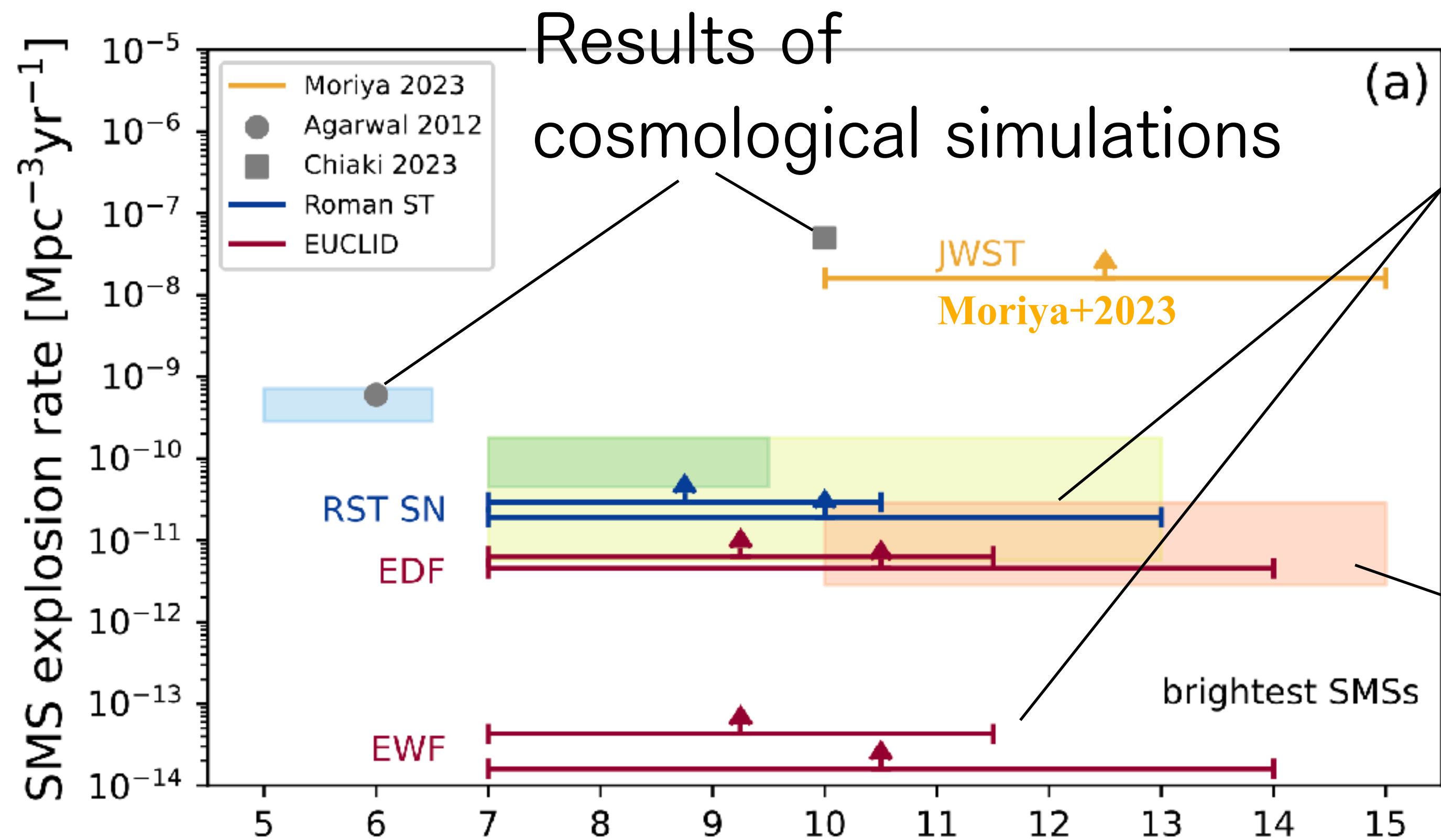
Most bright cases ($E_{exp} \sim 10^{56}$ erg) can be observed at $z=20$ (if found).

~10 yr timescale at **source frame**. $\rightarrow \sim(1+z)$ 10 yr for us.

\rightarrow Will be observed as nearly persistent "transients"

Detection prospects

Most bright cases (SMS mass $\sim 10^5 M_\odot$)

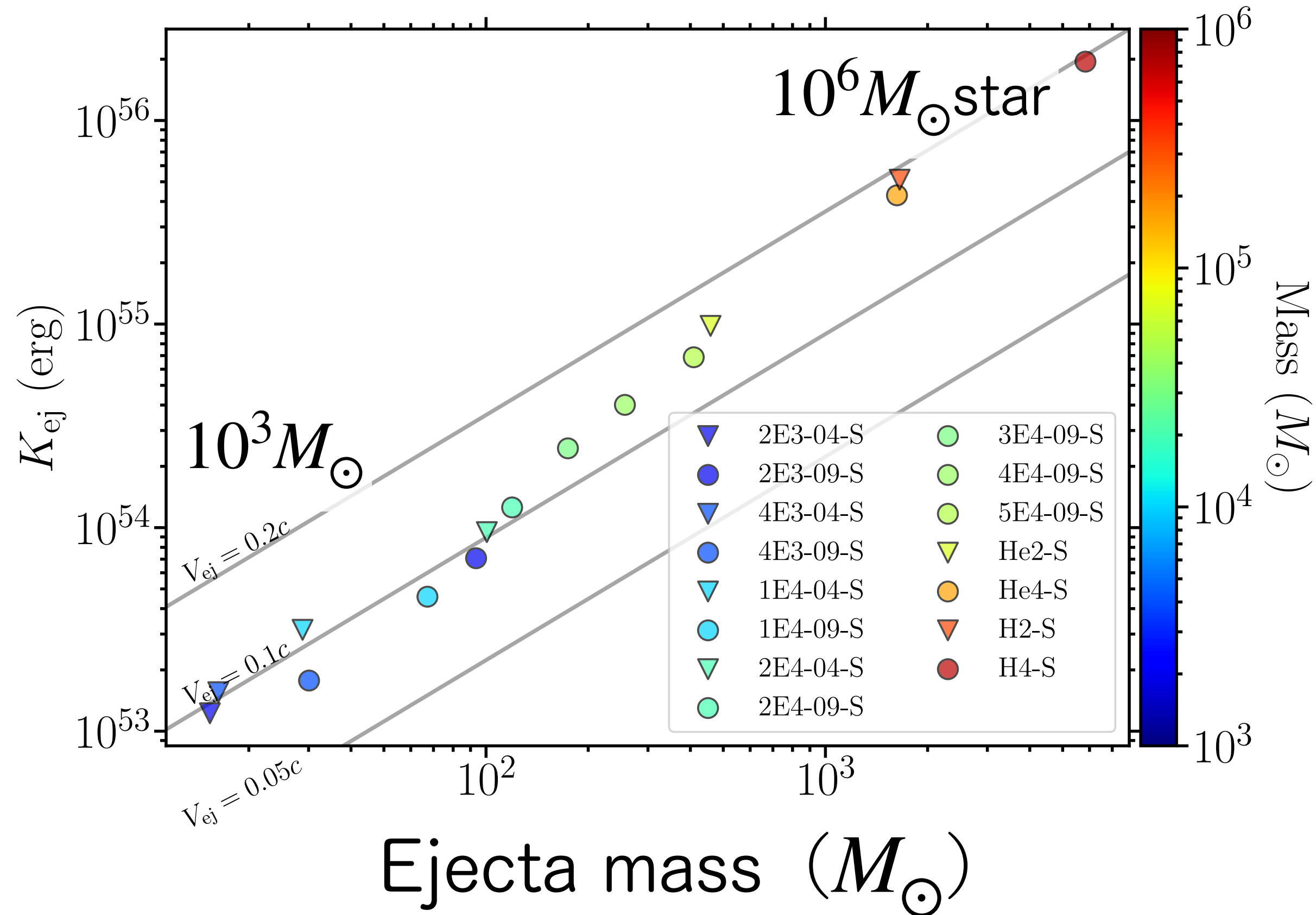


Upper limit that can be set by each survey program of Roman or Euclid.

SFR x (ratio of $> 10^5 M_\odot$ stars)
(Harikane+23)
assuming Salpeter's IMF

Lower mass (down to $10^3 M_\odot$)

Explosion energy (erg)



Slower average velocity
for smaller mass stars

Central region collapse earlier

→ Smaller BH and disk masses

→ Smaller energy available through bounce
compared to the stellar binding energy.

For $\sim 100 M_\odot$ star, the bounce cannot unbound matter
→ to normal collapsar

Summary

- Collapse of rotating normal massive star ($\sim 10M_{\odot}$)
 - Outflow from disk around BH can power SNe Ic-BL.
 - BZ process could work for powering GRB jet.
- Collapse of rotating SMS and VMS ($\sim 10^3 - 10^6 M_{\odot}$)
 - Explosion induced by disk bounce:
 $M_{\text{ej}} \approx 10^{-2}M$, $K_{\text{ej}} \approx 10^{-4}Mc^2$ $v/c \sim 0.2$ for high-mass end.
 - Scaled-up CSM-interacting SNe, can be target of JWST, Roman, EUCLID
 - Smaller average velocity for lower mass star.
Even lower mass... normal collapsar.

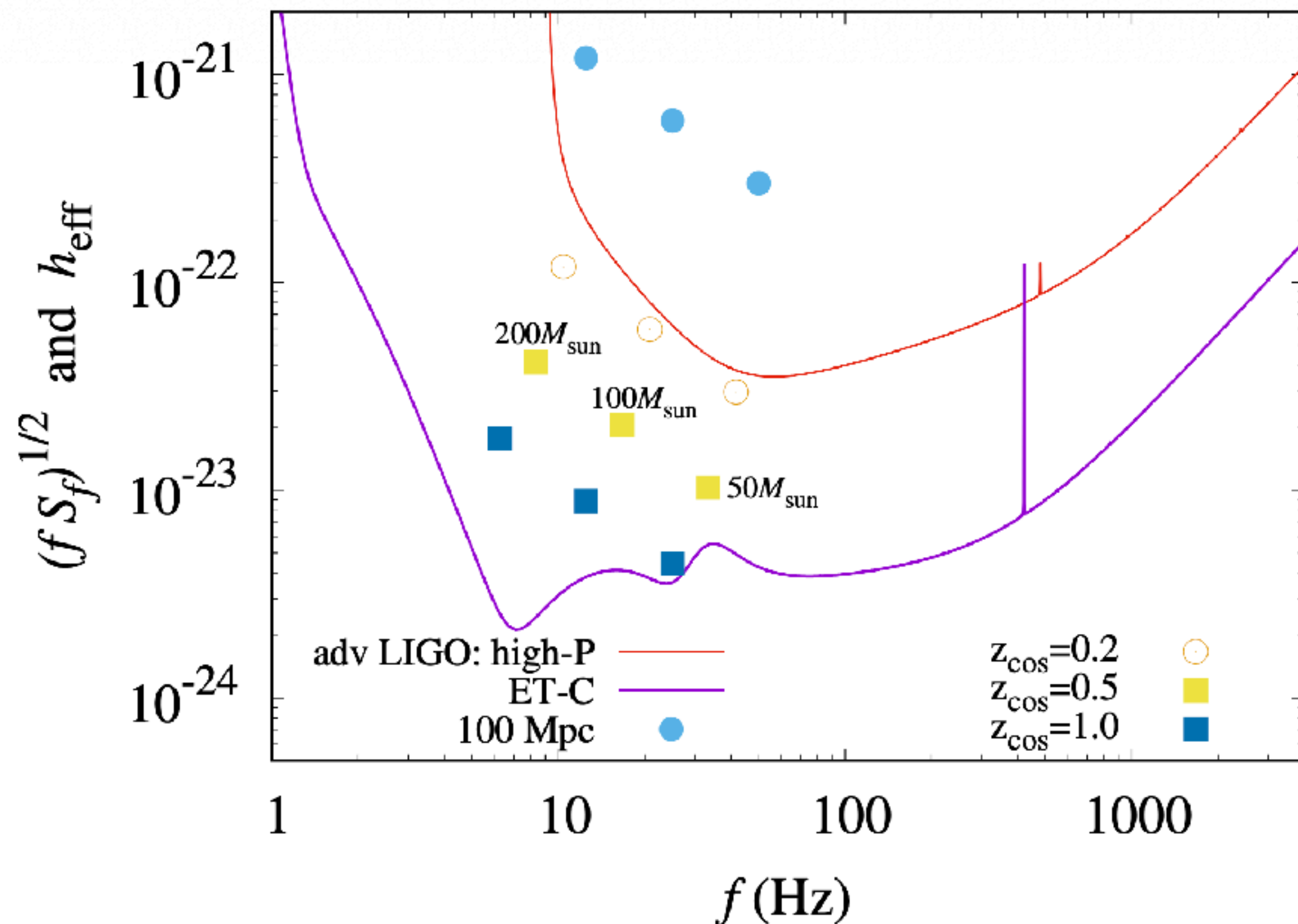
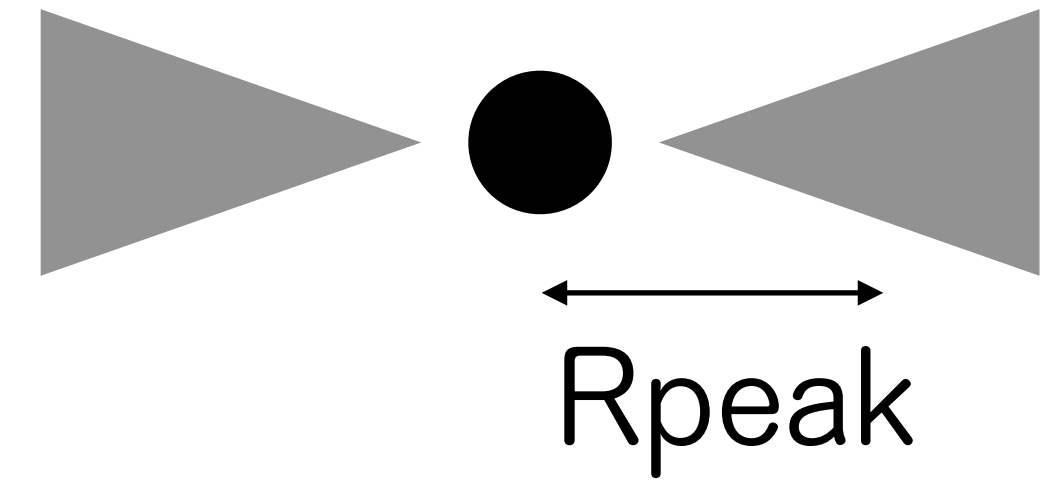
Backup slides

Lower mass (down to $10^3 M_\odot$)

Lower mass \rightarrow higher disk-to-BH mass ratio

\rightarrow Disk is unstable to non-axisymmetric deformation

\rightarrow Potential GW source



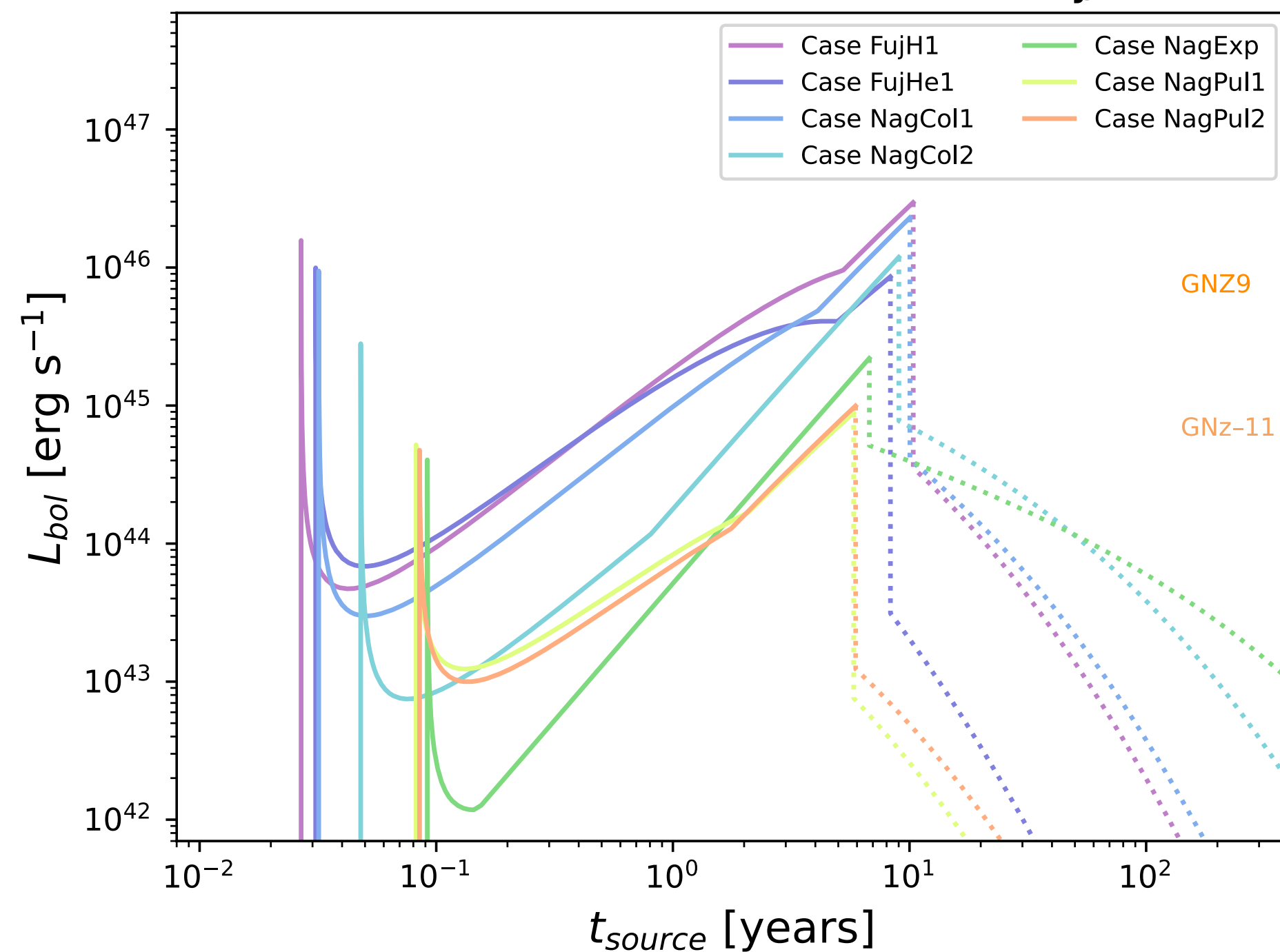
$$f_{\text{GW}} \approx 0.8 \frac{c}{\pi r_g} \left(\frac{R_{\text{peak}}}{r_g} \right)^{-3/2}$$

$\sim 30 \text{ Hz}$ for $100 M_\odot$ BH and $R_{\text{peak}} = 7 r_g$.

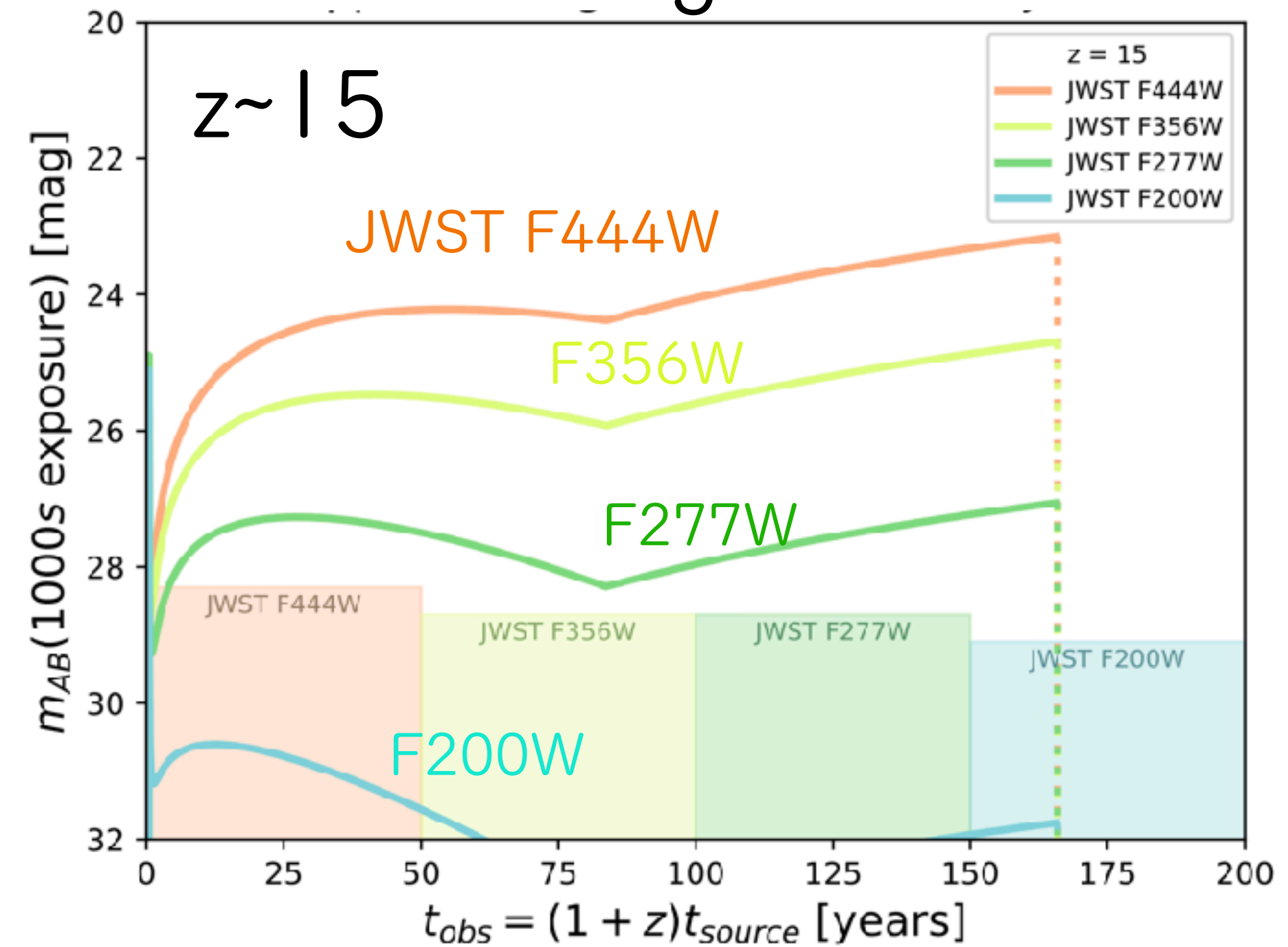
$$h_{\text{eff}} = \epsilon \frac{G M_{\text{disk}} r_g}{c^2 D R_{\text{peak}}} \approx 3 \times 10^{-22} \left(\frac{\epsilon}{0.2} \right) \left(\frac{M_{\text{disk}}}{20 M_\odot} \right) \times \left(\frac{R_{\text{peak}}}{7 r_g} \right)^{-1} \left(\frac{D}{100 \text{ Mpc}} \right)^{-1}$$

Detection prospects

Bolometric luminosity



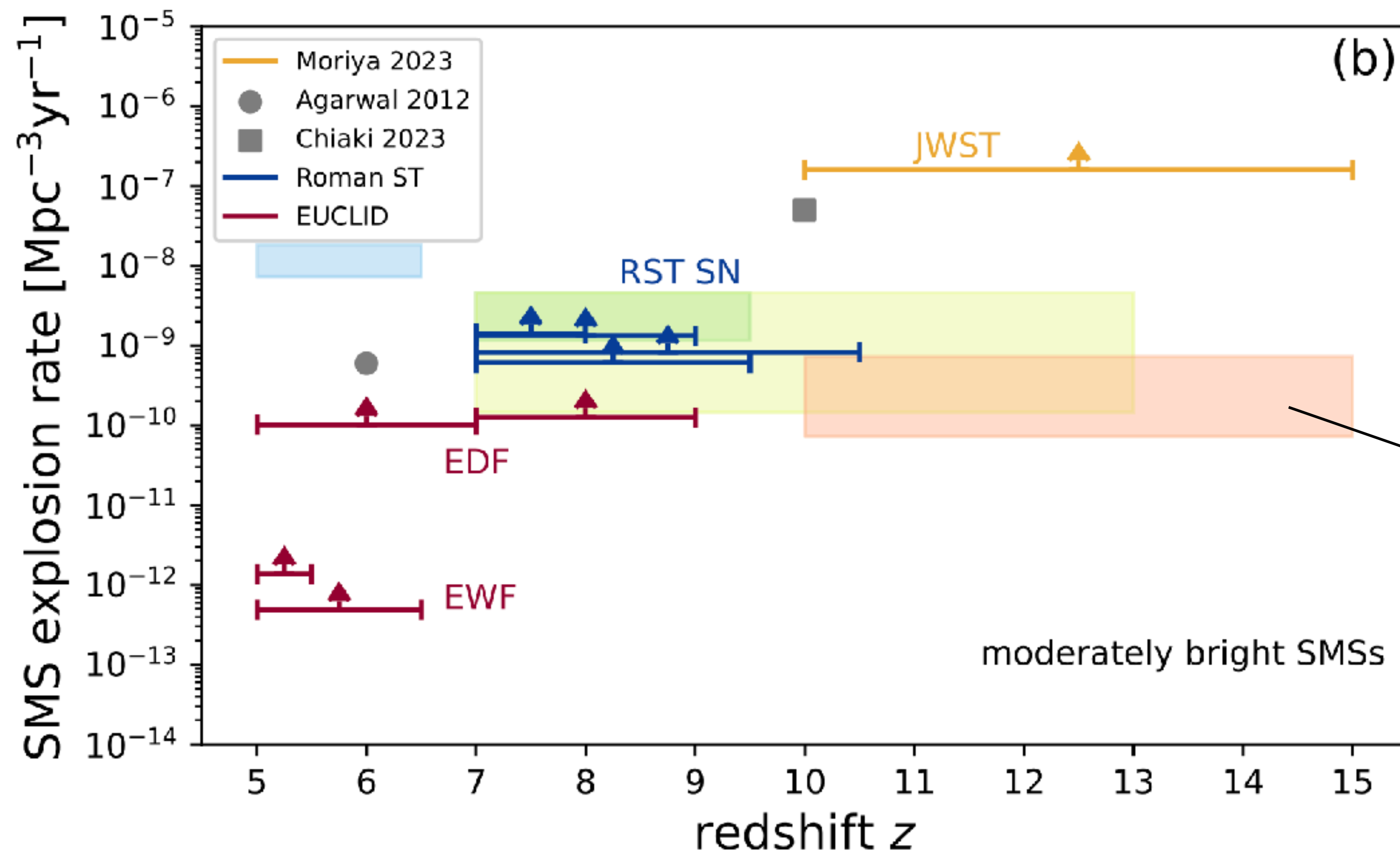
Magnitude



Even for more conservative cases ($E_{exp} \lesssim 10^{55} \text{ erg}$),
the explosion can be observed at $z \sim 15$

Detection prospects

More conservative cases (SMS mass $\sim 10^4 M_\odot$)

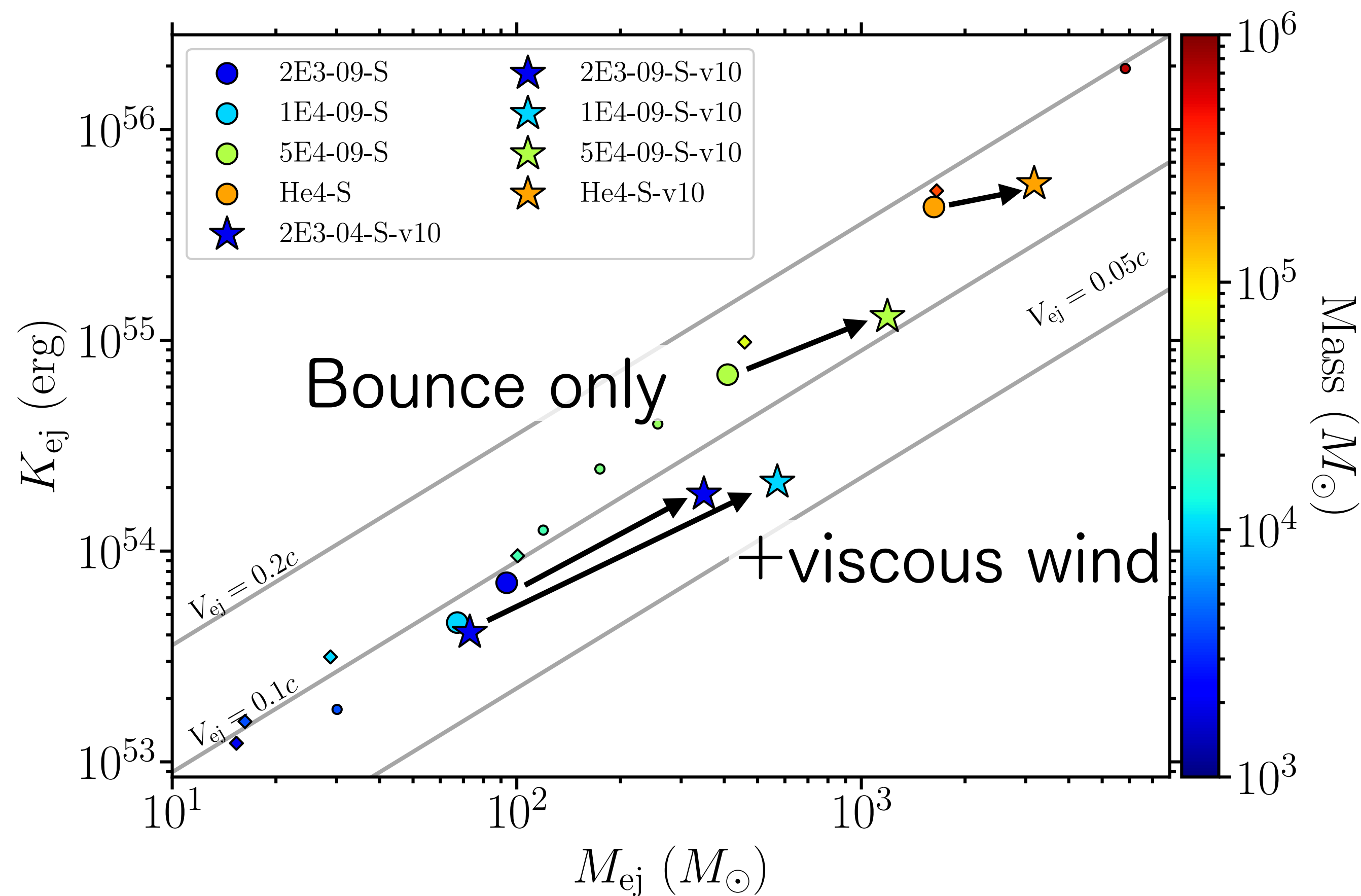


SFR (Harikane+23)

x ratio of $> 10^5 M_\odot$ stars
assuming Salpeter IMF

SMS explosion could be observed if SMS is formed at lower redshifts.

Effect of viscosity



Lower-mass star
 → less compact (large R/M)
 → envelope can have larger angular momentum
 → larger disk mass and its contribution