

Possibility of High Energy Neutrino Production in Super-Eddington Accreting Black Holes

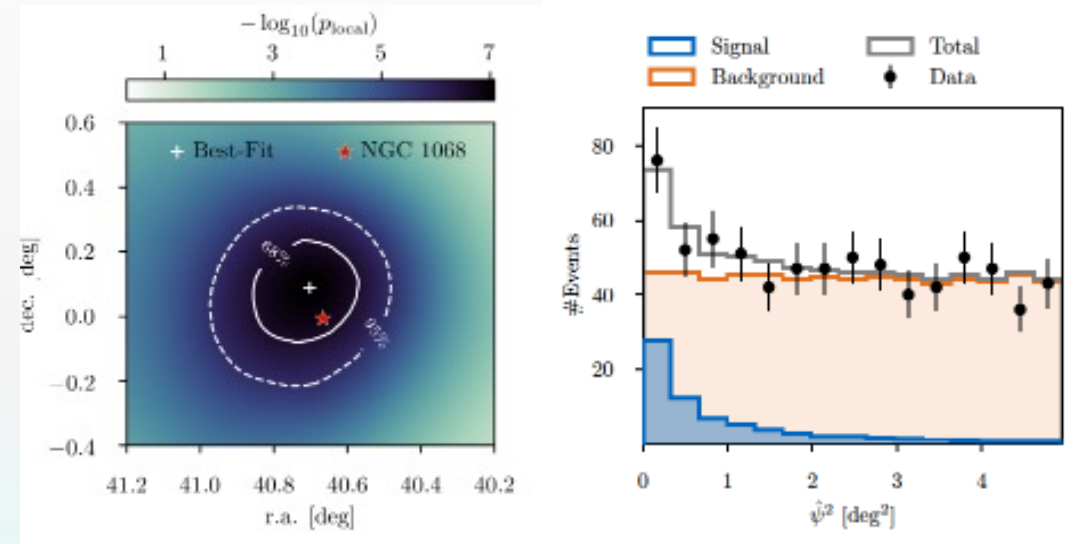
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High Energy Neutrinos from Seyfert Galaxies

- NGC 1068: $\sim 4.2\sigma$ evidence for TeV neutrino emission
 - Seyfert 2, starburst galaxy
 - no corresponding γ -ray emission
 - one of the (extinction-corrected) X-ray-brightest AGN in the sky
 - NGC 4151, NGC 7469,...
- pp or $p\gamma$ interactions in the vicinity of the black hole?
 - the photon density is high \rightarrow associated γ -ray photons would be absorbed via $\gamma\gamma \rightarrow e^-e^+$
- How to produce HE protons?
 - bright AGN \rightarrow dense accretion disk \rightarrow collisional \rightarrow non-thermal acceleration is impossible



The IceCube collaboration 2022, Science

X-ray Corona Model

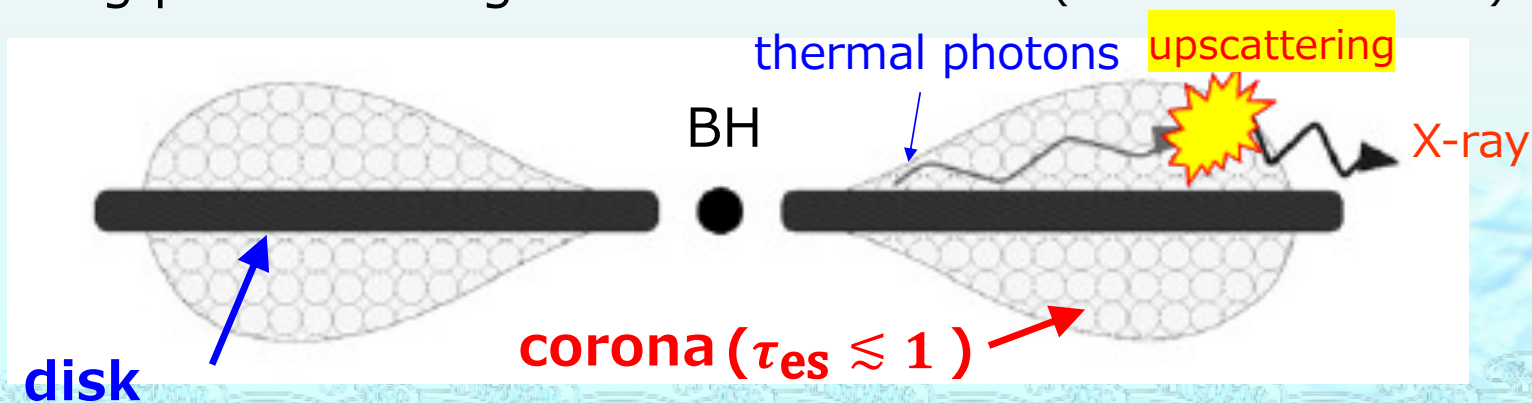
Liang & Nolan 1984; Haardt & Maraschi 1991

BH accretion flow = (1) dense, cool disk ($\sim 10^5$ K)
+ (2) tenuous, hot corona ($\sim 10^9$ K)

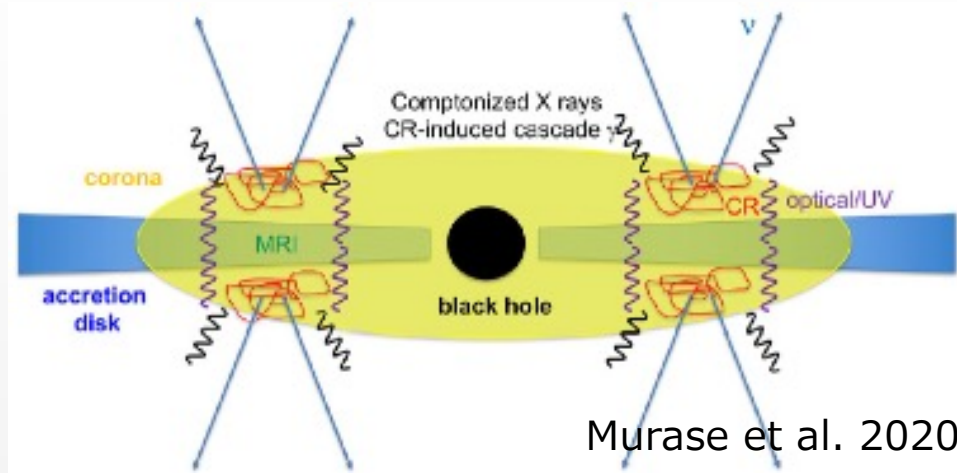
- (1) ... emit thermal soft photons (optical – UV)
- (2) ... produce hard power-law photons via Comptonization (X-ray)

→ account for the spectra of typical AGNs (and X-ray binaries)

- plasma loading : disk evaporation? (Liu et al. 1995 etc.)
- heating process: magnetic reconnection? (Liu et al. 2003 etc.)



Corona as a particle accelerator/neutrino emitter



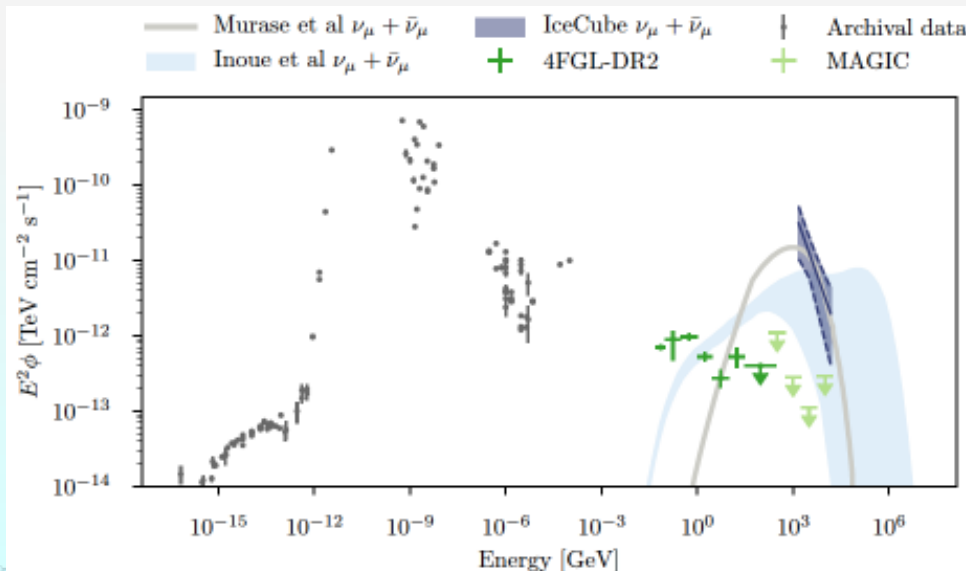
Murase et al. 2020; Kheirandish et al. 2021 etc.

- corona \sim collisionless plasma
- being turbulent due to magnetorotational instability (MRI)

→ Non-thermal particle acceleration is possible

→ HE neutrino production via $p\gamma$ interaction is expected

- fairly consistent with the neutrino emission from NGC 1068 reported by IceCube



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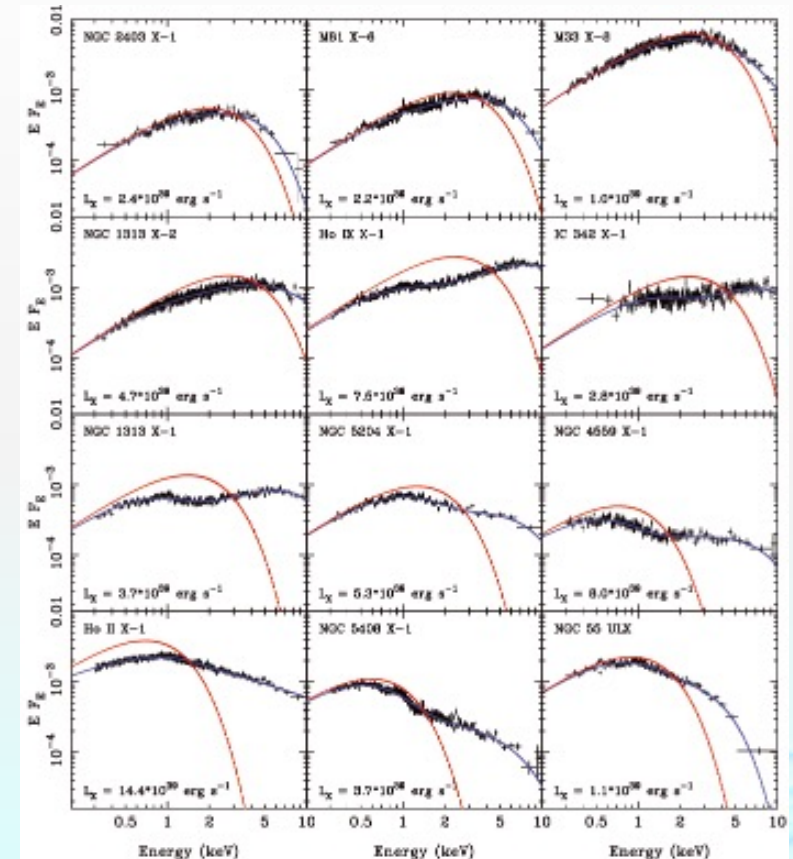
X-ray spectra of super-Eddington accreting BHs

X-ray spectra of ULXs (Gladstone+ 09)

e.g., Ultraluminous X-ray sources:

- (1) thermal + Comptonized
- (2) Comptonizing plasma (corona) should be relatively cool ($T_c \sim 10^7\text{K}$) and optically-thick ($\tau_c \sim 5\text{-}30$)

similar trends for GRS 1915+105 (Vierdayanti et al. 2010) and some narrow-line Seyfert 1 galaxies (Idogaki+ 2018 etc.)

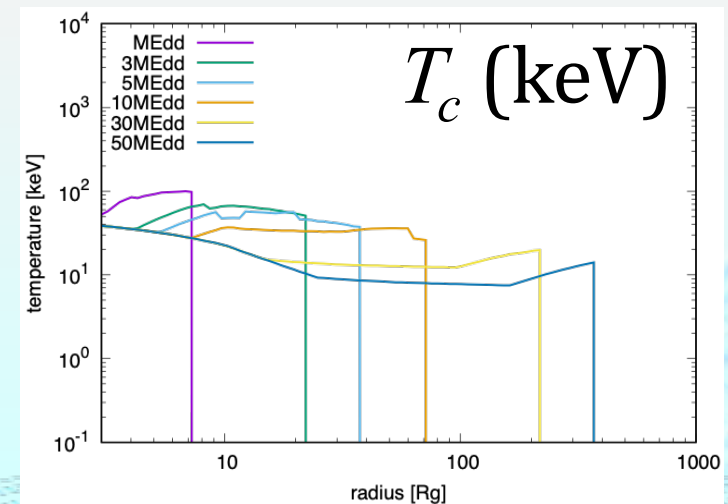
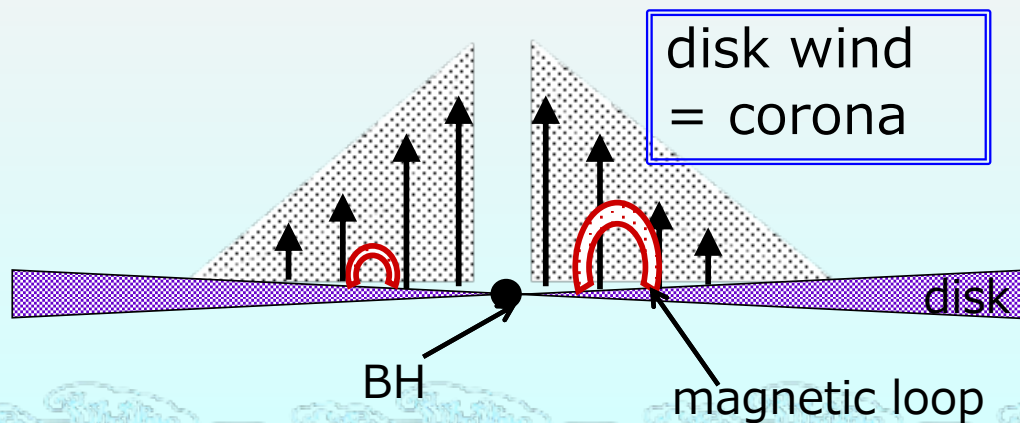


Common to super-Eddington BHs ?

wind-fed corona model for super-Eddington BHs

NK & Mineshige 2021

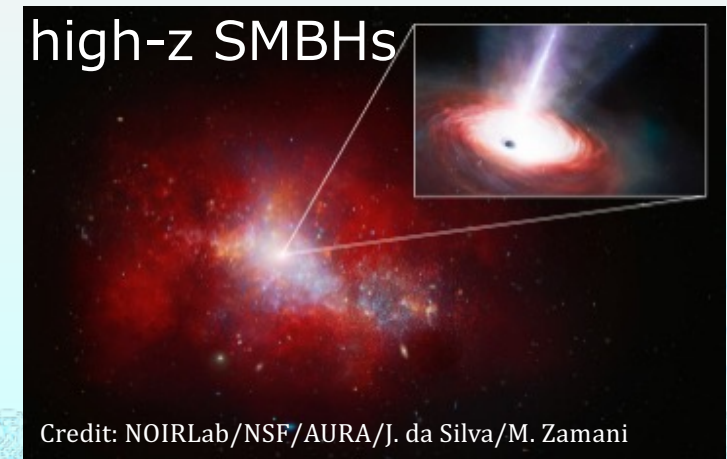
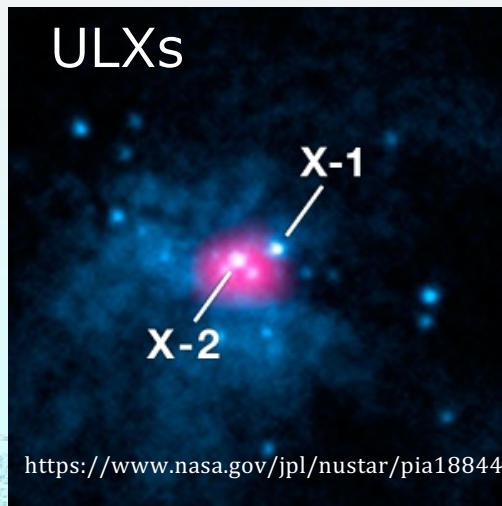
- radiation force $>$ gravity \rightarrow radiation-driven disk wind
 \rightarrow **efficient coronal plasma loading from the wind**
- disk mass loss: $\dot{M}_{\text{disk}} \propto r^s$ ($0 < s < 1$) \rightarrow **corona density**
- magnetic reconnection-heating \rightarrow **corona temperature**
(Liu et al. 2002)
- One can reproduce **optically-thick and cool corona** (\therefore **Comptonization**), as well as typical ULX spectra.



HE neutrinos from super-Eddington BHs?

- Previous works: HE neutrino production has been discussed mainly in RIAFs (LLAGNs) or coronae of standard disks
- But... super-Eddington accretors are also common:
 - ULXs, bright microquasars, Narrow Line Seyfert 1s,...
 - high-z AGNs are often supposed to be super-Eddington
 - **Their coronae are physically different**

→ **Let's investigate HE neutrinos from super-Eddington BHs**



Methods

- corona model: NK & Mineshige 2021
 - $n_c(r), T_c(r), H_c(r), U_{\text{ph}}(r) \rightarrow$ re-modeling into the one-zone corona
- neutrino spectrum calculations: Kimura et al. 2015
 - 2nd order Fermi acceleration of CR protons in the turbulent corona
 \rightarrow Fokker-Planck equation for CR p distribution function \mathcal{F}_p :

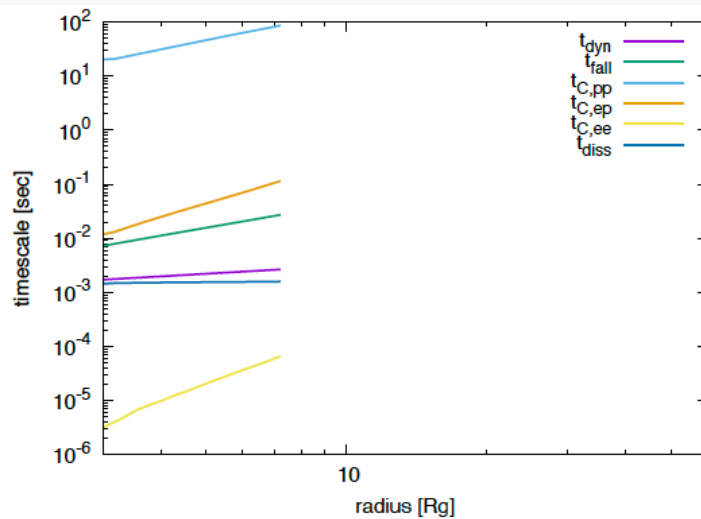
$$\frac{\partial \mathcal{F}_p}{\partial t} = \frac{1}{\varepsilon_p^2} \frac{\partial}{\partial \varepsilon_p} \left(\varepsilon_p^2 D_{\varepsilon_p} \frac{\partial \mathcal{F}_p}{\partial \varepsilon_p} + \frac{\varepsilon_p^3}{t_{p-\text{cool}}} \mathcal{F}_p \right) - \frac{\mathcal{F}_p}{t_{\text{esc}}} + \dot{\mathcal{F}}_{p,\text{inj}},$$

- consider both pp and $p\gamma$ interactions
- cooling processes: pp , $p\gamma$, p -synchrotron, Bethe-Heitler
- parameters:
 - $M_{\text{BH}} = 10M_{\odot}$ ($L_{\text{Edd}} = 1.26 \times 10^{39} \text{ erg s}^{-1}$)
 - $\dot{M} = \dot{M}_{\text{Edd}}, 3\dot{M}_{\text{Edd}}, 5\dot{M}_{\text{Edd}}, 10\dot{M}_{\text{Edd}}$

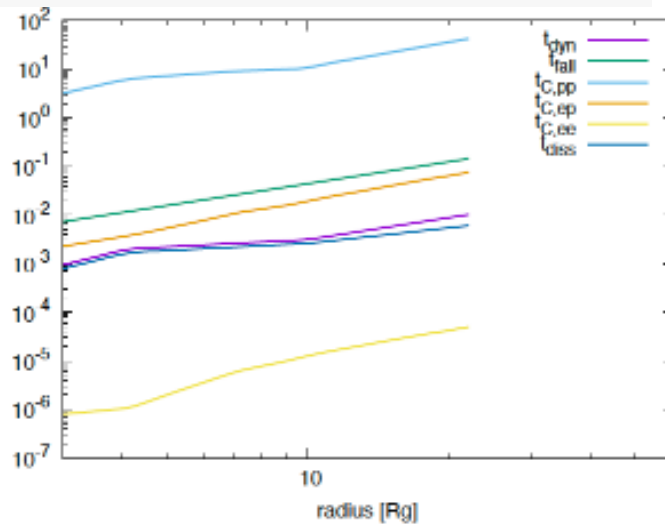
Is the corona collisionless?

$$t_{C,pp} = \frac{\sqrt{2\pi}}{n_p \sigma_T c \ln \Lambda} \left(\frac{m_p}{m_e} \right)^2 \theta_p^{3/2} : \text{Coulomb relaxation timescale}$$

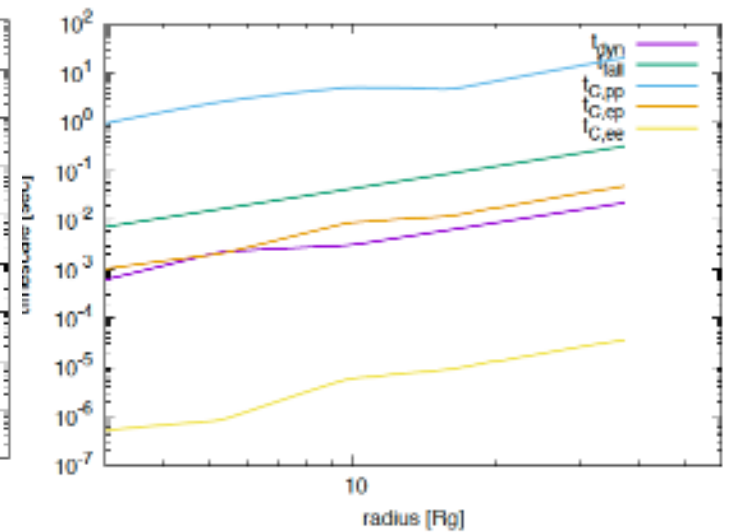
$$\dot{M} = \dot{M}_{\text{Edd}}$$



$$\dot{M} = 3\dot{M}_{\text{Edd}}$$

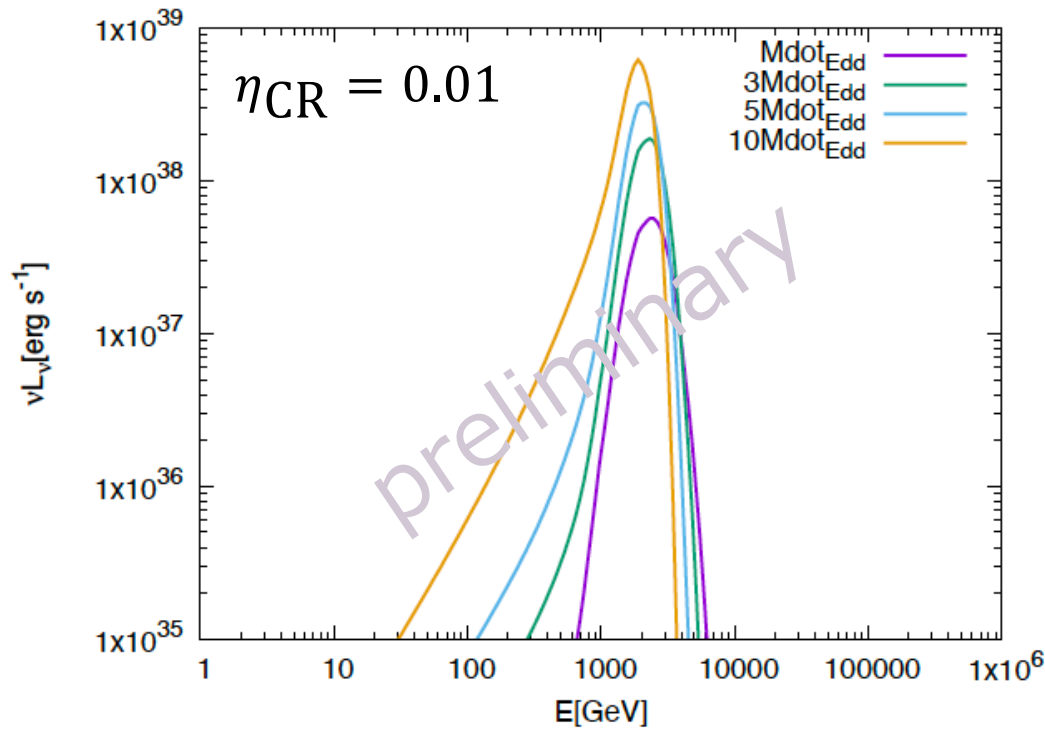


$$\dot{M} = 5\dot{M}_{\text{Edd}}$$



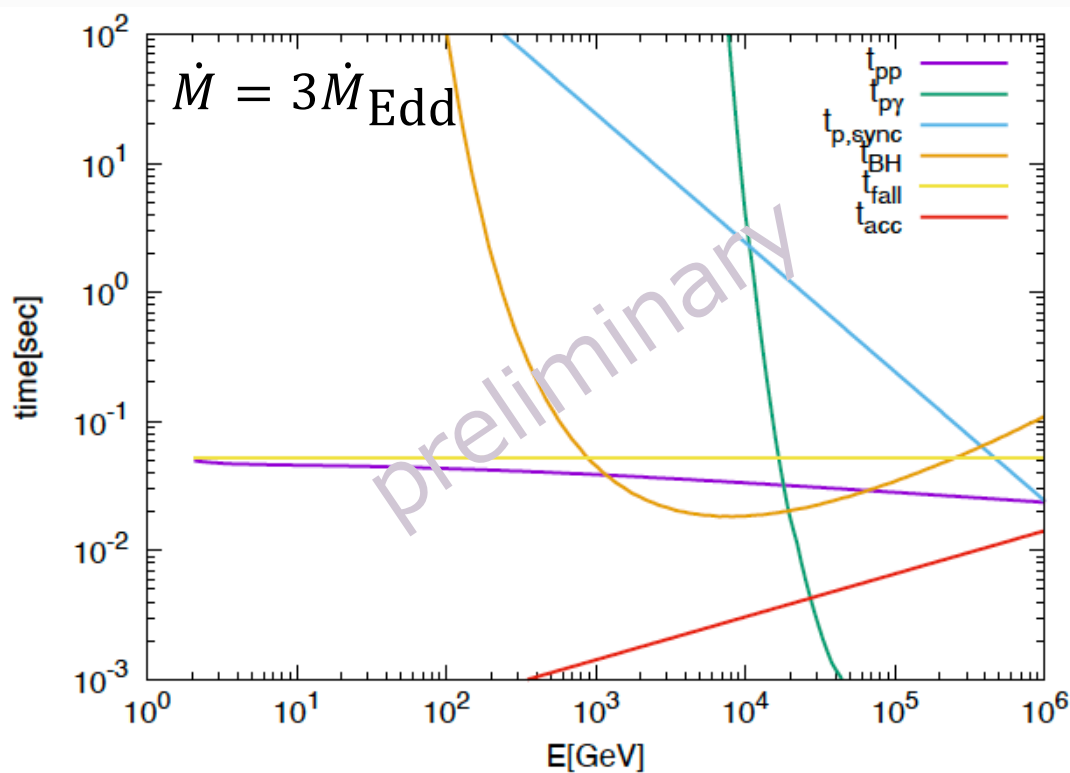
$t_{C,pp} \gg t_{\text{fall}}, t_{\text{dyn}} \rightarrow$ the corona can be regarded as collisionless

Results: HE neutrino spectra



- $\varepsilon_\nu L_\nu(\varepsilon_\nu) \gtrsim 10^{38} \text{ erg s}^{-1}$
- Sharp peak and steep cutoff at $\gtrsim \text{TeV}$
- Lower peak energy for larger \dot{M}
- The peak becomes broader for larger \dot{M}

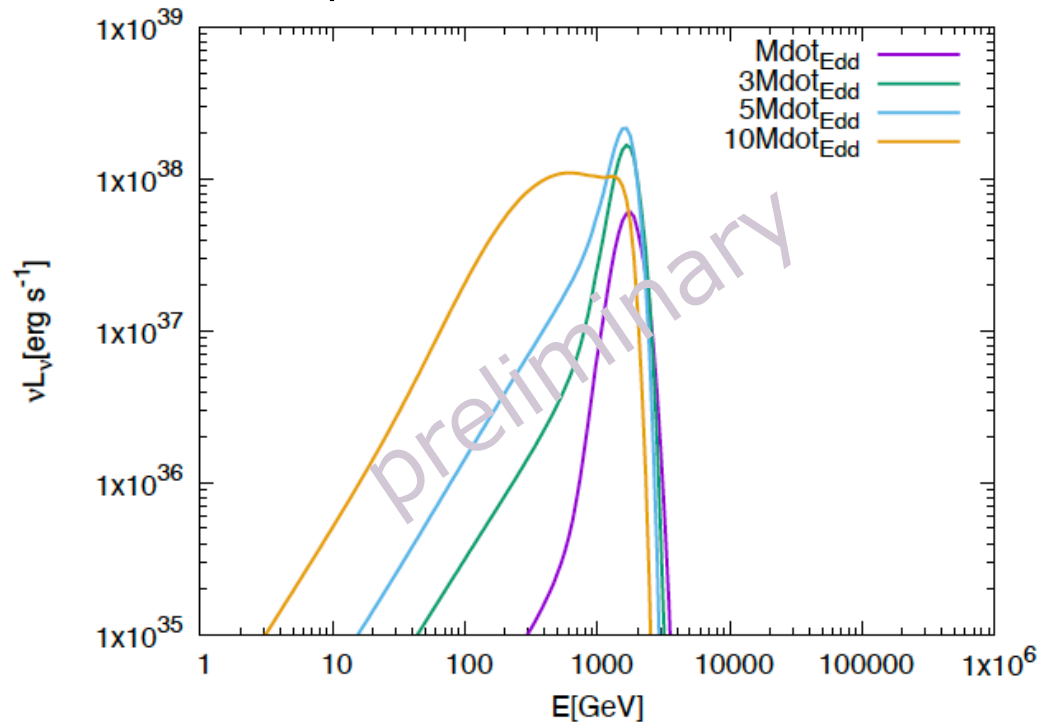
timescales



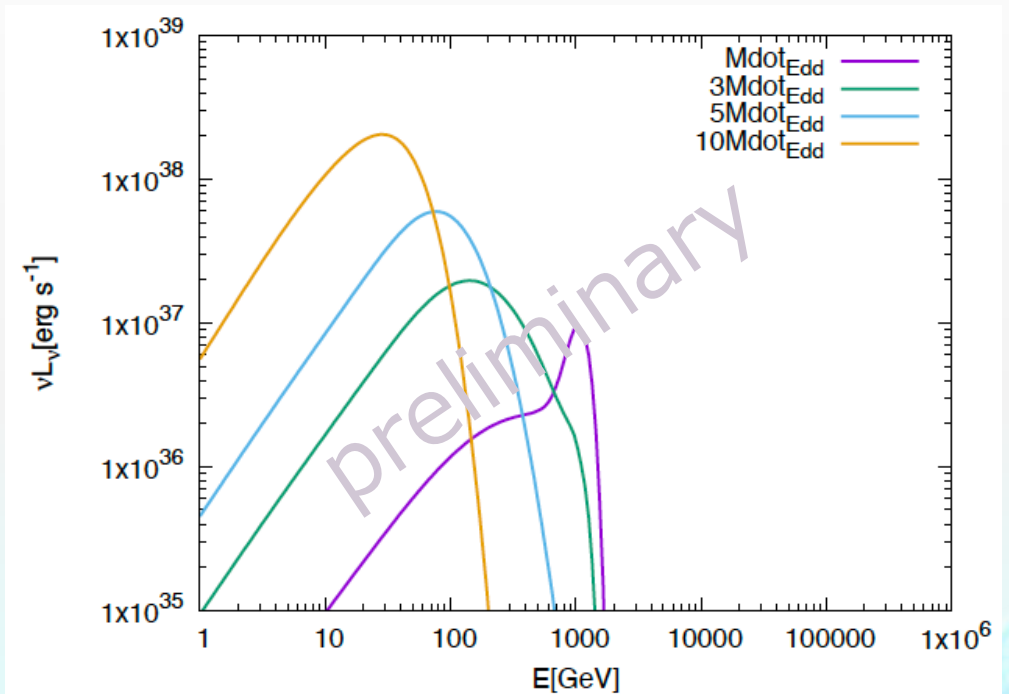
- $p\gamma$ interaction timescale becomes steeply efficient at $\gtrsim 10$ TeV
- $t_{\text{acc}} \sim t_{p\gamma}$ at $\gtrsim 10$ TeV
- due to the high corona density, pp interaction is also effective for high \dot{M}

Cases with weaker magnetic fields

1/3 of the fiducial value



1/15 of the fiducial value



Discussions

- HE neutrino spectra from ULXs have sharp cutoffs at $\sim \text{TeV}$ (or lower when B being weaker)
 - optically-thick compared to the corona in sub-Edd BHs
 \rightarrow higher X-ray photon density \rightarrow efficient $p\gamma$ cooling
 - $\dot{M} \uparrow \rightarrow$ higher photon density & denser corona
 $\rightarrow p\gamma/pp$ cooling \rightarrow lower and broader peak
 - weaker $B \rightarrow$ longer $t_{\text{acc}} \rightarrow$ lower and broader peak
- detectable?
 - fiducial: $\varepsilon_\nu L_\nu(\varepsilon_\nu) \gtrsim 10^{38} \text{ erg s}^{-1}$ at $\sim \text{TeV}$
 $\rightarrow F_\nu \sim 10^{-11} \text{ GeV cm}^{-2} \text{ s}^{-1}$ @ 10 Mpc
 - ~ 500 ULXs (including candidates) in $D < 10 \text{ Mpc}$ (Bernadich+2022) \rightarrow marginally detectable by IceCube (10 yrs) by stacking?

Summary

- **Accretion disk coronae in active galactic nuclei are the promising high energy neutrino emitters**
 - collisionless & turbulent plasma \rightarrow non-thermal particle acceleration \rightarrow HE ν production via pp or $p\gamma$
- **How about the coronae in super-Eddington BHs?**
 - ULXs, microquasars, NLS1s, high- z AGNs, etc.
 - plasma loading by radiation-driven wind \rightarrow optically-thick and cool corona (different from sub-Edd BHs)
 - collisionless condition \rightarrow OK!
 - **particle acceleration & neutrino production in ULX coronae**
 $\rightarrow \epsilon_\nu L_\nu(\epsilon_\nu) \sim 10^{38} \text{ erg s}^{-1}$ at $\sim \text{TeV}$ with a sharp cutoff
 - detectable by stacking?
 - What about other super-Edd BHs (NLS1s, high- z AGNs)?